CO781 / QIC 890:

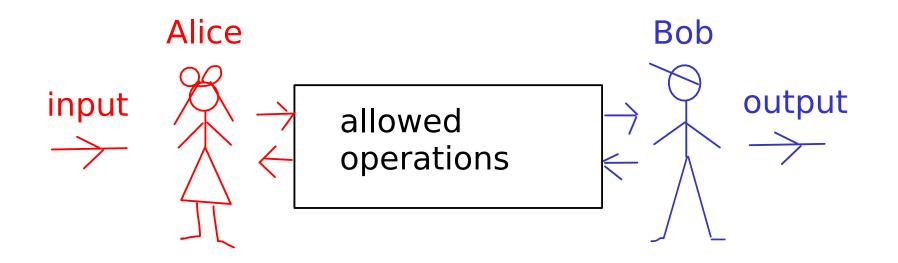
Theory of Quantum Communication

Topic 1, part 1

What is communication of data?

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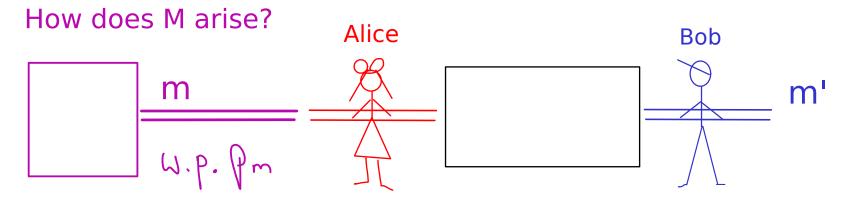
Simplest scenario: one <u>sender</u>, one <u>receiver</u>



Intuitively: if input data, output data are "similar" the data is communicated from Alice to Bob.

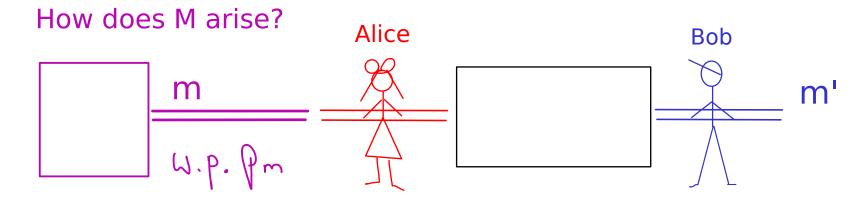
NB Local operations are "free".

e.g.1. Data: classical message $m \in \{0,1\}$ outcome of some random variable (rv) M



Rules: Alice and Bob do not know m before it arises, their initial state (shared or product) is independent of m. Only after Alice receives m may her operations depend on m. Bob's operations cannot depend on m, but they operate on "data coming from Alice that depends on m".

e.g.1. Data: classical message $m \in \{0,1\}$ outcome of some random variable (rv) M

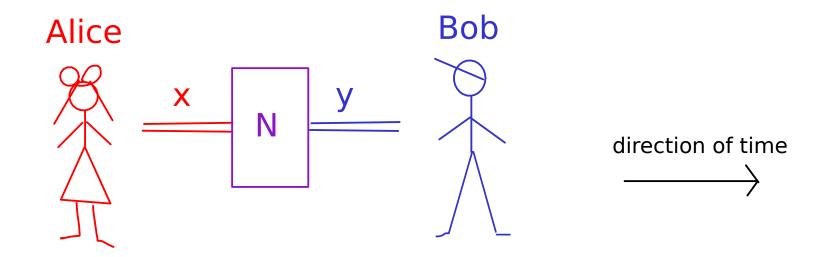


Intuitively: the data is communicated if for each possible m received by Alice, Bob's output m'=m.

Bob's output m' is a rv, call it M'.

If the data is communicated properly, necessary that: (i) M' approx M, (ii) if a rv T correlates with M, then, T correlates with M' the same way.

Def: A noisy classical channel N from Alice to Bob with input alphabet X & output alphabet Y is specified by Pr(Y=y | X=x) for all possible x, y.



Def: the noiseless classical channel over X has $X \subseteq Y$ and Pr(Y=y | X=x) = δ_{xy} i.e., Pr(Y=y | X=x) = $\begin{cases} 1 \text{ if } y=x, \\ 0 \text{ otherwise.} \end{cases}$

e.g.1. Data: classical message $m \in \{0, 1\}$

Concept: protocol. If Alice and Bob can use *once* a noiseless channel N of |X|=2, Alice can communicate M to Bob by choosing channel input x=m.

Concept: the channel N is a resource consumed in the above protocol. (Alice & Bob use <u>1 cbit.</u>)

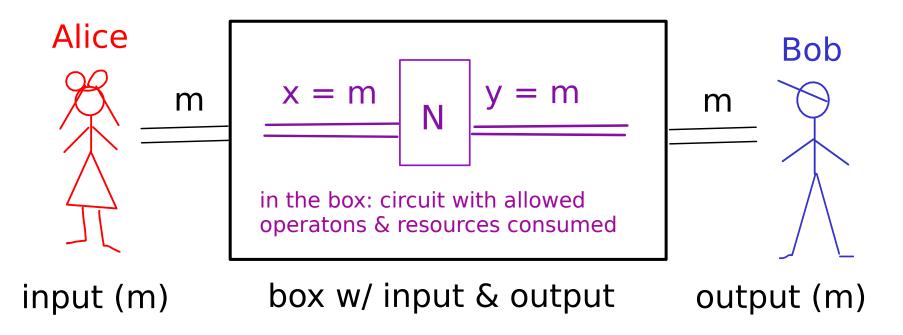
Concept: the ability to communicate M is a resource produced by the protocol (and can in turns be consumed as a subroutine elsewhere). (Alice & Bob create 1 cbit.)

 \notin the protocol simulates a noiseless channel (M->M')

Generalization: for |X| = d, log d cbits are consumed etc. Sometimes add direction of comm to cbit. e.g., 1 cbit \rightarrow

e.g.1. Data: classical message $m \in \{0, 1\}$

Diagramatic representation of a protocol:



Convention: the diagram holds for all m Formally : input is a mixture over m, apply linearity

e.g.1. Data: classical message $m \in \{0, 1\}$

$$\forall \mathcal{M} \quad |\mathsf{m}\rangle\langle\mathsf{m}|_{\mathcal{A}} \xrightarrow{\mathsf{enc}} |\mathsf{m}\rangle\langle\mathsf{m}|_{X} \xrightarrow{\mathsf{N}} |\mathsf{m}\rangle\langle\mathsf{m}|_{Y} \xrightarrow{\mathsf{dec}} |\mathsf{m}\rangle\langle\mathsf{m}|_{\mathcal{B}}$$

Question: is the following a reasonable description of communicating M to Bob?

$$\frac{1}{2} \Pr[m] \langle m|_{A} \xrightarrow{enc} \sum_{m=0}^{1} \Pr[m] \langle m|_{X}$$

$$\frac{N}{N} \xrightarrow{\sum_{m=0}^{1}} \Pr[m] \langle m|_{Y}$$

$$\frac{dec}{m=0} \xrightarrow{\sum_{m=0}^{1}} \Pr[m] \langle m|_{B}$$

It is necessary, but not sufficient.

e.g.1. Data: classical message $m \in \{0, 1\}$

$$\forall \mathcal{M} \quad |\mathsf{m}\rangle\langle\mathsf{m}|_{\mathcal{A}} \xrightarrow{\mathsf{enc}} |\mathsf{m}\rangle\langle\mathsf{m}|_{\chi} \xrightarrow{\mathsf{N}} |\mathsf{m}\rangle\langle\mathsf{m}|_{\chi} \xrightarrow{\mathsf{dec}} |\mathsf{m}\rangle\langle\mathsf{m}|_{\mathcal{B}}$$

Question: is the following a reasonable description of communicating M to Bob?

$$\frac{1}{2} \Pr |m\rangle \langle m|_{A} \xrightarrow{\text{enc}} \sum_{m=0}^{1} \Pr |m\rangle \langle m|_{X}$$

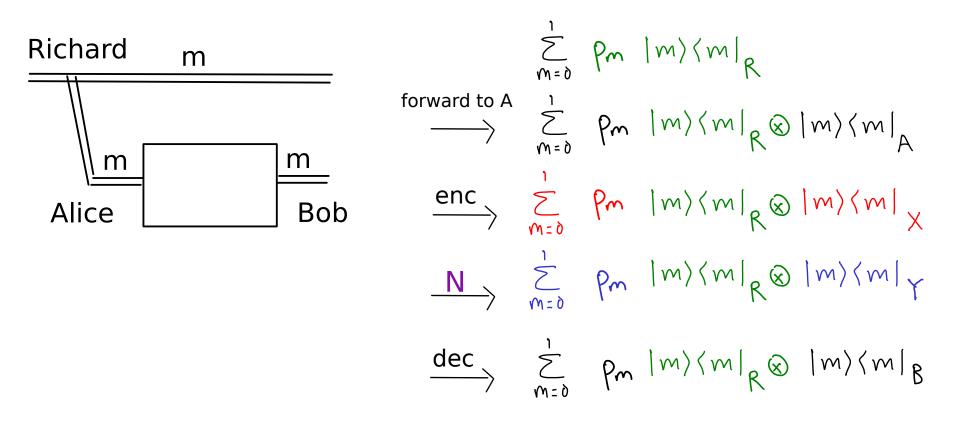
$$\frac{N}{N} \geq \sum_{m=0}^{1} \Pr |m\rangle \langle m|_{Y}$$

$$\frac{\text{dec}}{M=0} \sum_{m=0}^{1} \Pr |m\rangle \langle m|_{B}$$

It is necessary, but not sufficient.

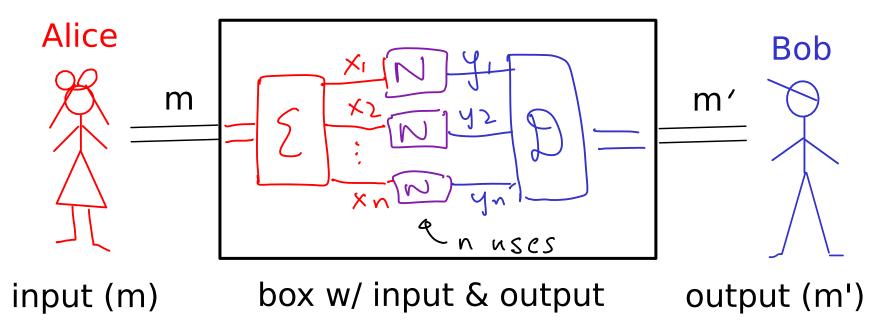
- 1. Bob could have created this without Alice.
- 2. Cannot impose correctness (e.g., p0 = p1 = 1/2, 0 < -> 1)

e.g.1. Data: classical message $m \in \{0,1\}$ Model: a reference Richard is holding a copy of m:



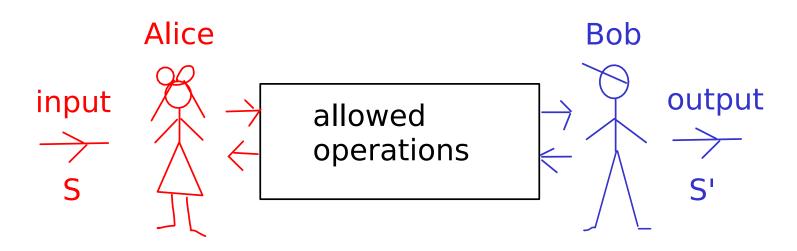
Ex: if T is correlated with M (as defined by a joint distribution pr(tm), how does T correlate with M'?Ex: what if Alice also keeps a copy of m?

- e.g.1++. Data: classical message m $\in \{1, 2, 3, ..., 2^n\}$
- A protocol to transmit data through a noisy channel:



Shannon's noisy coding theorem: optimize r.

e.g.2. Quantum data: state (variable) on a system S



One definition of communication of quantum data: for each input state ho on S, the output on S' is ho , (or approximately so)

(We will see weaker definitions later, e.g., quantum data compression or remote state preparation.)

Def: A quantum channel N from a d1-dim system X to a d2-dim system Y is a function from d1xd1 matrices to d2xd2 matrices that are (1) linear, (2) trace-preserving, (3) completely positive. (aka TCP maps, Q ops etc)

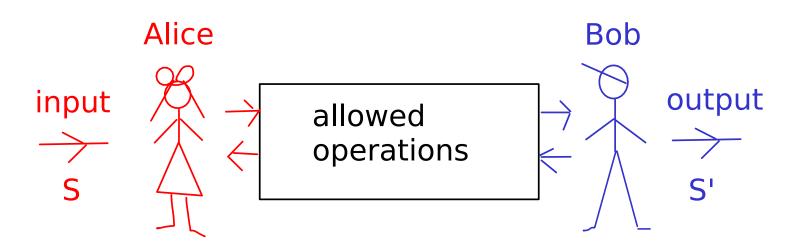
Def: A noiseless quantum channel N on a d-dim sys is given by the identity map on dxd matrices (can also embed in an output space with dim bigger than d).

Similar to e.g.1, can communicate a 2-dim Q system (creating 1 qbit) by using a noiseless quantum channel on a 2-dim system (consuming 1 qbit).

(Extensions: log d qbits, adding direction)

Recall: quantum channels and their representations

e.g.2. Quantum data: state (variable) on a system S



Alternative definition of comm of quantum data: each input in S is mapped to the output on S' (approximately) by the noiseless quantum channel.

e.g.2. Quantum data: state (variable) on a system S Def: the diamond norm distance between two channels N_1 , N_2 is given by

 $\| N_{1} - N_{2} \|_{\diamond} := \max \| I \otimes N_{1} (I + X + I) - I \otimes N_{2} (I + X + I) \| t$ $| / / I + Y_{RS} / I + Y_{$ from S to S' trace distance $\dim(R) = \dim(S)$

NB M1, M2 same dim, Hermitian. Schatten 1-norm || M1 - M2 || is the sum of the absolute values of the eigenvalues of M1-M2. $|| \dots || t = 1/2 || \dots || 1.$ If M1, M2 are density matrices on sys S, referee draws one at random, prepares the state on S and gives the system S to Alice (she doesn't know which state), then, max prob for Alice to say whether the state is M1 or M2 is $\frac{1}{3} + \frac{1}{9} \parallel M_{1} - M_{2} \parallel \frac{1}{4}$.

e.g.2. Quantum data: state (variable) on a system S Def: the diamond norm distance between two channels N_1 , N_2 is given by

$$\| N_{1} - N_{2} \|_{\diamond} := \max \| \mathbb{I} \otimes N_{1} (|\mathsf{T} \times \mathsf{T}|) - \mathbb{I} \otimes N_{2} (|\mathsf{T} \times \mathsf{T}|) \| \mathbf{t}$$

$$| \mathsf{T} \times \mathsf{RS}$$

$$from S to S' \qquad | dim(R) = dim(S) \qquad trace distance$$

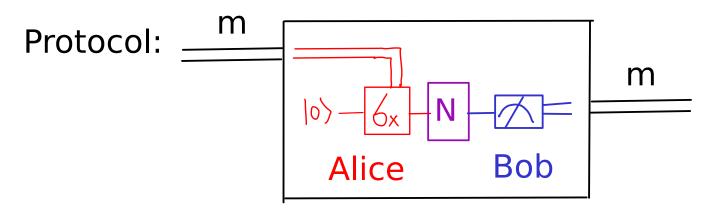
Consequence: If a communication protocol approximates the noiseless channel from S to S' to diamond norm distance \leqslant \pounds , then,

 $\forall \, | \Psi \rangle$ on system RS, the protocol yields a state on RS'

with trace distance less than \in from $|\Psi\rangle$.

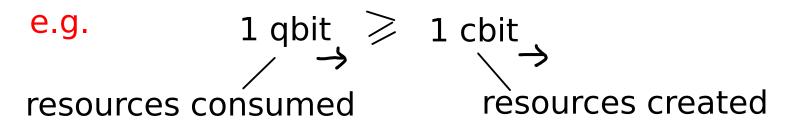
i.e., this notion of approx preserves correlation with S.

e.g. 3. Alice can communicate 1 classical bit to Bob using a noiseless quantum channel N on a 2-dim sys.

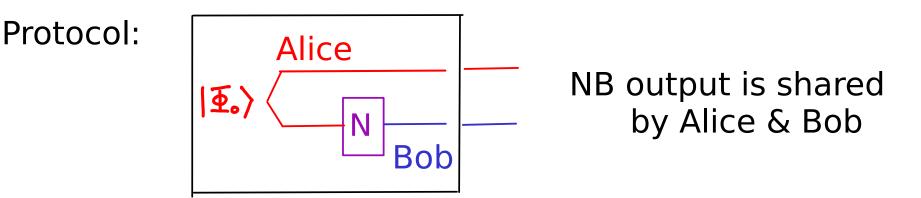


This consumes "1 qbit" (the ability to send a 2-dim sys) and creates "1 cbit".

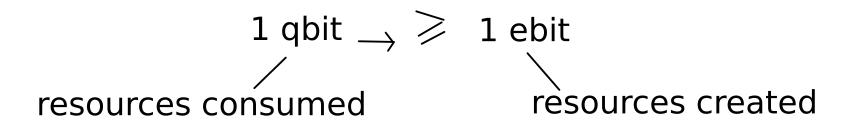
Concept: a resource inequality holds if there exists a protocol converting the resources (later: need approx).



e.g. 4. Alice and Bob can share $|\Phi_{\circ}\rangle = \frac{1}{\sqrt{2}}(|\circ\circ\rangle + |11\rangle)$ using a noiseless quantum channel N on a 2-dim sys.



This consumes "1 qbit" and creates "1 ebit" (a maximally entangled state of local dim 2). Resource inequality:



NB Above requires qbit to be entanglement preserving.

Exercise:

What is meant by the resource inequality

1 qbit $\leftarrow \geq 1$ ebit ?

Is it true? If so, provide a protocol. If not, provide a proof that no such protocol exists.

e.g 5 Superdense coding (SD) (Bennett-Wiesner "92")

Suppose Alice and Bob share the state $\frac{1}{25} \stackrel{\circ}{\subseteq} 1381$ and Alice can send an s-dimensional quantum system to Bob. Then, Alice can communicate any of t=s² messages to Bob.

Data: classical Allowed resources: noiseless q channel, entanglement Proof: for simplicity, first consider s=2. Suppose Alice & Bob share the state $|\Phi_{\circ}\rangle = \frac{1}{\sqrt{2}}(|\circ\circ\rangle + |\cdots\rangle)$ so that Alice (Bob) holds the first (second) qubit.

Recall the Pauli matrices:

$$\mathcal{E}_{\mathcal{O}} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \quad \mathcal{E}_{\mathcal{X}} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \mathcal{E}_{\mathcal{Y}} = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \mathcal{E}_{\mathcal{Z}} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

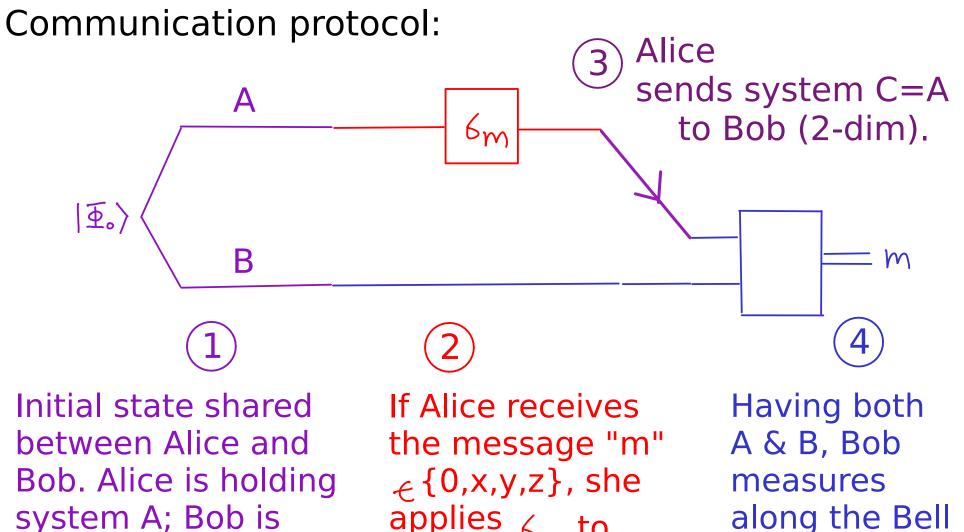
Suppose Alice wants to communicate a message m from the set $\{o, \chi, y, z\}$.

If her message is m, she applies 6_m to A. The shared state $|\Phi_0\rangle$ on AB is transformed by $6_m \otimes I$.

For $ \Phi_{\circ}\rangle = \frac{1}{\sqrt{2}} (00\rangle + 11\rangle)$
$\delta_{o} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}, \delta_{\chi} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \delta_{y} = \begin{pmatrix} 0 & -\overline{\iota} \\ \overline{\iota} & 0 \end{pmatrix}, \delta_{\Xi} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$
$ \Phi_{\circ}\rangle = 6_{\circ} \otimes \mathbb{I} \Phi_{\circ}\rangle = \frac{1}{\sqrt{2}} (00\rangle + 11\rangle)$
$ \Phi_{\mathbf{x}}\rangle = 6_{\mathbf{x}} \otimes \mathbb{I} \Phi_{\mathbf{x}}\rangle = \frac{1}{\sqrt{2}} (10\rangle + 01\rangle)$
$ \Phi_y\rangle = 6_y \otimes I \Phi_o\rangle = \frac{1}{\sqrt{2}}(i io\rangle - i oi\rangle)$
$ \Phi_{z}\rangle = 6_{z} \otimes \mathbb{I} \Phi_{o}\rangle = \frac{1}{\sqrt{2}} (00\rangle - 11\rangle)$

These 4 states are mutually orthogonal, and form the "Bell basis". NB Alice operates on a 2-dim system A, but the joint state on AB tranverses to 1 out of 4 possible distinguishable (ortho) states.

If Alice sends C=A to Bob, he has AB in the state $|\Phi_v\rangle_{.}$ He can measure AB along the Bell basis to find v !

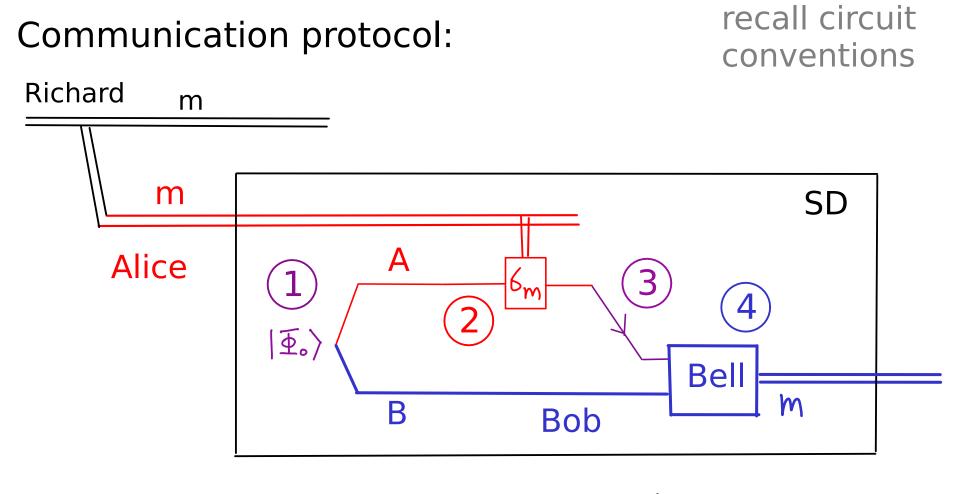


holding system B.

applies 6m to qubit A.

along the Bell basis; outcome is m.

1 ebit

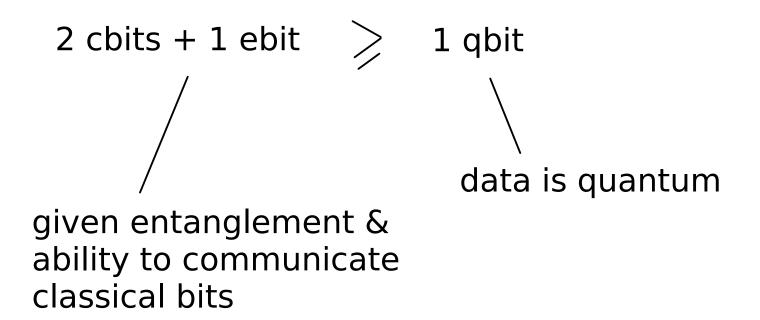


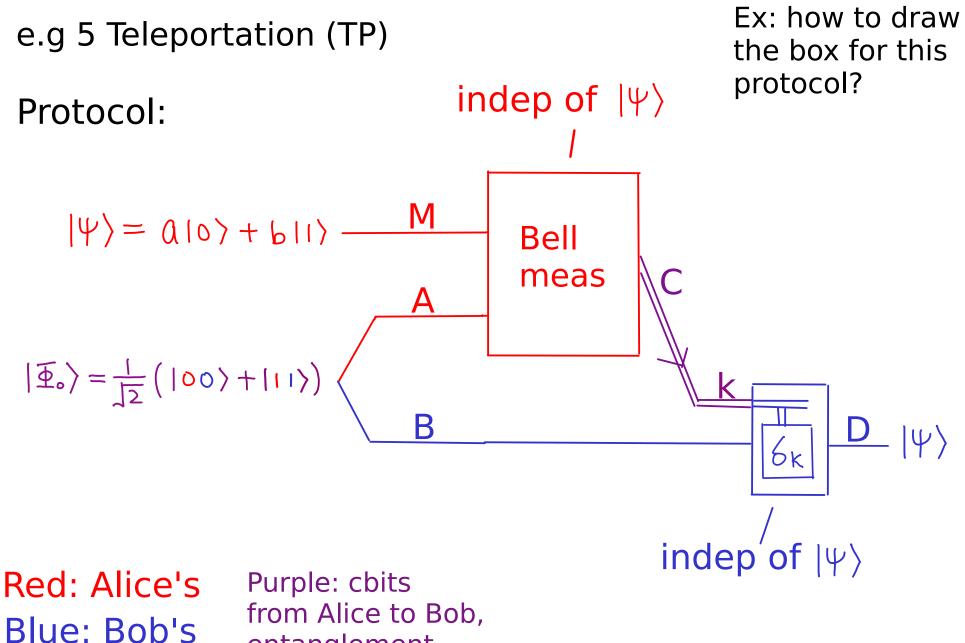
Resource inequality: 1 ebit + 1 qbit > 2 cbits \rightarrow \rightarrow

NB the ebit can be pre-shared, or created by Bob through back communication. "Superdense" refers to the communication of 1 in s^2 distinguishable messages by sending an s-dim sys.

e.g 5 Teleportation (TP) (Bennett-Brassard-Crepeau-Jozsa-Peres-Wootters 93)

In the language of resource inequality:





entanglement

Proof: main mathematical tool: Expressing an 8-dim quantum state in 2 ways.

$$(a | o \rangle + b | 1 \rangle)_{M} \frac{1}{\sqrt{2}} (| o o \rangle + | 1 | \rangle)_{AB}$$

$$= (a | o o o \rangle + a | o 1 | \rangle + b | 1 o o \rangle + b | 1 | 1 \rangle)_{MAB} \frac{1}{\sqrt{2}}$$

$$= \frac{1}{\sqrt{2}} (| o o \rangle + | 1 | \rangle)_{MA} (a | o \rangle + b | 1 \rangle)_{B} \frac{1}{2}$$

$$= \frac{1}{\sqrt{2}} (| o o \rangle - | 1 | \rangle)_{MA} (a | o \rangle - b | 1 \rangle)_{B} \frac{1}{2}$$

$$= \frac{1}{\sqrt{2}} (| o o \rangle - | 1 | \rangle)_{MA} (a | o \rangle - b | 1 \rangle)_{B} \frac{1}{2}$$

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$$= \frac{1}{\sqrt{2}} (| o o \rangle - | 1 | \rangle)_{MA} (a | o \rangle - b | 1 \rangle)_{B} \frac{1}{2}$$

$$+ \frac{1}{\sqrt{2}} (|0|\rangle - ||0\rangle)_{MA} (\alpha ||\rangle - b |0\rangle)_{B} \frac{1}{2} = \alpha |0|| + b |00\rangle$$

$$\begin{split} & (\Psi) \\ & (A|0\rangle + b|1\rangle)_{M} \stackrel{1}{\sqrt{2}} (100\rangle + 111\rangle)_{AB} = \\ & (\Xi_{\bullet}) \rightarrow \frac{1}{\sqrt{2}} (100\rangle + 111\rangle)_{MA} (A|0\rangle + b|1\rangle)_{B} \stackrel{1}{\sqrt{2}} (1\Psi) \\ & (\Xi_{\bullet}) \rightarrow \frac{1}{\sqrt{2}} (100\rangle - 111\rangle)_{MA} (A|0\rangle - b|1\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & (\Xi_{\bullet}) \rightarrow \frac{1}{\sqrt{2}} (101\rangle + 110\rangle)_{MA} (A|1\rangle + b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle + 110\rangle)_{MA} (A|1\rangle + b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (101\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (100\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (100\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} \stackrel{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} (100\rangle - 110\rangle)_{MA} (A|1\rangle - b|0\rangle)_{B} (\Phi) \\ & + \frac{1}{\sqrt{2}} (E_{\bullet}) (\Psi) \\ & + \frac{1}{\sqrt{2}} ($$

If Alice measures MA along the Bell basis, each outcome $k \leftarrow \{0, x, y, z\}$ occurs with prob 1/4, and postmeasurement state is $|\Phi_k\rangle_{MA} \otimes G_k |\Psi\rangle_B$.

If Alice sends k to Bob, he can apply \mathcal{G}_{k} to B, turning $\mathcal{G}_{k} | \Psi \rangle_{B}$ to $| \Psi \rangle_{B}$.

Remarks:

- 0. What is teleported, the body or the soul ?
- 1. $|\Psi\rangle$ takes infinitely many bits to describe. Also, given a copy, Alice cannot learn a description of $|\Psi\rangle$
- 2. But Alice's operations are independent of $|\psi\rangle$ and only 2 bits of info need to be sent.
- 3. Generalizes to higher dimension.
- 4. Preserves global state / correlations
- Ex: for any state $|\Psi\rangle_{RM}$ where M is 2-dim, show that

147RM = Q (do) 10) + b (di) (1)

for some unit vectors $|\langle q_0 \rangle$, $|\langle q_1 \rangle$ on R, and $|\langle q_1 \rangle^2 + |\langle b_1 \rangle^2 - |$. If Alice teleports M to Bob (with identity map on R), repeat the above analysis to show that the final state on RB is $|\langle q \rangle$.

Summary: in each communication problem, need to clarify the data type, how it arises, what the output state should be (in terms of the input state), what resources are available to sender and receiver etc.

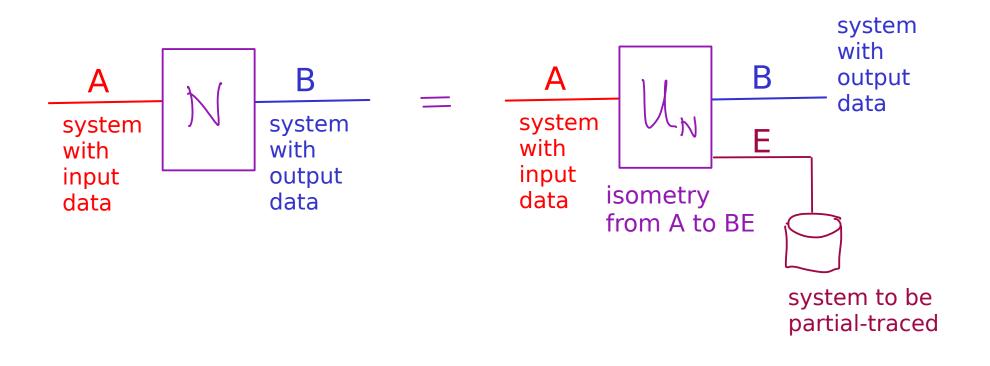
We saw examples how resources are used, and how various notions of "correct communication" are achieved.

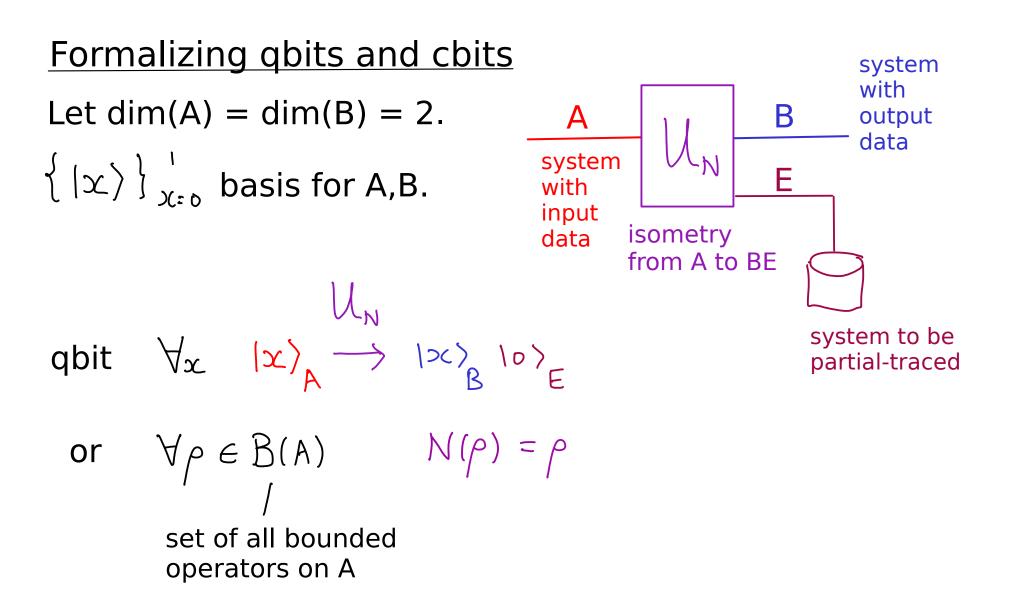
We have considered resources cbits, qbits, ebits that satisfy all properties we want, and behave as we want as a component in any protocol.

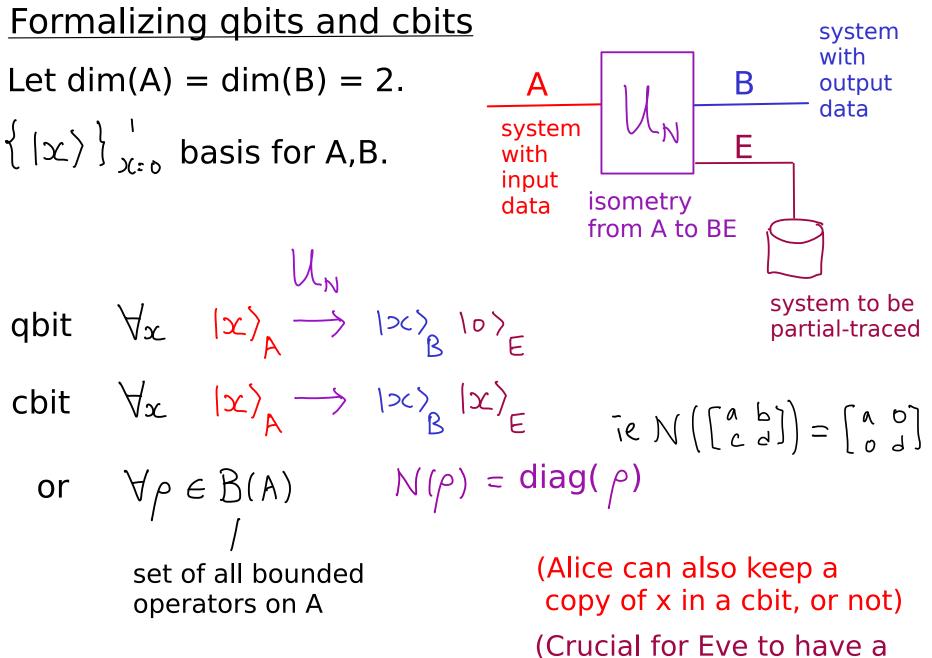
Formalizing qbits and cbits

One can model the transformation from the input system to the output system as a quantum channel.

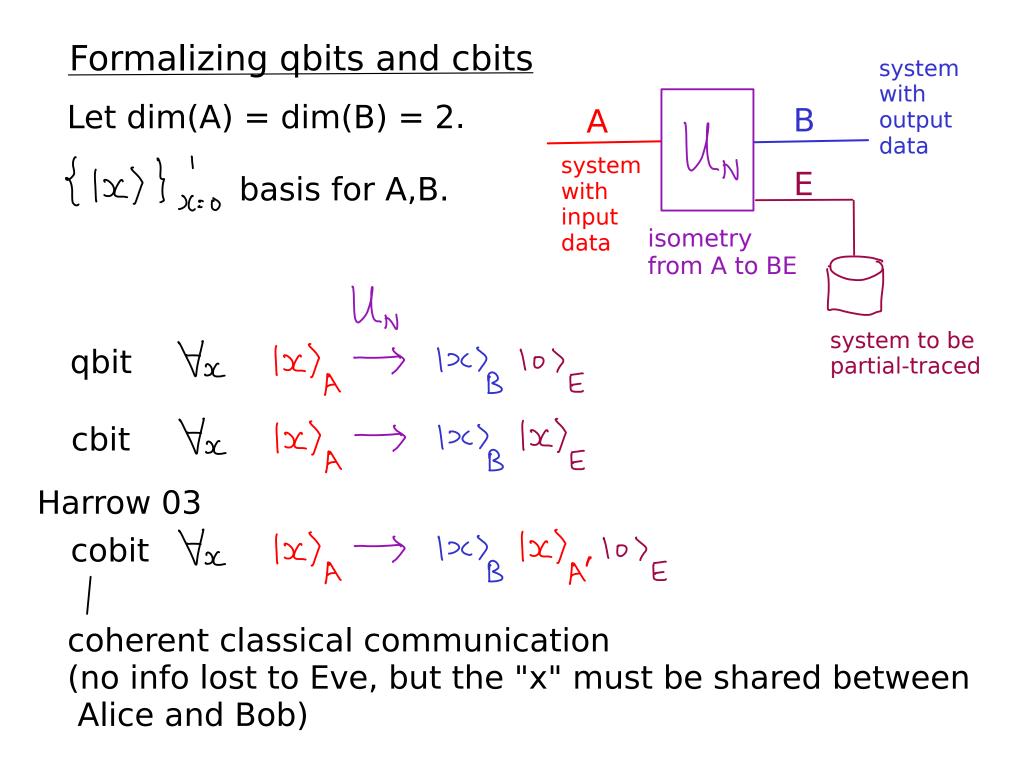
Recall that any quantum channel has a Stinespring dilation (also called the isometric extension or unitary representation):

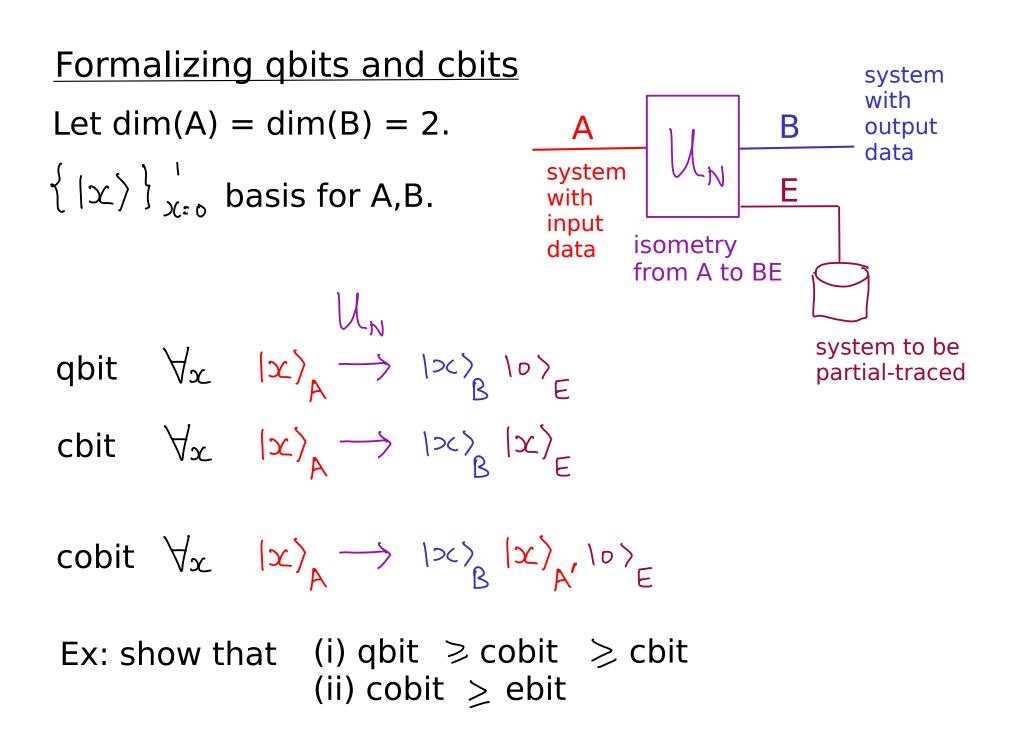






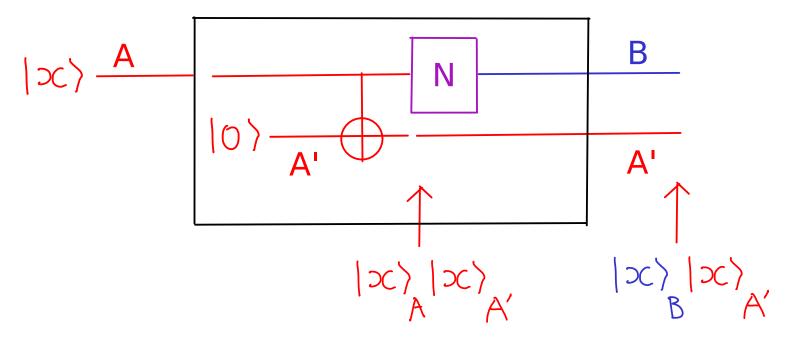
copy that no one can take back -- def of being classical)





e.g., How to convert 1 qbit into 1 cobit? In particular, how to get an additional output sys A' ?

Ans: local operations are allowed. Before applying the channel, Alice "makes a copy" of A onto A' "coherently"



and if input at A is in superposition the state on AA' is in superposition and final state on BA' in superposition

<u>Composable communication:</u>

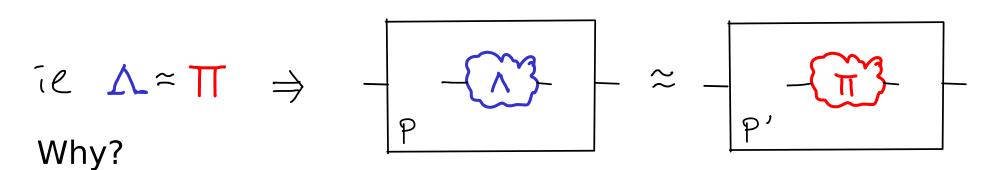
Let Λ be a communication resource (e.g., qbit or cbit).

Let \prod be a protocol approximately creating the resource (same input and output space).

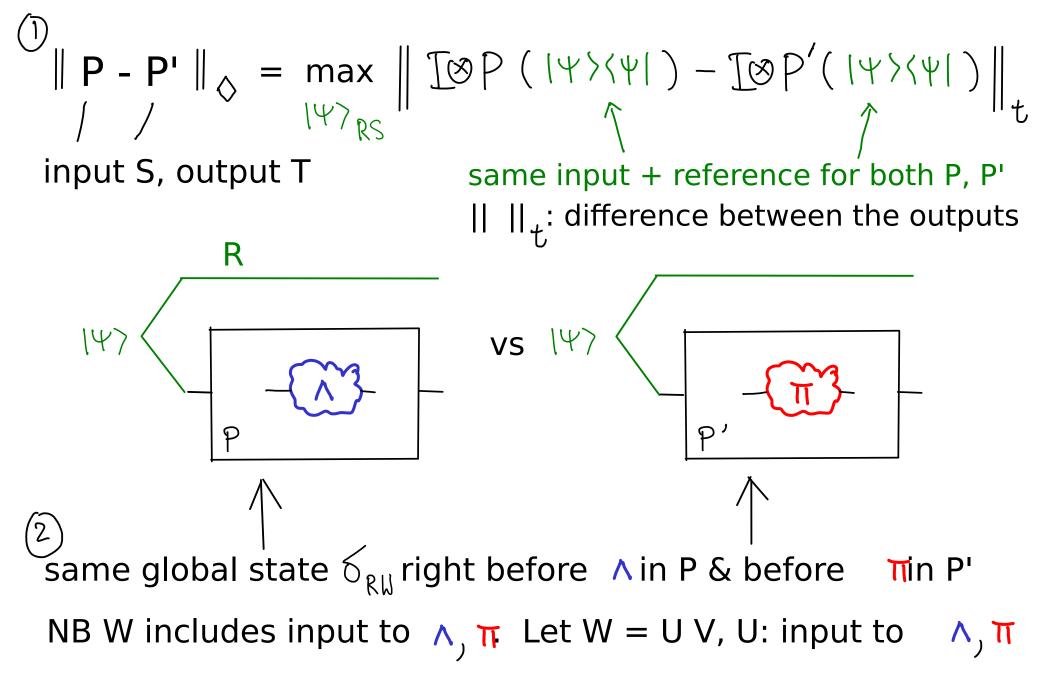
If $\| \bigtriangleup - \Pi \|_{\Diamond} \leq \varepsilon$

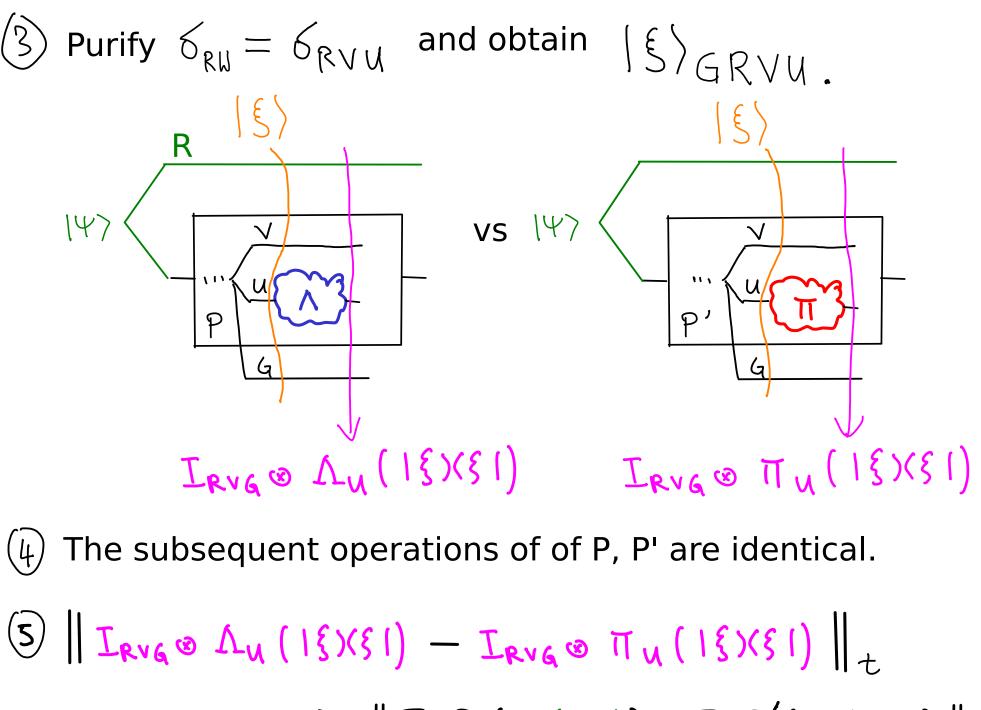
then, we have the following desirable composable property:

Let P be any protocol consuming \bigwedge as a resource. Replace \bigwedge by \prod in P, and call the resulting protocol P'. Then, $\| P - P' \|_{\diamondsuit} \leq 2$. NB, \land , Π , P, P' all TCP maps



Proof sketch for composability of resources with approximation in diamond norm:





 $\geq \left\| I \otimes P\left(|\psi\rangle\langle\psi| \right) - I \otimes P'(|\psi\rangle\langle\psi|) \right\|_{t}$

why? because same op $I_{RVG} \otimes \Lambda_{U}(|\{\rangle\langle \xi|) \longrightarrow I \otimes P(|\Psi\rangle\langle \Psi|)$ $I_{RVG} \otimes \Pi_{U}(|\{\rangle\langle \xi|) \longrightarrow I \otimes P'(|\Psi\rangle\langle \Psi|)$

and trace distance is monotonic decreasing if the same operation is applied to the two states to be compared (cf Watrous lectures, or obsorb the operation as the distinguishing measurement of the pre-operation state)

(i) Since
$$\| \wedge - \Pi \|_{\mathcal{O}} \leq \varepsilon$$
,
 $\| I_{RVG} \otimes \Lambda_{U} (|\xi)(\xi|) - I_{RVG} \otimes \Pi_{U} (|\xi)(\xi|) \|_{t} \leq \varepsilon$

) Combining the conclusions of (6) & (7) proves the claim.

Final remarks: composable resources are strong resources. They are harder to create, but are very useful if given. They're simpler to consider.

We will see both composable resources, and weaker form of communication in this course (the latter are easlier to obtain, but are less useful) so these tends to be "end goals" of our quantum information processing tasks. They are often more specific and awkward to define ...