

Using renormalization Hopf algebra intuition on symmetric functions

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Why am I talking in this session?

and with a section called “Renormalization Hopf algebras”!?

In 2014-2015 I was on sabbatical.
I spent some of that time talking to Steph.

I wanted to understand why the renormalization Hopf algebra world seemed to have a different flavour than the rest of the combinatorial Hopf algebra world.

One paper came out of this [arXiv:1511.06337](https://arxiv.org/abs/1511.06337).
It had errors (all due to me).
Now it is corrected and has appeared: E-JC **24**(3) (2017) #P3.10.

The Connes-Kreimer Hopf algebra

Take a polynomial algebra of rooted trees and use the following coproduct:

$$\Delta(t) = \sum_{\substack{S \subseteq V(t) \\ S \text{ antichain}}} \left(\prod_{s \in S} t_s \right) \otimes \left(t \cdot \prod_{s \in S} t_c \right)$$

where $t_s =$ subtree rooted at s

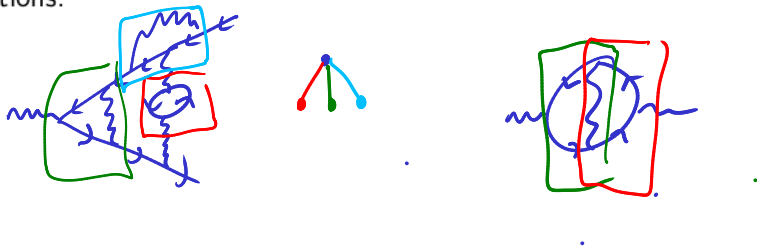
Eg:

$$\Delta(\text{rooted tree}) = 1 \otimes \text{rooted tree} + \dots + \text{rooted tree} \otimes 1 + 2 \cdot \dots$$

Renormalization Hopf algebras of Feynman diagrams

Feynman diagrams are graphs which tell a story of particle interactions.

Eg:



Take a polynomial algebra of them and use the following coproduct:

$$\Delta(G) = \sum_{\gamma} \gamma \otimes G/\gamma$$

product of divergent 1PI subgraphs

Important things from this perspective

- Add-a-root/graph insertion, Hochschild 1-cocycles.
- SubHopf algebras with one generator in each degree.
- Combinatorial specifications/Dyson-Schwinger equations.
- Feynman rules/character maps.
- Iterated coproduct, and pulling out particular terms.

Shape Hopf algebra

Use two layers.

Take the polynomial algebra generated by ^{connected} shapes, including skew.

Identify disconnected shapes with the monomial of components.

Make the coproduct match Schur functions:

~~$$\Delta(\lambda/\mu) = \sum_{\mu \leq \eta \leq \lambda} (\eta/\mu) \otimes (\lambda/\eta)$$~~

$$\Delta(\lambda/\mu) = \sum_{\mu \leq \eta \leq \lambda} (\eta/\mu) \otimes (\lambda/\eta)$$

And nothing else.

... and symmetric functions

The shape Hopf algebra is the Hopf algebra of intervals in Young's lattice.

It is not cocommutative.

The symmetric function Hopf algebra is a quotient.

It is a cocommutative quotient.

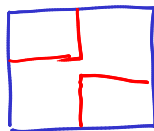
That's all we need to know.

A complicated proof of an easy fact

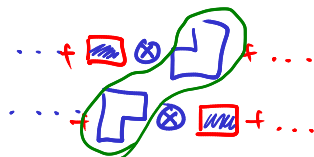
If α is a shape let α^* be the 180° rotation.

Their Schur functions are the same, written $\alpha \sim \alpha^*$

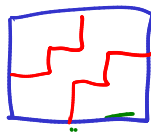
Proof.



$$\Delta(\square) = \dots$$



repeat

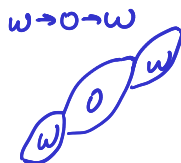
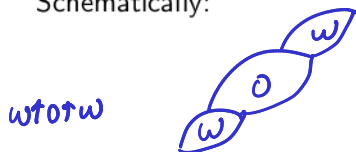


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Cocommutativity was the key.

McNamara, van Willigenburg WOW shapes

Schematically:

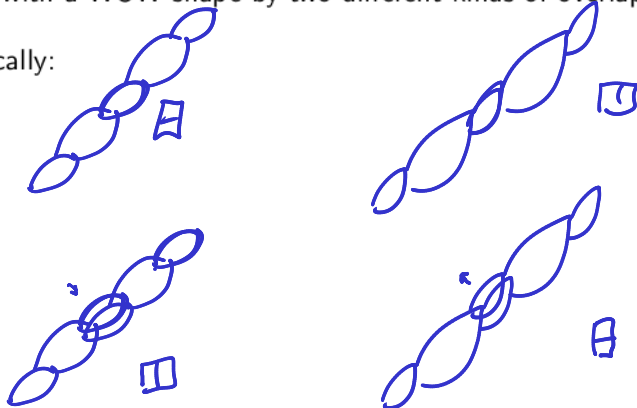


(These are only two of their four cases.)

Composition

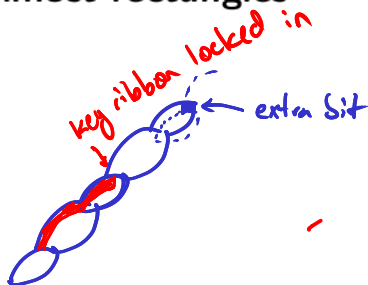
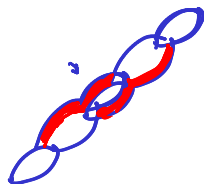
Compose with a WOW shape by two different kinds of overlap.

Schematically:



Key ribbons and almost rectangles

Schematically:



Note that the top and bottom key ribbons may differ.

Result

Theorem


Let β be a partition shape that is a rectangle with the lower right corner box removed. Let γ be a $W \rightarrow O \rightarrow W$ or $W \uparrow O \uparrow W$ shape with no loose end ribbons. Then

$$\beta \circ \gamma \sim \beta^* \circ \gamma.$$

What are loose end ribbons?

Prove the result by cocommutativity pulling out the key ribbon \otimes stuff and stuff \otimes key ribbon terms (with many details).

Conclusion

- This is much more specialized than what McNamara and van Willigenburg can prove about WOW-built skew Schur function identities.
- It can prove some cases that they can't.
- Generalizations? 
- What about all those other renormalization Hopf algebra things?

