## Embezzlement and Applications

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## Local operations:



uncorrelated outcomes
Bob $\mathrm{X}^{\mathrm{T}}|0\rangle-\sqrt{\mathrm{v}} \rightarrow \square \mathrm{b}^{\prime}$

## Entanglement:



## No free entanglement:



No free entanglement even with a catalyst:


No free entanglement even with a catalyst:


## IMPOSSIBLE

## APPROX POSSIBLE

Embezzlement of entanglement:
Any state $|\phi\rangle$ can be embezzled to any accuracy w/ some $|\psi\rangle$.
Theorem. $\forall \varepsilon>0, \forall d,|\phi\rangle_{A^{\prime} B^{\prime}} \in C^{d} \otimes C^{d}$

$$
\begin{aligned}
& \exists \mathrm{N},|\psi\rangle_{\mathrm{AB}} \in \mathrm{C}^{\mathrm{N}} \otimes \mathrm{C}^{N}, \\
& \exists \mathrm{U}, \mathrm{~V} \text { s.t. }\left(\mathrm{U}_{\mathrm{AA}^{\prime}} \otimes \mathrm{V}_{\mathrm{BB}}\right)|\psi\rangle_{\mathrm{AB}}|00\rangle_{\mathrm{A}^{\prime} B^{\prime}} \approx^{\varepsilon}|\psi\rangle_{\mathrm{AB}}|\phi\rangle_{\mathrm{A}^{\prime} \mathrm{B}^{\prime}}!
\end{aligned}
$$

van Dam \& Hayden 2002

- conceived such possibility !
- one $|\psi\rangle$ (universal)

$$
|\psi\rangle \propto \sum_{\mathrm{k}=1}^{\mathrm{N}}(1 / \mathrm{k})|\mathrm{k}\rangle_{\mathrm{A}}|\mathrm{k}\rangle_{\mathrm{B}}
$$

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## van Dam \& Hayden 2002

- conceived such possibility !
- one $|\psi\rangle$ (universal) fits all ( $\forall$ 2-party $|\phi\rangle$ of fixed d)


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## Alternative (\& obvious) embezzlement scheme

Want: $\left(U_{A^{\prime}} \otimes V_{B B^{\prime}}\right)|\psi\rangle_{A B}|00\rangle_{A^{\prime} B^{\prime}} \approx^{\varepsilon}|\psi\rangle_{A B}|\phi\rangle_{A^{\prime} B^{\prime}}$ L, Toner, Watrous 08


Given: $A^{\prime} B^{\prime},|\phi\rangle$ what $\mathrm{AB},|\psi\rangle$ ?

## Alternative (\& obvious) embezzlement scheme

Want: $\left(U_{A A A^{\prime}} \otimes V_{B B^{\prime}}\right)|\psi\rangle_{A B}|00\rangle_{A^{\prime} B^{\prime}} \approx{ }^{\varepsilon}|\psi\rangle_{A B}|\phi\rangle_{A^{\prime} B^{\prime}}$
Choose: $A=A_{1} \ldots A_{n}, B=B_{1} \ldots B_{n}, \forall i, A_{i} \sim A^{\prime}, B_{i} \sim B^{\prime}$


$$
\begin{aligned}
|00\rangle_{A^{\prime} B^{\prime}} & \otimes|\psi\rangle_{\mathrm{AB}} \quad \longrightarrow|\phi\rangle_{\mathrm{A}^{\prime} B^{\prime}} \otimes
\end{aligned} \begin{array}{|} 
& \left|\psi^{\prime}\right\rangle_{\mathrm{AB}} \approx^{\varepsilon}|\psi\rangle_{\mathrm{AB}} \text { if } \mathrm{n}=1 / \varepsilon \\
\propto \sum_{\mathrm{r}=1^{n-1}|00\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}}<\sum_{\mathrm{r}=1^{n-1}|00\rangle^{\otimes r+1}|\phi\rangle^{\otimes n-r-1}}
\end{array}
$$

## Summary of the embezzlement scheme

$$
\overbrace{C \sum_{r=1^{n-1}}|00\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}}^{|\psi\rangle_{A B}|00\rangle_{A^{\prime} B^{\prime}}} \leftrightarrow \overbrace{C \sum_{r=2^{n}}|00\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}}^{\left|\psi^{\prime}\right\rangle_{A B}|\phi\rangle_{A^{\prime} B^{\prime}}} \approx \varepsilon \quad|\psi\rangle_{A B}|\phi\rangle_{A^{\prime} B^{\prime}}
$$

- $\operatorname{dim}(A B)=\operatorname{dim}\left(A^{\prime} B^{\prime}\right)^{(1 / \varepsilon)}$ (close to optimal)
- works $\forall|\eta\rangle_{A^{\prime} B^{\prime}} \rightarrow|\phi\rangle_{A^{\prime} B^{\prime}}$ using $|\psi\rangle=C \sum_{r=1^{n-1}}|\eta\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}$
- works for multipartite $|\eta\rangle \&|\phi\rangle$
- works for other reason why $|\eta\rangle \nLeftarrow|\phi\rangle$.

References for embezzlement:

- van Dam and Hayden, 0201041
- Leung, Toner and Watrous, 0804.4118
- Leung and Wang, 1311.6842
- Connes and Stormer, J functional analysis 28, 187 (1978)
$\infty$-dim generalization, self-embezzlement:
- Haagerup, Scholz and Werner (in preparation)
- Cleve, Liu, Paulsen, 1606.05061
- Cleve, Collins, Liu, Paulsen, 1811.12575

Mismatched descriptions of what to embezzle:

- Steurer, Dinur, Vidick, 1310.4113

Open problems on embezzlement:

1. van Dam - Hayden sch
unitaries depends on $|\phi\rangle$ bipartite states

LTW scheme
catalyst depends on $|\phi\rangle$
unitaries independent of $|\phi\rangle$
multi-partite states

LTW scheme can use a universal catalyst: tensor product of catalysts for an $\varepsilon$-net of target states and a fixed initial state.

For embezzlement of multipartite state, is there a more efficient universal catalyst?
2. L, Wang 2013 showed that finite-dim embezzlement catalyst is essentially unique for universal embezzlement in the bipartite setting. Same for multi-partite setting?

## Outline:

1. Embezzlement
2. Approximate violation of conservation laws \& macroscopically controlled coherent operations
3. Finite Bell inequality that cannot be violated maximally with finite amount of entanglement
4. Quantum reverse Shannon theorem
Local operations

Superselection rules
Entanglement
Conserved quantities (charge, spin etc)

SSR: Restricted Hamiltonian or unitary that are block-diagonal
"Block index" is conserved

Local operations $\longrightarrow$ Superselection rules
Entanglement $\qquad$ Conserved quantities (charge, spin etc)

Embezzlement
Generic recipe to approx an otherwise forbidden transformation

Suppose $|\eta\rangle \nLeftarrow|\phi\rangle$, say, because $|\eta\rangle,|\phi\rangle$ contain different amount of a conserved quantity.
Cyclic permutation conserves the quantity (allowed).
Using $|\psi\rangle=C \sum_{r=1}{ }^{n-1}|\eta\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}$ one can perform

$$
\begin{aligned}
|\psi\rangle|\eta\rangle & =C \sum_{r=1}^{n-1}|\eta\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}|\eta\rangle \\
& \rightarrow C \sum_{r=1} 1^{n-1}|\eta\rangle^{\otimes r+1}|\phi\rangle^{\otimes n-r-1}|\phi\rangle \approx^{\varepsilon}|\psi\rangle|\phi\rangle
\end{aligned}
$$

and "violate" the conservation law!

Furthermore, the approx transformation is coherent, and can be performed / not in superposition.

Conditioned on $1^{\text {st }}$ register being $|1\rangle$, apply $|\psi\rangle|\eta\rangle \rightarrow^{\varepsilon}|\psi\rangle|\phi\rangle$

$$
(\mathrm{a}|0\rangle|\gamma\rangle+\mathrm{b}|1\rangle|\eta\rangle)|\psi\rangle \leftrightarrow^{\varepsilon}(\mathrm{a}|0\rangle|\gamma\rangle+\mathrm{b}|1\rangle|\phi\rangle)|\psi\rangle
$$

Thus $|\psi\rangle$ makes the superselection rule irrelevant.

## Application: macroscopically-controlled gates

e.g., $|0\rangle_{\mathrm{s}}$ : spin down (ground state)
$|1\rangle_{\mathrm{s}}$ : spin up (excited state)
"X gate": a $|0\rangle_{\mathrm{S}}+\mathrm{b}|1\rangle_{\mathrm{S}} \leftrightarrow \mathrm{a}|1\rangle_{\mathrm{S}}+\mathrm{b}|0\rangle_{\mathrm{S}}$ but $|0\rangle_{\mathrm{S}} \psi|1\rangle_{\mathrm{S}}$
Allowed: $|r\rangle_{L}|0\rangle_{S} \leftrightarrow|r-1\rangle_{L}|1\rangle_{S}$ where $|k\rangle_{L}=k$-photon state in laser beam.

But changes in \# photon in laser beam decoheres the spin.
Solution: use $|\psi\rangle_{L}=\sum_{r=1}{ }^{n-1}|r\rangle_{L}$ :

$$
\begin{aligned}
&|\psi\rangle_{\mathrm{L}}\left(\mathrm{a}|0\rangle_{\mathrm{S}}+\mathrm{b}|1\rangle_{\mathrm{S}}\right) \leftrightarrow \underbrace{\sum_{r=1^{n-1}}|r-1\rangle_{\mathrm{L}}}_{\approx|\psi\rangle_{\mathrm{L}}} a|1\rangle_{\mathrm{S}} \\
&\longrightarrow \underbrace{\sum_{r=1}^{\mathrm{n}-1} \mid \mathrm{r}+1}_{\approx|\psi\rangle_{\mathrm{L}}}\rangle_{\mathrm{L}} \mathrm{~b}|0\rangle_{\mathrm{S}} \\
& \approx|\psi\rangle_{\mathrm{L}}\left(\mathrm{a}|1\rangle_{\mathrm{S}}+\mathrm{b}|0\rangle_{\mathrm{S}}\right) \text { nearly coherent } \mathrm{X} \text { gate }
\end{aligned}
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$$
|\psi\rangle_{L}\left(\mathrm{a}|0\rangle_{\mathrm{S}}+\mathrm{b}|1\rangle_{\mathrm{S}}\right) \leftrightarrow \approx|\psi\rangle_{\mathrm{L}}\left(\mathrm{a}|1\rangle_{\mathrm{S}}+\mathrm{b}|0\rangle_{\mathrm{S}}\right)
$$

In the lab, we use the coherent state $|\psi\rangle_{L} \propto \sum_{r=1}^{n-1} \alpha r / \sqrt{ }(r!)|r\rangle_{L}$ !

## Local operations <br> 

Entanglement $\qquad$ Conserved quantities (charge, spin etc)

Embezzlement
Generic recipe to approx an otherwise forbidden transformation

Principle: use catalyst to introduce a large uncertainty of the conserved quantity to enable approximately violation of conservation law
$|\psi\rangle \propto \sum_{r=1}{ }^{n-1}|00\rangle^{\otimes r}|\phi\rangle^{\otimes n-r}$
Uncertainty in \# of copies of $|00\rangle$ vs $|\phi\rangle$

$$
|\psi\rangle_{L} \propto \sum_{r=1^{n-1}}|r\rangle_{L}
$$

Uncertainty in photon \#

## More on conservation laws

Kitaev, Mayers, \& Preskill (0310088) investigated (in response to Popescu) if superselection rules (SSR) help quantum crypto by restricting adversarial behavior:
superposition of diff charges possible if a charge reservoir (a condensate ~ catalyst) is accessible, and SSR cannot enhance quantum cryptography.

Bartlett, Rudolph, and Spekkens (0610030) generalized the above, by connection to "reference frames" which are like the catalyst in this talk.

Embezzlement $\rightarrow$ arbitrary unitary despite SSR ?
Latter solved by Popescu, Sainz, Short, Winter (1804.03730)
1-party result, does not give embezzlement ...

## Outline:

1. Embezzlement
2. Approximate violation of conservation laws \& macroscopically controlled coherent operations
3. Finite Bell inequality that cannot be violated maximally with finite amount of entanglement
4. Quantum reverse Shannon theorem

Embezzlement based Bell inequality that cannot be violated maximally with finite amount of entanglement

Embezzlement based nonlocal game that cannot be played optimally with finite amount of entanglement

Non-closure of quantum correlations via embezzlement

References:

- Leung, Toner, Watrous (0804.4118)
- Ji, Leung, Vidick (1802.04926)
- Coladangelo (1904.02350)


## Nonlocal games



Players can coordinate before the game noncommunicating once the game starts

## Nonlocal games

Goal: max prob(winning) Does entanglement help?


## e.g., GHZ game

$$
k=3, q \in R \in\left\{\begin{array}{l}
(x, x, x),(y, y, x) \\
(y, x, y),(x, y, y)
\end{array}\right\}
$$



## e.g., GHZ game

$$
k=3, q \in R\left\{\begin{array}{l}
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$$



Without entanglement, winning prob $\leq 3 / 4$.
With a GHZ state, each party measures $\sigma_{x / y}$, winning prob $=1$ !
"Rigid" - unique optimal strategy (mod local isometries), robust.

| Nonlocal games | Bell experiments |
| :--- | :--- |
| Questions to players | Measurement settings |
| Answers from players | Measurement outcomes |
| Prob(win) $\rightarrow$ payoff function | Bell inequality |
| Classical strategy | Local hidden variable model |
| shared randomness |  |

Entangled strategy has strictly higher winning prob than classical

Violation of Bell inequality

## Why nonlocal games?

Computational complexity -
Effects of entanglement in interactive proof systems
Physics -
QM vs local hidden variable model
Crypto -
QKD via rigidity (uniqueness of optimal solution)

Here: how much entanglement is needed to win optimally?
Conjecture since 2009: for some games with finitely many Q\&A, more entanglement always strictly increases the winning prob.

## Proofs:

Numerical evidence: Pal-Vertesi 09 (I3322)
Existential: Slofstra (+Vidick) 17, Dykema-Prakash-Paulsen 17
Robust: dim lower bound vs deviation from optimal
Explicit: Ji, L, Vidick 18, Coladangelo-Stark 18, Coladangelo 19

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Explicit: Ji, L, Vidick 18, Coladangelo-Stark 18, Coladangelo 19
JLV18, C19 (elementary proof + physical understanding + exponentially stronger dim bound):
Turn a game from L, Toner, Watrous 08 into nonlocal games LTW game has 2 parties, each with 3 -dim quantum question and 2-dim quantum answer, based on embezzlement.
JLV18: 3 parties, each with 12 questions and 8 or 4 answers
C19: 2 parties, 5 or 6 questions and 3 answers each

## The possibility \& impossibility of embezzlement

Qualitative no-go thm: $|\psi\rangle_{A B}|00\rangle_{A^{\prime} B^{\prime}} \nLeftarrow|\psi\rangle_{A B}|\phi\rangle_{A^{\prime} B^{\prime}}$
Possibility of approximate embezzlement:
poor "continuity" of no-go thm
Poor continuity still limits how well one can embezzle
-- high accuracy requires more dim in the catalyst !

## Limits to embezzlement of entanglement

Theorem (from Fannes ineq):

$$
\begin{aligned}
& \text { If } \varepsilon>0,|\phi\rangle_{A^{\prime} B^{\prime}} \in C^{d} \otimes C^{d},|\psi\rangle_{A B} \in C^{N} \otimes C^{N}, \\
& \text { and } \exists \mathrm{U} \text {, V s.t. }\left\langle\left.\psi\right|_{A B}\left\langle\left.\phi\right|_{A^{\prime} B^{\prime}}\left(U_{A A^{\prime}} \otimes V_{B B^{\prime}}\right) \mid \psi\right\rangle_{A B} \mid 00\right\rangle_{A^{\prime} B^{\prime}} \geq 1-\varepsilon
\end{aligned}
$$

then $\varepsilon \geq 8[\mathrm{E}(|\phi\rangle) /(\log \mathrm{N}+\log \mathrm{d})]^{2}$

## "Nonlocal games" with quantum Qns \& Ans


$Q_{1}, \cdots, Q_{k}, A_{1}, \cdots, A_{k}$ : quantum sys
Initial state on $\mathrm{R}_{1}, \cdots, \mathrm{Q}_{\mathrm{k}}$ pure

2-outcome POVM meas /
known to players

Embezzlement game that cannot be won with finite entanglement
2-player
LTW08


Initial state on RXY:

$$
\frac{1}{\sqrt{2}}-\left[|0\rangle|00\rangle+|1\rangle \frac{(|11\rangle+|22\rangle)}{\sqrt{2}}\right]_{R X Y}
$$

Possible strategy:
if $X(Y)$ in $\operatorname{span}\{|1\rangle,|2\rangle\}$ then reverse-embezzle $|11\rangle+|22\rangle \rightarrow|11\rangle$.
Winning prob $\rightarrow 1$.

No other way to win: direct proof prob(winning) $<1-\log ^{-2} \operatorname{dim}(E)$

## Turning embezzlement game into a nonlocal game:

Regev and Vidick (1207.4939):
Referee's state $R$ and answers AB classical
Questions XY remain quantum
Difficulty: distributing the initial state

## Turning embezzlement game into a nonlocal game:



## Turning embezzlement game into a nonlocal game (JLV18):



1. referee $\rightarrow$ 3rd player Victor initial state on XYR $\rightarrow$ shared entanglement

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1. referee $\rightarrow$ 3rd player Victor initial state on XYR $\rightarrow$ shared entanglement
2. replace measurement by a rigidity test of the GHZ state
3. Real referee $R$ uses questions+winning conditions to enforce correct initial state \& evolution.

## Resulting game:

3-player, 12 questions each
3-bit answer from Victor, 2 bits from Alice \& Bob each

1. Suffices for Victor, Alice, Bob to share entangled state with $3, \mathrm{O}(1 / \varepsilon), \mathrm{O}(1 / \varepsilon)$ qubits to win wp > 1- $\varepsilon$.
2. Necessary for the entangled state to have at least $\Omega\left(\varepsilon^{-1 / 32}\right)$ qubits (exp that of Slofstra-Vidick-17).
3. Verification of increasing dim based on "1 test".

## Turning embezzlement game into a nonlocal game (C19):

Goal: forcing the players to convert $\frac{(|11\rangle+|22\rangle)}{\sqrt{2}}$ into $|11\rangle$
Referee conducts one of 3 possible games $G_{1}, G_{2}, G_{3}$ :
$\mathrm{G}_{1}$ can only be won close-to-optimally with a state close to

$$
\frac{|00\rangle+|11\rangle+|22\rangle}{\sqrt{3}}
$$

$\mathrm{G}_{2}$ can only be won close-to-optimally with a state close to

$$
\frac{|00\rangle+\sqrt{ } 2|11\rangle}{\sqrt{3}}
$$

$\mathrm{G}_{3}$ ensures that the states above live in the same Hilbert space!

## Open problems on nonlocal games \& quantum games:

1. Is I3322 a game that will proof the conjecture in 2009 ?
2. Are there other physical reasons for requiring unbounded amount of entanglement to optimize Bell ineq violation?
3. The embezzlement (quantum) game shows: LU-assisted by entanglement is not a closed set for 3 input and 2 output dimensions? What is the min dim for non-closure?
4. For LU-assisted by entanglement, if we allow approximations, is there a bound on the sufficient entanglement that depends only on the input/output dims?
5. For nonlocal games, is there a bound on entanglement independent of the game but depends only on the approx and the \# qns and ans?
6. Applications of the embezzlement game or nonlocal game derived from it? e.g., J LV18, C19 games verify increasing dims.

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## Quantum reverse Shannon theorem:

Quantum Shannon theorem:
Simulate noiseless channel using noisy channel at the best rate
Capacity $\mathrm{C}(\mathrm{N})=$ \# qubits sent per channel use


Quantum Reverse Shannon theorem:
Simulate noisy channel $N$ using noisless channel at the rate $1 / C(N)$


Why?? If true, any channel N can simulate any other channel M at optimal rate - $\mathrm{C}(\mathrm{N}) / \mathrm{C}(\mathrm{M})$ ( N simulates I which simulates M ) so any channel N is characterized by $\mathrm{C}(\mathrm{N})$ !

## Quantum reverse Shannon theorem:

- Bennett, Devetak, Harrow, Shor, Winter (0912.5537)
- Berta, Christandl, Renner (0912.3805 - alternative proof)

Proved for tensor-product inputs when entanglement is free but different inputs consume different amount of entanglement so a superposition of inputs is decohered.

Idea: embezzle away the left-over entanglement to keep the coherence of a superposition of inputs!

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