

Welcome to

Quantum Morphisms

# Some math basics (+ some notation)

bra-ket notation

$|\psi\rangle$  - a column vector (a "ket")

$\langle\psi|$  - a row vector, the conjugate transpose of  $|\psi\rangle$   
(a "bra")

$|i\rangle \in \mathbb{C}^n$  -  $i^{\text{th}}$  standard basis vector (usually  $i \in \{1, \dots, n\}$  or  $\{0, \dots, n-1\}$ )

$|x\rangle \in \mathbb{C}^Y$  - standard basis vector indexed by  $x \in Y$

$\langle\psi|\varphi\rangle$  - inner product of  $|\psi\rangle$  &  $|\varphi\rangle$  (a scalar)

$|\psi\rangle\langle\varphi|$  - outer product of  $|\psi\rangle$  &  $|\varphi\rangle$  (a matrix)

$\mathbb{F}^{m \times n}$  -  $m \times n$  matrices over  $\mathbb{F}$

Important types of matrices (in  $\mathbb{C}^{n \times n}$ )

• Hermitian:  $M^* = M$  i.e.  $M_{ij} = \overline{M_{ji}}$

• Positive semidefinite (psd): Hermitian &  $\langle\psi|M|\psi\rangle \geq 0$   
 $M \succeq 0$   $\forall |\psi\rangle \in \mathbb{C}^n$

• Unitary:  $M^*M = MM^* = I$

These are all normal i.e.  $M^*M = MM^*$

## Spectral Decomposition

$$\langle \psi_i | \psi_j \rangle = \delta_{ij} := \begin{cases} 1 & \text{if } i=j \\ 0 & \text{if } i \neq j \end{cases}$$

There exists an orthonormal basis  $|\psi_1\rangle, \dots, |\psi_n\rangle \in \mathbb{C}^n$  s.t.

$$M = \sum_{i=1}^n \lambda_i |\psi_i\rangle \langle \psi_i| \quad \lambda_i \in \mathbb{C} \text{ are the eigenvalues of } M$$

Often write  $M = \sum_{\lambda} \lambda E_{\lambda}$  where  $E_{\lambda} = \sum_{i: \lambda_i = \lambda} |\psi_i\rangle \langle \psi_i|$

Then  $\underbrace{E_{\lambda}^2 = E_{\lambda}^* = E_{\lambda}}_{\text{projection}}, \quad E_{\lambda} E_{\mu} = \delta_{\lambda\mu} E_{\lambda} \quad \& \quad \sum_{\lambda} E_{\lambda} = I$

$M$  Hermitian  $\Rightarrow$  real eigenvalues

$$\sum_{\lambda} \sqrt{\lambda} E_{\lambda}$$

$M$  psd  $\Rightarrow$  non-negative eigenvalues ( $M^{1/2}$  well-defined)

$M$  unitary  $\Rightarrow$  eigenvalues of unit modulus,  $\lambda = e^{i\theta} \quad 0 \leq \theta < 2\pi$

## Quantum States

A quantum state is a description of a quantum system

"state space"

A pure quantum state is a unit vector  $|\psi\rangle \in \mathbb{C}^n$

Some PHYSICS: an electron can have spin up  $|\uparrow\rangle$  or down  $|\downarrow\rangle$   
or any superposition of the two:  $\alpha|\uparrow\rangle + \beta|\downarrow\rangle \quad |\alpha|^2 + |\beta|^2 = 1$

What (linear) operations map states to states?

$M: \mathbb{C}^n \rightarrow \mathbb{C}^n$  must satisfy  $\langle \psi | M^* M | \psi \rangle = 1 \quad \forall |\psi\rangle \in \mathbb{C}^n$

Exercise: Show this implies  $M^* M = I$

& thus  $M$  is unitary.

## General quantum states

Suppose we know that a given system is in state  $|0\rangle \in \mathbb{C}^2$  w/ probability  $1/2$  + in state  $|1\rangle \in \mathbb{C}^2$  w/ prob  $1/2$ .

How could we describe this?

$$\frac{1}{2}|0\rangle + \frac{1}{2}|1\rangle = \frac{1}{2}(|0\rangle + |1\rangle) = \begin{pmatrix} 1/2 \\ 1/2 \end{pmatrix}$$

$\begin{pmatrix} 1/2 \\ 1/2 \end{pmatrix}$  not a unit vector!

Instead we use **density matrices**

$$\rho \in \mathbb{C}^{n \times n} \text{ s.t. } \rho \geq 0 + \text{Tr}(\rho) = 1$$

pure state

$$|\psi\rangle \in \mathbb{C}^n \\ \propto |\psi\rangle$$

density matrix

$$|\psi\rangle\langle\psi| \in \mathbb{C}^{n \times n} \\ \propto |\psi\rangle\langle\psi|$$

mixed state/ensemble

$$|\psi_i\rangle \text{ w/ prob } p_i \longrightarrow \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

density matrix

$$p_i \text{ w/ prob } p_i \longrightarrow \sum_i p_i \rho_i$$

$$\text{Above example: } \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1| = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Exercise: Show that if  $|\psi_0\rangle + |\psi_1\rangle$  are an orthonormal basis for  $\mathbb{C}^2$ , then the density matrix for the ensemble

$|\psi_i\rangle$  w/ prob  $1/2$  for  $i=0,1$  is  $\frac{1}{2}I$ . Do for  $\mathbb{C}^n$ .

# Measurements

A quantum measurement is a tuple  $\mathcal{M} = (M_1, \dots, M_k)$  with  $M_i \in \mathbb{C}^{m \times n}$  s.t.  $\sum_{i=1}^k M_i^* M_i = I$ .

When a state  $\rho \in \mathbb{C}^{n \times n}$  is measured with  $\mathcal{M}$  it results in

① a classical outcome  $i \in [k]$  w/ prob

$$\text{Tr}(M_i \rho M_i^*) = \underbrace{\text{Tr}(\rho M_i^* M_i)}_{\geq 0}$$

$$\sum_i \text{Tr}(\rho M_i^* M_i) = \text{Tr}(\rho) = 1$$

pure case:

$$\text{Tr}(M_i |\psi\rangle\langle\psi| M_i^*) = \langle\psi| M_i^* M_i |\psi\rangle = \|M_i |\psi\rangle\|^2$$

② and the corresponding post-measurement state

$$\frac{1}{\text{Tr}(M_i \rho M_i^*)} M_i \rho M_i^* = M_i \rho^{1/2} (M_i \rho^{1/2})^*$$

$$\rho = \frac{1}{n} I \quad M_i = |i\rangle\langle i| \quad \frac{1}{n} \text{Tr}(|i\rangle\langle i| |i\rangle\langle i|) = \frac{1}{n} \text{Tr}(|i\rangle\langle i|) = \frac{1}{n}$$

pure case:  $\frac{1}{\|M_i |\psi\rangle\|} M_i |\psi\rangle$

Note: The outcome probabilities only depend on the operators  $M_i^* M_i$ .

The  $M_i^* M_i$  form a positive operator valued measure (POVM), i.e. a tuple  $(P_1, \dots, P_k)$  with  $P_i \in \mathbb{C}^{n \times n}$

$$\text{s.t. } P_i \geq 0 \quad \forall i \quad + \quad \sum_i P_i = I.$$

$$(P_1^{1/2}, \dots, P_k^{1/2})$$

Example: Full basis measurement of a pure state

$$|\psi\rangle \in \mathbb{C}^n, \quad |\varphi_1\rangle, \dots, |\varphi_n\rangle \in \mathbb{C}^n \quad \text{ONB} \quad M_i = |\varphi_i\rangle\langle\varphi_i|$$

[Note  $M_i^* M_i = |\varphi_i\rangle\langle\varphi_i| |\varphi_i\rangle\langle\varphi_i| = |\varphi_i\rangle\langle\varphi_i| = M_i = M_i^*$   
special case of projective measurement:  $M_i^* = M_i = M_i^2$ .

Obtain outcome  $i$  w/ prob

$$\langle\psi| M_i^* M_i |\psi\rangle = \langle\psi| \varphi_i\rangle\langle\varphi_i| \psi\rangle = |\langle\varphi_i|\psi\rangle|^2$$

Post-measurement state:

$$\begin{aligned} \frac{1}{\|M_i|\psi\rangle\|} M_i|\psi\rangle &= \frac{1}{\| |\varphi_i\rangle\langle\varphi_i|\psi\rangle \|} |\varphi_i\rangle\langle\varphi_i|\psi\rangle \\ &= |\varphi_i\rangle \quad \text{up to a phase.} \end{aligned}$$

# Composite Systems

System A:  $|\psi_A\rangle \in \mathbb{C}^{d_A}$  ( $\rho_A \in \mathbb{C}^{d_A \times d_A}$ )

System B:  $|\psi_B\rangle \in \mathbb{C}^{d_B}$  ( $\rho_B \in \mathbb{C}^{d_B \times d_B}$ )

Combined system:

$$|\psi_A\rangle \otimes |\psi_B\rangle \in \mathbb{C}^{d_A} \otimes \mathbb{C}^{d_B} \cong \mathbb{C}^{d_A d_B}$$

$$\rho_A \otimes \rho_B \in \mathbb{C}^{d_A \times d_A} \otimes \mathbb{C}^{d_B \times d_B} \cong \mathbb{C}^{d_A d_B \times d_A d_B}$$

$$A \otimes B = \begin{pmatrix} A_{11}B & A_{12}B & \dots \\ A_{21}B & \dots & \dots \\ \vdots & \dots & A_{nn}B \end{pmatrix}$$

$$(A \otimes B)(C \otimes D) = AC \otimes BD$$

whenever  $AC + BD$  exist

A pure state  $|\psi\rangle \in \mathbb{C}^{d_A} \otimes \mathbb{C}^{d_B}$  is entangled if it cannot be written as

$$|\psi\rangle = |\psi_A\rangle \otimes |\psi_B\rangle$$

for states  $|\psi_A\rangle \in \mathbb{C}^{d_A}$ ,  $|\psi_B\rangle \in \mathbb{C}^{d_B}$ .

A state  $\rho \in \mathbb{C}^{d_A \times d_A} \otimes \mathbb{C}^{d_B \times d_B}$  is entangled if it cannot be written as

$$\rho = \sum_{i=1}^k p_i \rho_A^i \otimes \rho_B^i$$

for states  $\rho_A^i \in \mathbb{C}^{d_A \times d_A}$ ,  $\rho_B^i \in \mathbb{C}^{d_B \times d_B}$  & prob dist  $p_1, \dots, p_k$ .

Exercise: Show that  $\frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \in \mathbb{C}^2 \otimes \mathbb{C}^2$  is entangled.  
 $= |0\rangle \otimes |0\rangle$       $|\psi\rangle = |\psi\rangle \otimes |\varphi\rangle$   
 $|\psi\rangle \otimes |\varphi\rangle$

## Joint Measurements

$$\rho \in \mathbb{C}^{d_A \times d_A} \otimes \mathbb{C}^{d_B \times d_B}$$

Alice "has"  $\mathbb{C}^{d_A \times d_A}$      Bob "has"  $\mathbb{C}^{d_B \times d_B}$

Alice performs measurement  $A = (E_1, \dots, E_k)$  on her system

Equivalent to "global" measurement  $(E_i \otimes I, \dots, E_k \otimes I)$

Outcome  $i$  w/ prob  $\text{Tr}((E_i \otimes I) \rho (E_i^* \otimes I))$

post-measurement state  $(E_i \otimes I) \rho (E_i^* \otimes I)$  normalized

If Bob also measures with  $B = (F_1, \dots, F_r)$ :

outcome  $i, j$  w/ prob

$$\text{Tr}((E_i \otimes F_j) \rho (E_i^* \otimes F_j^*))$$

post meas state

$$(E_i \otimes F_j) \rho (E_i^* \otimes F_j^*) \text{ normalized}$$

Exercise: Show that doing both measurements "at the same time" is equivalent to doing one & then the other.



Example:  $|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \in \mathbb{C}^2 \otimes \mathbb{C}^2$

Measure first qubit in standard basis.

Outcome 0:  $\frac{1}{2} \text{Tr}(|0\rangle\langle 0| \otimes I) (|00\rangle\langle 00| + |00\rangle\langle 11| + |11\rangle\langle 00| + |11\rangle\langle 11|)$   
 $= \frac{1}{2} \text{Tr}(|0\rangle\langle 0|) = \frac{1}{2}$

Outcome 1:  $\frac{1}{2}$

post measurement state:  $M_0: |\psi\rangle = |0\rangle\langle 0| \otimes I \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$   
 $= |00\rangle$

" :  $|11\rangle$

Alice measures her qubit w/ full basis meas  $|\varphi_0\rangle, |\varphi_1\rangle$

Bob " his " " " " " "  $|\overline{\varphi_0}\rangle, |\overline{\varphi_1}\rangle$

