

# Dyson-Schwinger equations and Renormalization Hopf algebras

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## Unfolding some recursive equations

Lets get our intuition going

$$X = \mathbb{I} + xB_+(X^2)$$

What does this count?

$$X = \mathbb{I} + xB_+(X^3)$$

What does this count?

$$X = \mathbb{I} - xB_+\left(\frac{1}{X}\right)$$

What does this count?

1

## Answers

$$X = \mathbb{I} + xB_+(X^2)$$

counts computer science binary trees (separate slots for left and right children).

$$X = \mathbb{I} + xB_+(X^3)$$

counts ternary trees with separate slots for left, middle, and right children.

$$X = \mathbb{I} - xB_+\left(\frac{1}{X}\right)$$

counts plane rooted trees.

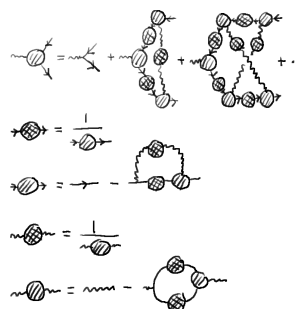
## Dyson-Schwinger equations combinatorially

As the simple tree examples, or systems

$$X^r(x) = \mathbb{I} - \sum_{k \geq 1} x^k p_r(k) B_+^{k,r}(X^r(x) Q(x)^k)$$

where  $Q(x) = \prod X^r(x)^{s_r}$  and  $r$  runs over the different external leg structures.

Example: QED



## Dyson-Schwinger equations analytically

Example from Broadhurst and Kreimer [3].

$$X(x) = \mathbb{I} - xB_+ \left( \frac{1}{X(x)} \right).$$

along with

$$F(\rho) = \frac{1}{q^2} \int d^4k \frac{k \cdot q}{(k^2)^{1+\rho}(k+q)^2} - \dots \Big|_{q^2=\mu^2}$$

gives  $(X \mapsto G, B_+ \mapsto F)$

$$G(x, L) = 1 - \frac{x}{q^2} \int d^4k \frac{k \cdot q}{k^2 G(x, \log k^2)(k+q)^2} - \dots \Big|_{q^2=\mu^2}$$

where  $L = \log(q^2/\mu^2)$ . The (analytic) Dyson-Schwinger equation for a bit of massless Yukawa theory.

4

## Dyson-Schwinger equations physically

Equations of motion, analogous to the classical differential equations of motion.

By expanding in the coupling constant Dyson-Schwinger equations give perturbation theory.

But Dyson-Schwinger equations also contain non-perturbative information if we can extract it. Broadhurst and Kreimer [3] solved

$$G(x, L) = 1 - \frac{x}{q^2} \int d^4k \frac{k \cdot q}{k^2 G(x, \log k^2)(k+q)^2} - \dots \Big|_{q^2=\mu^2}$$

where  $L = \log(q^2/\mu^2)$  parametrically with

$$G(x, L) = \frac{\sqrt{x}}{\exp(p^2)\text{erfc}(p)} \quad q^2 = \mu^2 \left( \frac{\text{erfc} p}{\text{erfc} p_0} \right)^{1/2}$$

Other physical perspectives: <http://web.mit.edu/redingtn/www/netadv/Xdysonschw.html>

5

## Dyson-Schwinger equations and $B_+$

The key is  $B_+$ .

All the Hopf algebras we're interested in are generated by one or more  $B_+$  and so are the solutions of Dyson-Schwinger equations or quotients thereof.

$B_+$  is a 1-cocycle

$$\Delta B_+ = (\text{id} \otimes B_+) \Delta + B_+ \otimes \mathbb{I}$$

A subpiece comes from the branches, or is the whole thing. Unique decomposition.

$(\mathcal{H}_{rt}, B_+)$  is universal for Hopf algebras with a 1-cocycle. Connes, Kreimer: [4].

6

## $B_+$ and the universal law

The 1-cocycle property is the cohomological way to say unique decomposition.

Rooted trees are nice due to the unique decomposition of a tree into its root and the forest of its subtrees:  $B_+$ . For unlabelled trees,  $\mathbf{T}(x) = \sum t(n)x^n$ ,

$$\mathbf{T}(x) = x \exp \left( \sum_{m \geq 1} \mathbf{T}(x^m)/m \right).$$

Which by Pólya's classical analysis gives the asymptotics

$$t(n) \sim C \rho^{-n} n^{-3/2}$$

Asymptotics of the form  $C \rho^{-n} n^{-3/2}$  are ubiquitous for classes of rooted trees with recursive definitions, hence the term universal law.

7

## Operators giving the universal law

How ubiquitous? Let  $\mathcal{O}$  be the set of operators on power series built out of

1.  $\mathbf{E}(x, \cdot)$  such that
  - (a)  $\mathbf{E}(x, y)$  has nonnegative coefficients and zero constant term,
  - (b)  $\mathbf{E}(a, b) < \infty \Rightarrow \exists \epsilon > 0, \mathbf{E}(a + \epsilon, b + \epsilon) < \infty$ ,
  - (c)  $\exists R > 0, [x^i y^j] \mathbf{E}(x, y) \leq R^{i+j}$ .
2.  $\text{MSet}_M$  and  $\text{Seq}_M$  for all  $M \subseteq \mathbb{Z}^{>0}$ .
3.  $\text{DCycle}_M$  and  $\text{Cycle}_M$  for  $\sum_{m \in M} 1/m = \infty$  or  $M$  finite.

using scalar multiplication from  $\mathbb{R}^{\geq 0}$ , addition, multiplication, and composition, and where if  $\text{MSet}_M$ ,  $\text{DCycle}_M$ , or  $\text{Cycle}_M$  appear then scalars and coefficients of  $\mathbf{E}$  must be integers.

8

**Theorem 1. [Bell, Burris, – [1]]** Let  $\Theta \in \mathcal{O}$  such that

- $\Theta$  is nonlinear
- $[x^n] \Theta(\mathbf{A}(x))$  depends only on  $[x^i] \mathbf{A}(x)$  for  $i < n$ .

Let  $\mathbf{A}(x)$  be a power series

- with nonnegative coefficients
- with zero constant term
- which diverges at its radius of convergence
- if  $\text{MSet}_M$ ,  $\text{DCycle}_M$ , or  $\text{Cycle}_M$  appear in  $\Theta$  then  $\mathbf{A}(x)$  has integer coefficients.

Then there is a unique  $\mathbf{T}(x)$  satisfying

$$\mathbf{T}(x) = \mathbf{A}(x) + \Theta(\mathbf{T}(x)).$$

The coefficients of  $\mathbf{T}$  satisfy the universal law on their support.

9

## $B_+$ and the first recursion

For an analytic Dyson-Schwinger equation write

$$G(x, L) = \sum \gamma_k(x) L^k \quad \gamma_k = \sum_{j \geq k} \gamma_{k,j} x^j$$

The Hochschild closedness of  $B_+$  is what permits us to rewrite the linearized coproduct which along with  $S \star Y$  gives the recursion ([5])

$$\gamma_k(x) = \frac{1}{k} \gamma_1(x) (1 + rx \partial_x) \gamma_{k-1}(x)$$

10

## $B_+$ and the second recursion

Again

$$G(x, L) = \sum \gamma_k(x) L^k \quad \gamma_k = \sum_{j \geq k} \gamma_{k,j} x^j$$

The properties of  $B_+$  don't care about connectedness which permits us to modify the primitives of the theory to

- reduce to one insertion place; univariate Mellin transforms.
- take away higher order behaviour of Mellin transforms; geometric series Mellin transforms.

which along with the other recursion gives ([6])

$$\gamma_{1,n} = p(n) + \sum_{j=1}^{n-1} (-rj - 1) \gamma_{1,j} \gamma_{1,n-j}$$

11

## $B_+$ and the growth of $\gamma_1$

$$\gamma_{1,n} = p(n) + \sum_{j=1}^{n-1} (-r^j - 1) \gamma_{1,j} \gamma_{1,n-j}$$

is what we were able to analyze to show that the primitives determine the growth of the whole theory.

In particular Lipatov bounds  $\gamma_{1,n} \leq c^n n!$  carry over.

## $B_+$ and sub Hopf algebras

Today's punchline, solutions to Dyson-Schwinger equations are sub Hopf algebras. Bergbauer, Kreimer [2].

In the example

$$X = \mathbb{I} + xB_+(X^2)$$

$$c_0 = \mathbb{I} \quad c_1 = \bullet \quad c_2 = 2\mathbb{I} \quad c_3 = \wedge + 4\mathbb{I}$$

$$c_4 = 4\mathbb{I} + 2\wedge + 8\mathbb{I}$$

check

$$\begin{aligned} \Delta c_4 &= 4(\mathbb{I} \otimes \mathbb{I} + \mathbb{I} \otimes \mathbb{I} + \bullet \otimes \bullet + \bullet \otimes \wedge + \bullet \otimes \mathbb{I} + \mathbb{I} \otimes \bullet + \mathbb{I} \otimes \bullet) \\ &\quad + 2(\wedge \otimes \mathbb{I} + \mathbb{I} \otimes \wedge + 2 \bullet \otimes \bullet + \bullet \otimes \mathbb{I} + \wedge \otimes \bullet) \\ &\quad + 8(\mathbb{I} \otimes \mathbb{I} + \mathbb{I} \otimes \bullet + \bullet \otimes \mathbb{I} + \bullet \otimes \bullet) \\ &= c_4 \otimes c_0 + c_0 \otimes c_4 + (2c_3 + 2c_1 c_2) \otimes c_1 \\ &\quad + (3c_1^2 + 3c_2) \otimes c_2 + 4c_1 \otimes c_3 \end{aligned}$$

12

13

## The sub Hopf algebra result

Let  $B_+^{d_n}$  be Hochschild 1-cocycles. Consider

$$X = \mathbb{I} + \sum x^n w_n B_+^{d_n}(X^{n+1})$$

write  $X = \sum x^n c_n$ . Then the Dyson-Schwinger equation has a unique solution

$$c_n = \sum w_m B_+^{d_m} \sum_{\substack{k_1 + \dots + k_m = n-m \\ k_i \geq 0}} c_{k_1} \dots c_{k_m}$$

and the  $c_n$  generate a sub Hopf algebra

$$\Delta c_n = \sum_{k=0}^n P_k^n \otimes c_k$$

where the  $P_k^n$  are homogeneous polynomials of degree  $n - k$  in the  $c_i$ , specifically

$$P_k^n = \sum_{\ell_1 + \dots + \ell_{k+1} = n-k} c_{\ell_1} \dots c_{\ell_{k+1}}$$

14

## The role of $B_+$ for the sub Hopf algebras

Bergbauer and Kreimer [2] give a very natural operadic proof and an elementary proof consisting of a triple induction.

The inductive proof has the advantage of showing explicitly the use of the Hochschild 1-cocycle property of  $B_+$  and that no deep facts are needed.

15

## References

- [1] Jason Bell, Stanley Burris, and Karen Yeats, Counting Rooted Trees. *Elec. J. Combin.* **13** (2006), #R63. (Also arXiv:math.CO/0512432.)
- [2] C. Bergbauer and D. Kreimer, Hopf algebras in renormalization theory. *IRMA Lect. Math. Theor. Phys.* **10** (2006), 133-164. (Also arXiv:hep-th/0506190.)
- [3] D.J. Broadhurst and D. Kreimer, Exact solutions of Dyson-Schwinger equations . . . . *Nucl.Phys. B* **600**, (2001), 403-422. (Also arXiv:hep-th/0012146).
- [4] A. Connes and D. Kreimer. Hopf algebras, renormalization and noncommutative geometry. *Comm. Math. Phys.* **199** (1998), 203-242. (Also arXiv:hep-th/9808042)
- [5] Dirk Kreimer and Karen Yeats, An Étude in non-linear Dyson-Schwinger Equations. *Nucl. Phys. B Proc. Suppl.*, **160**, (2006), 116-121. (Also arXiv:hep-th/0605096.)
- [6] Dirk Kreimer and Karen Yeats, Recursion and Growth Estimates in Renormalizable Quantum Field Theory. arXiv:hep-th/0612179.