CO781 / QIC 890:

Theory of Quantum Communication

Prelude to topics 4-6

Quantum entropies and some important properties

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This lecture and set of slides are highlights from the lecture notes for the F2016 offering, lecture 10 (relinked in the F2020 website).

Please read through the notes offline.

<u>Important conventions (I)</u>

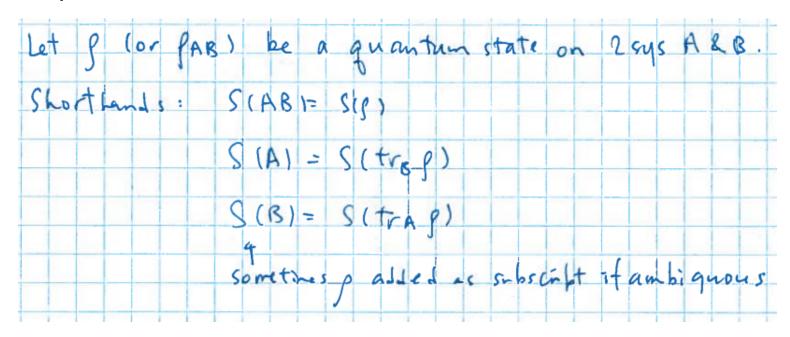
Recall the von Neumann entropy S(p) is the Shannon entropy of the spectrum.

Let
$$\int_{\substack{\text{spectral} \\ \text{decomp}}}^{} = U D U^{\dagger}, \quad \log \int_{\substack{\text{decomp}}}^{} = U (\log D) U^{\dagger}$$

So:
$$S(\rho) := -\text{tr } D \log D = -\text{tr} \left[\begin{array}{c} \lambda_1 & 0 \\ \lambda_2 & \log \lambda_2 \end{array} \right] \left[\begin{array}{c} \log \lambda_1 & 0 \\ \log \lambda_2 & \log \lambda_d \end{array} \right]$$

= -tr
$$U D U^{\dagger} U (log D) U^{\dagger}$$
 = -tr $\rho log \rho$ cyclic property of trace

Important conventions (II)



Definition: Conditional entropy (on a state on AB)

S(A|B) := S(AB) - S(B)

Definition: Quantum mutual information (on a state on AB)

S(A:B) := S(A) + S(B) - S(AB)

Similarities:

Quantum

N/A

Classical

 $H(X|Y) := \sum_{y} p(y) H(X|Y=y)$

$$S(A|B) := S(AB) - S(B)$$

chain rule:

S(AB) = S(B) + S(A|B)

H(X|Y) = H(XY) - H(Y)

H(XY) = H(Y) + H(X|Y)

$$S(A:B) := S(A) + S(B) - S(AB)$$

Klein's inequality (see notes)

$$S(A:B) \geqslant 0$$
 "=" iff $\rho = \rho_A \otimes \rho_B$

subadditivity:

$$S(AB) \leq S(A) + S(B)$$

conditioning reduces entropy:

$$S(A:B) = S(B) - S(B|A)$$

so,
$$S(B) \geqslant S(B|A)$$

I(X:Y) = H(X) + H(Y) - H(XY)

$$I(X:Y) \geqslant 0$$
 "=" iff X,Y indep

$$H(XY) \leq H(X) + H(Y)$$

conditioning reduces entropy:

$$H(Y) \geqslant H(Y|X)$$

sym(A < -> B)

Examples:

$$S(AB) - S(B)$$
 $S(A) + S(B) - S(AB)$

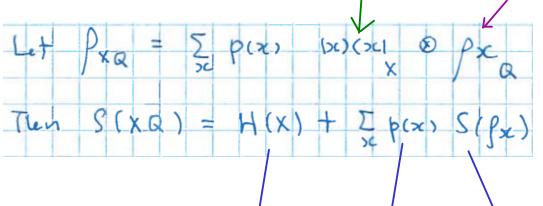
ρ	S(A)	S(B)	S(AB)	S(A B)	S(A:B)	
(I) A S TOXOLE	1	0	1	1	0	
10×0 (12 0 (2))B	0	1	1	0	0	
+ 1 100X001	1	1	1	0	1	
克(100)tup) (亚X里)		1	0	-1	2	

- * have B, do we know everything about A? not if we measure A in X-basis ...
- * quantum conditional entropy can be negative NOT a convex combination of entropies ...
- * S(AB) can be \geq , = , \leq S(B)

Special properties for classical-quantum systems

1. Entropy of a classical-quantum system:

each state on the classical system X is projector onto a basis states
each state on the quantum system Q is labelled by one x, arbitrary otherwise



same format as one drawn of an ensemble of states

$$\mathcal{E} = \{ p(sc), psc \}$$

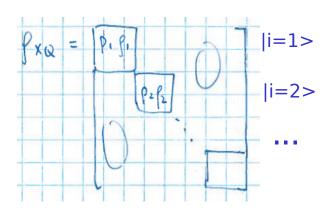
Shannon entropy of classical part

/ von Neumann entropy of Q given x average over X

* Special case in which
$$S(Q|X) = S(XQ) - S(X) = \sum_{SC} p(X) S(f_{X})$$
 where conditioning has the same interpretation as the classical case.

on X a classical system

Proof:
$$p_{XQ} = \sum_{i} p_{i} |_{i \times i} |_{x} \otimes p_{i}|_{Q}$$



If
$$\rho_1$$
 has eigenvalues λ_{11} , λ_{12} ,....
 ρ_2 has eigenvalues λ_{21} , λ_{22} ,...

then ρ has eigenvalues $\rho_1 \lambda_{11}$, $\rho_1 \lambda_{12}$,...
 $\rho_2 \lambda_{21}$, $\rho_2 \lambda_{22}$...

$$S(p) = -\sum_{i j} P(i \lambda_{i}) \log (P(i \lambda_{i}))$$

$$= -\sum_{i j} P(i \lambda_{i}) (\log P(i + \log \lambda_{i}))$$

$$= -\sum_{i j} P(i \lambda_{i}) \log P(i - \sum_{i j} P(i \lambda_{i}) \log \lambda_{i})$$

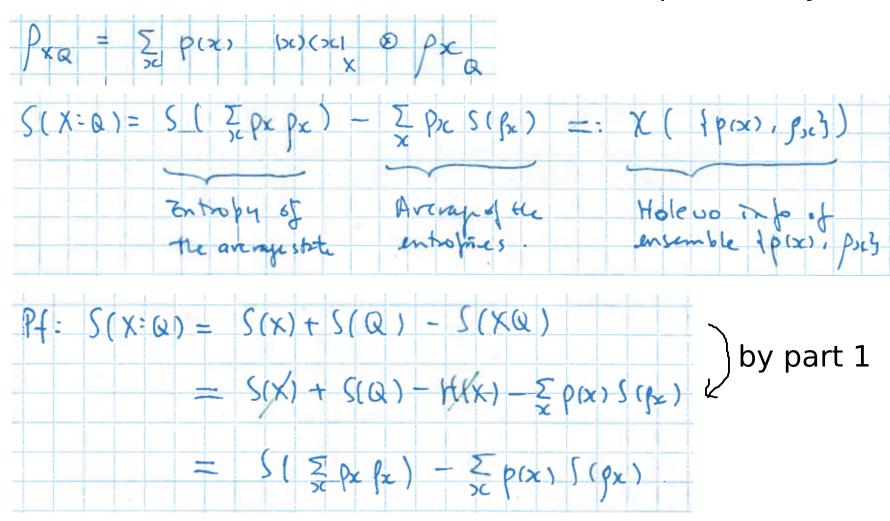
$$= -\sum_{i j} P(i \lambda_{i}) \log P(i - \sum_{i j} P(i \lambda_{i}) \log \lambda_{i})$$

$$= -\sum_{i j} P(i \log P(i + \sum_{i j} P(i \log P(i)))$$

$$= -\sum_{i j} P(i \log P(i + \sum_{i j} P(i \log P(i)))$$

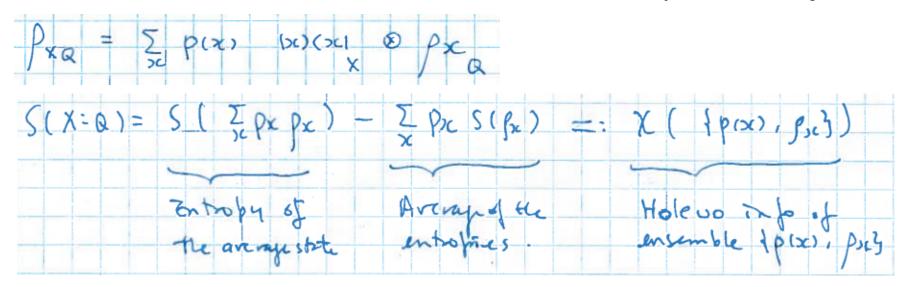
Special properties for classical-quantum systems

2. Quantum mutual information of a classical-quantum system:



Special properties for classical-quantum systems

2. Quantum mutual information of a classical-quantum system:



Corollary for arbitrary state

Mixing increases entropy (concavity of entropy)

If
$$p_{x} > 0$$
, $\sum_{x} p_{x} = ($, p_{x} are density matrices
then $S(\sum_{x} p_{x} p_{x}) \ge \sum_{x} p_{(x)} S(p_{x})$

Proof: since S(X:Q) above non-negative.

Strong subadditivity (SSA) and 4 more equivalent statements



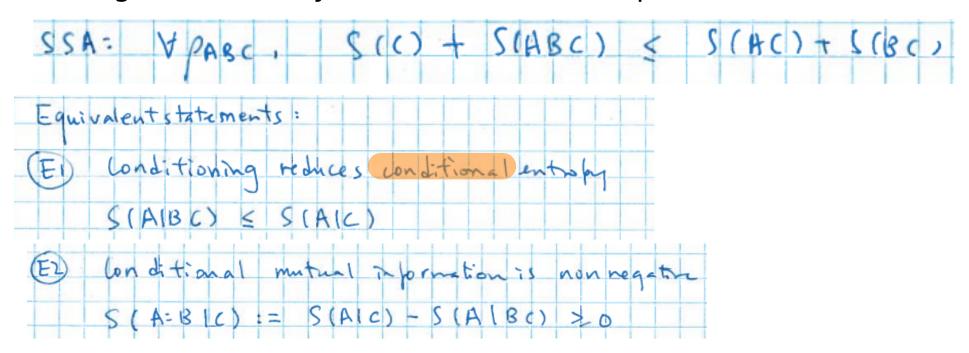
NB. System C is special.

NB. If dim C = 1, SSA reduces to SA:

$$0 + S(AB) \leq S(A) + S(B)$$

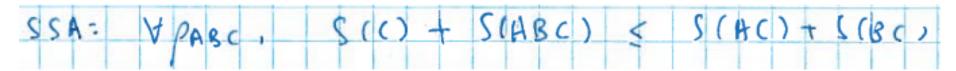
Proofs (several listed in the notes, possible term project)

Strong subadditivity (SSA) and 4 more equivalent statements

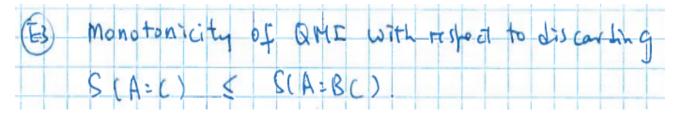


E1, E2 are basic rewritings of SSA.

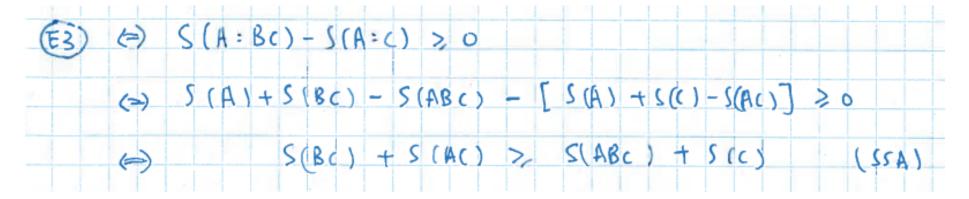
Strong subadditivity (SSA) and 4 more equivalent statements



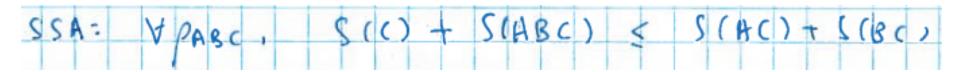
Further equivalent statements:



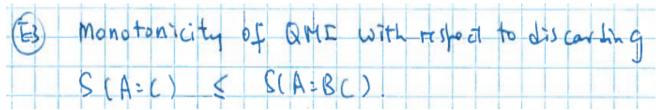
Proof (E3 <=> SSA):

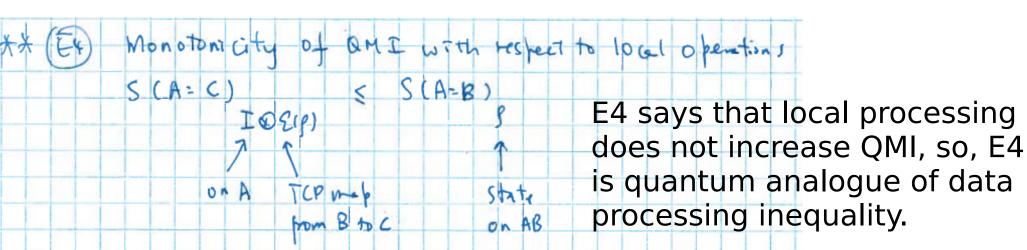


Strong subadditivity (SSA) and 4 more equivalent statements



Further equivalent statements:



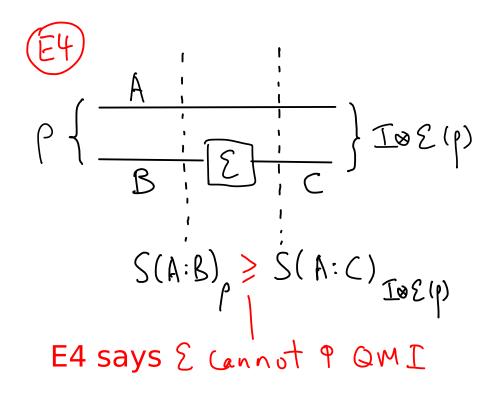


Proof (E3<=>E4): Assume E4 is true. Applying E4 with B=CD, $\{ cal E \} = partial trace of D gives E3.$ Proof (E3<=>E4): Assume E3 is true.

Stinespring dilation unitary representation

First examine what E4 says:

Now express \mathcal{E} as isometric extension



So, if E3 holds, E4 also holds.

How much do the entropy and QMI change if a system B is added or removed?

The Araki-Lieb inequality

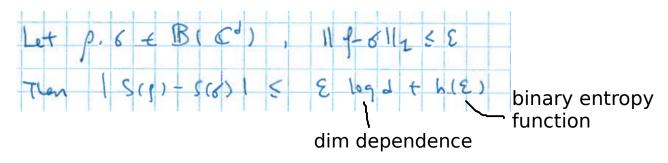
(a)
$$| S(AB) - S(A) | \leq S(B)$$

(b)
$$| S(A:BC) - S(A:C) | \leq 2 S(B)$$

Proof (see notes)

How much does the entropy change if the underlying state changes a little?

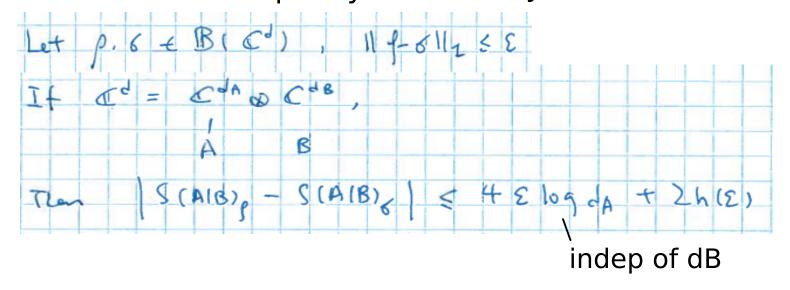
Fannes inequality (Continuity of S)



What about S(A:B) = S(A) + S(B) - S(AB)?

Can use Fannes inequality on each term ... suboptimal if dB large

Fannes-Alicki inequality (Continuity of conditional entropy)



Then, S(A:B) = S(A) - S(A|B) continuous (so is quantum capacity)!