

Syllabus:

When QM meets information theory (wk 1-3)

QM & immediate consequences for info processing

Noiseless Q computation of classical problems (wk 4-7)

Q circuits, universality, Q algorithms

* Q computation in the presence of noise (wk 8-10)

(8) Modelling noise: mixed state QM & Q operations

(9) Quantum error correction

(10) Reliable Q computation using noisy components

Q computation in the presence of adversary (wk 11-12)

Q cryptography

8. Modelling noise: mixed state QM & Q operations

States:

- (b) States on composite systems (NC 8.3.1, 2.5, KLM 3.5.2)
- (a) Noisy quantum data (NC 2.4, KLM 3.5.1)

Evolution:

- (c) Most general (noisy) quantum dynamics
- (d) Characterizations (NC 8.2, KLM 3.5.3)
- (e) Important examples (NC 8.3)

Measurements (reading exercise):

- (f) POVM measurements (NC 2.2.6, KLM A8)
- (g) Trace distance, indistinguishability,
Helstrom-Holevo theorem (NC 9.2, KLM A8)

(a) Noisy quantum data on S alone

(b) States on composite systems R&S

(b)

Pure state
QM on 2
systems RS

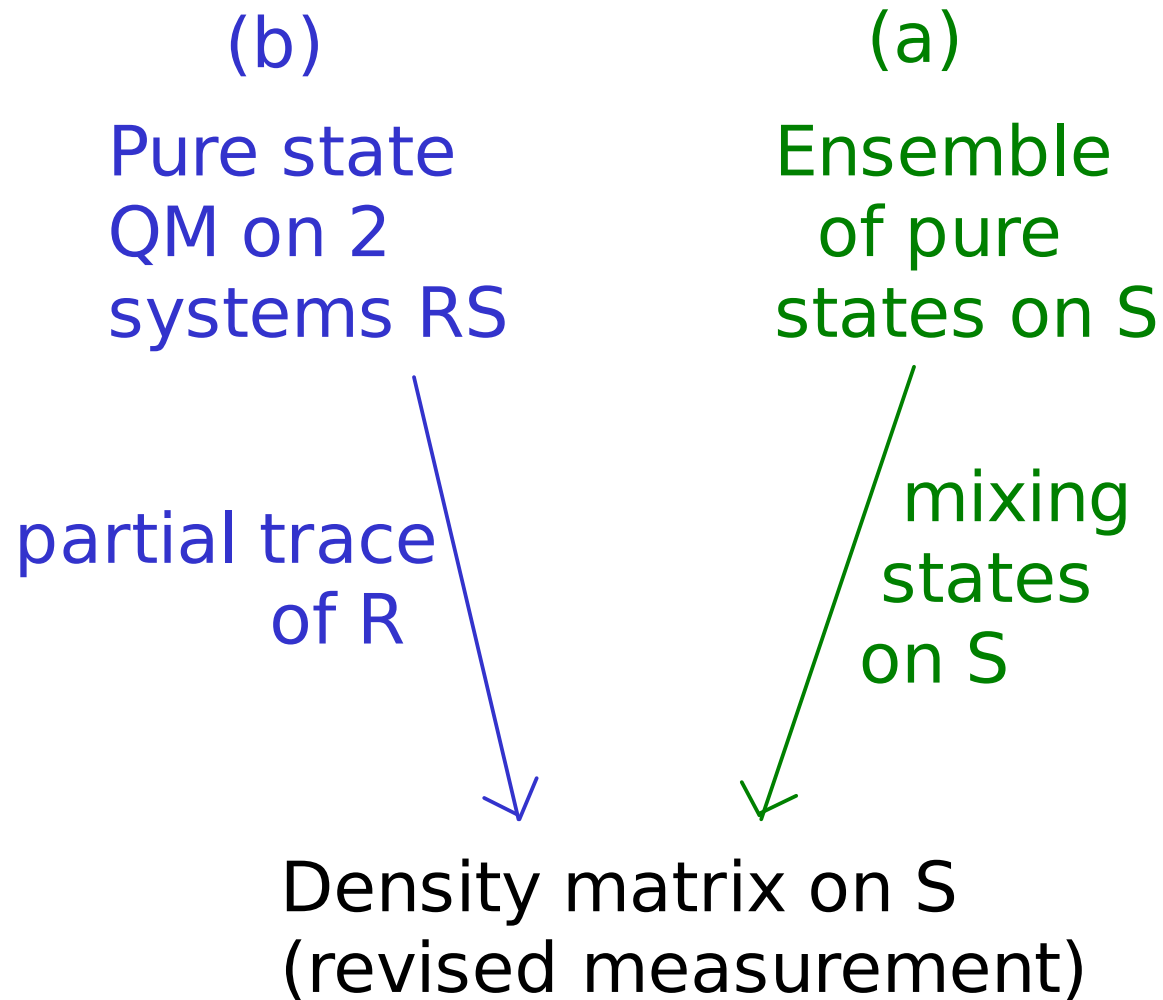
partial trace
of R



Density matrix on S
(revised measurement)

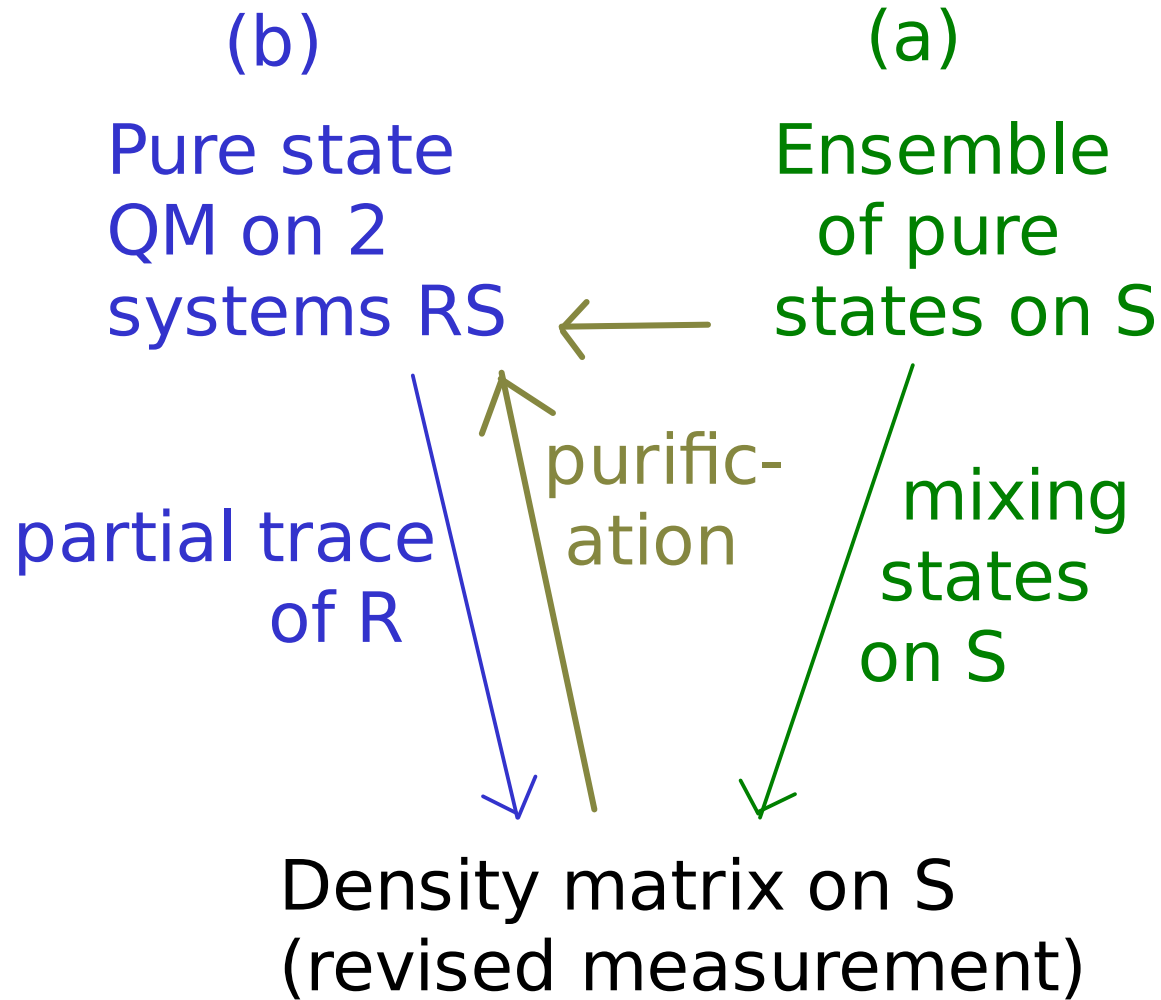
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(b) States on composite systems R&S

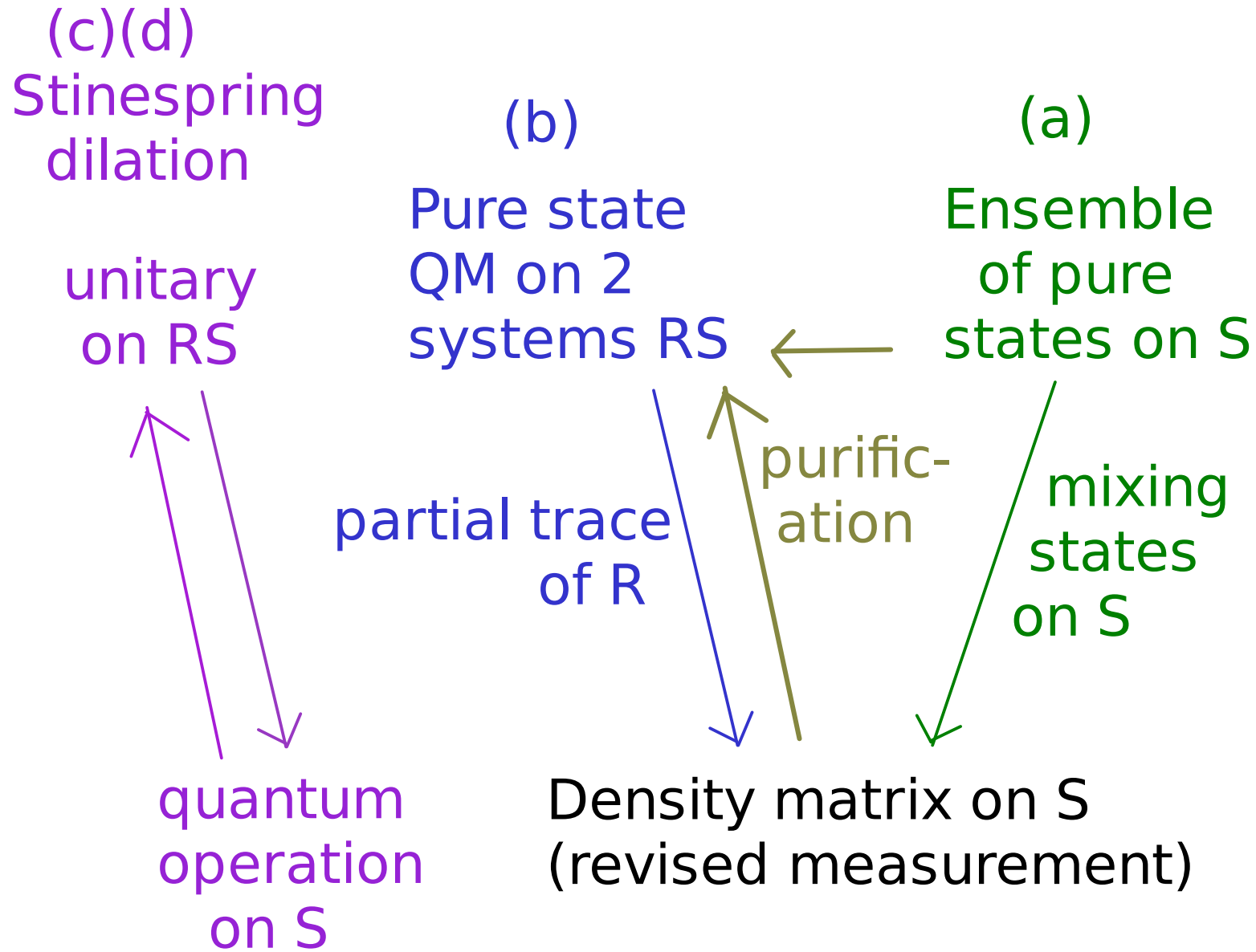


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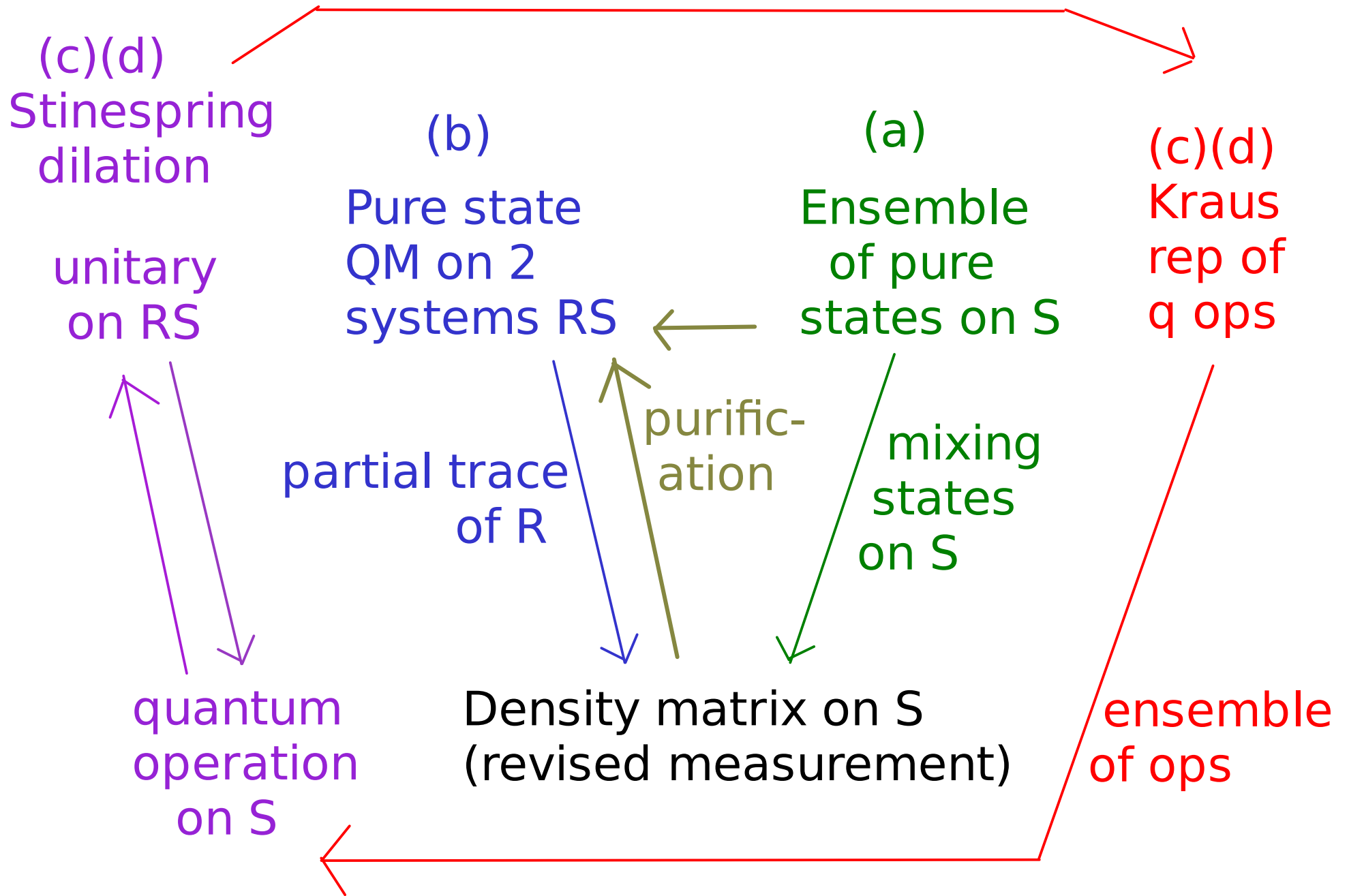
(b) States on composite systems R&S



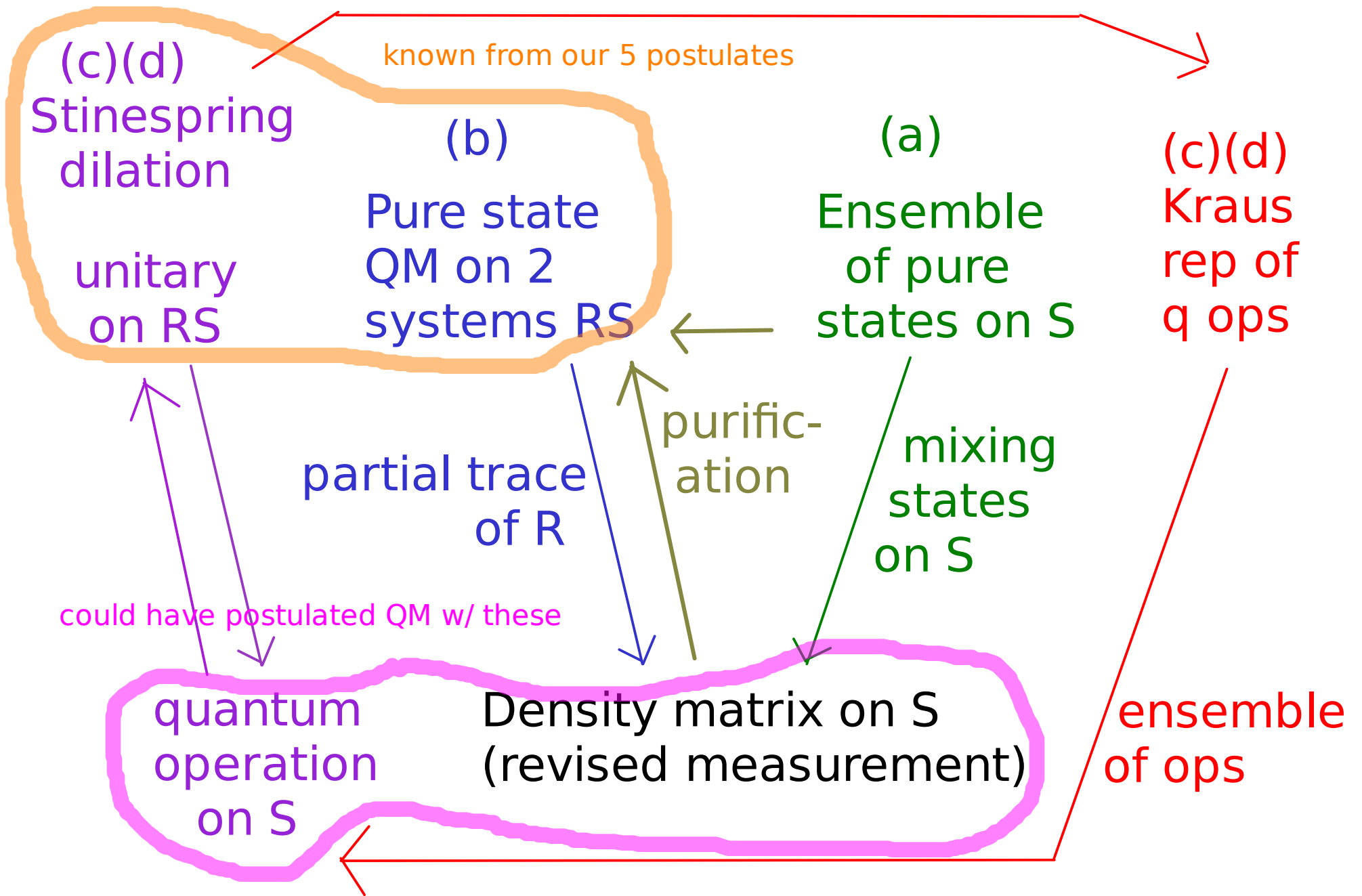
(c)(d) Most general (noisy) quantum dynamics



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(c)(d) Most general (noisy) quantum dynamics



either way, once proved, we use all of the above methods

(b)

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QM on 2
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Density matrix on S
(revised measurement)

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but we operate (unitaries and measurements) on S .

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Lemma: for any column vector v , $\|v\|^2 = \text{tr} v v^t$.

Proof:

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Lemma: for any column vector v , $\|v\|^2 = \text{tr } v v^\dagger$.

Proof: $v = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_d \end{pmatrix}$, $v v^\dagger = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_d \end{pmatrix} \begin{pmatrix} v_1^* & v_2^* & \dots & v_d^* \end{pmatrix} = \begin{pmatrix} v_1 v_1^* & \dots & v_1 v_d^* \\ v_2 v_1^* & \dots & v_2 v_d^* \\ \vdots & \dots & \vdots \\ v_d v_1^* & \dots & v_d v_d^* \end{pmatrix}$

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$$\text{tr} v v^\dagger = v_1 v_1^* + \dots + v_d v_d^* = \|v\|^2$$

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$$= \text{tr} P_k (\nu_0 \nu_0^\dagger + \nu_1 \nu_1^\dagger)$$

where $\nu_0 = \langle 0| \otimes I |\Psi\rangle = \langle 0| \otimes I (a|00\rangle + b|11\rangle) = a|0\rangle_S$
 $\nu_1 = \langle 1| \otimes I |\Psi\rangle = \langle 1| \otimes I (a|00\rangle + b|11\rangle) = b|1\rangle_S$

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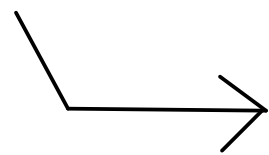
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revised
formula
for prob
of meas

revised way to rep
quantum data on S



$$= \text{tr } P_k \left(a^2 |0\rangle\langle 0| + b^2 |1\rangle\langle 1| \right) \text{ (tr linear)}$$

$$= \text{tr } P_k \left(v_0 v_0^\dagger + v_1 v_1^\dagger \right)$$

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General: $|\Psi\rangle = \sum_i a_i |\bar{i}\rangle |\eta_i\rangle$ on RS.

real ortho-normal unit vector on S

(derived in topic03-02.pdf p5-8)

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revised
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$$= \text{tr} P_k \left(\underbrace{\sum_i a_i^2 |\eta_i\rangle \langle \eta_i|}_{\rho_S} \right) \text{ where } a_i |\eta_i\rangle_S = \langle i| \otimes I |\Psi\rangle.$$

ρ_S : density matrix on S
dxd if $d = \dim(S)$

revised rep
for quantum
data on S

trace 1, positive semidefinite

The partial trace

Revised formulation of QM:

Revised description of quantum state:

$$|\Psi\rangle = \sum_i a_i |\bar{i}\rangle |\eta_i\rangle \longrightarrow |\Psi\rangle\langle\Psi| \longrightarrow \sum_i a_i^2 |\eta_i\rangle\langle\eta_i| = \rho_S$$

1. outer product 2. partial trace

revised description of measurement:

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Define partial trace (describing a state on S from a state on RS) so postmeasurement states & dynamics also makes sense.

The partial trace

$d \times d$

Recall the trace of a matrix M is the sum of all the diagonal elements. In the Dirac notation:

$$\text{tr } M = \text{tr} \left(M \sum_{i=1}^d |i\rangle\langle i| \right) = \sum_{i=1}^d \langle i | M | i \rangle$$

insert identity

$\{|i\rangle\}$ basis

tr is cyclic and linear

- in terms of a basis
- can use any basis

basis independent

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insert identity
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d dim, basis $\{|i\rangle\}$

Definition: the partial trace of system B, denoted tr_B , is defined on matrices acting on systems AB as

$$\text{tr}_B M = \sum_{i=1}^d \underbrace{\left(\begin{array}{c} I \\ A \end{array} \otimes \begin{array}{c} \langle i | \\ B \end{array} \right) M \left(\begin{array}{c} I \\ A \end{array} \otimes |i\rangle \right)}_{d_A \times d_A} \quad d_A d \times d_A d$$

d_A dim

i.e., trace B and do nothing on A

Question:

Suppose $|\Psi\rangle_{RS} = a|00\rangle + b|1+\rangle$ and we apply partial trace to R, what is the density matrix on S?

(a) $\begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix}$ (b) $\begin{pmatrix} a^2 + \frac{b^2}{2} & \frac{b^2}{2} \\ \frac{b^2}{2} & \frac{b^2}{2} \end{pmatrix}$ (c) $\begin{pmatrix} a^2 & \frac{ab}{\sqrt{2}} \\ \frac{ab}{\sqrt{2}} & b^2 \end{pmatrix}$

Please do not scroll down ...

Useful from previous slides:

$$|\Psi\rangle = \sum_i a_i |i\rangle |\eta_i\rangle \longrightarrow |\Psi\rangle\langle\Psi| \longrightarrow \sum_i a_i^2 |\eta_i\rangle\langle\eta_i| = \rho_S$$

1. outer product 2. partial trace

$$\text{tr}_B M = \sum_{i=1}^d \underbrace{(\mathbb{I} \otimes \langle i|)}_{\substack{| \\ A \\ | \\ B}} M (\mathbb{I} \otimes |i\rangle)$$

$d_A d \times d_A d$

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Exercise: show that, if we apply the partial trace to S instead, the density matrix, now on R, is (c)! Hint: rewrite the joint state on RS using method 2 with the basis $\{|0\rangle, |1\rangle\}$ for S.

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NB: partial tracing R is different from partial tracing S! Outcomes even have different dims if $\dim(R) \neq \dim(S)$.

Exercise: show that partial tracing R or S from any of the 4 Bell states give density matrices equal to $I/2$!

$$|\Phi_0\rangle_{RS} = \frac{|00\rangle + |11\rangle}{\sqrt{2}}$$

$$|\Phi_z\rangle_{RS} = \frac{|00\rangle - |11\rangle}{\sqrt{2}}$$

$$|\Phi_x\rangle_{RS} = \frac{|01\rangle + |10\rangle}{\sqrt{2}}$$

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So, the partial trace is a many-to-one mapping, and is mathematically irreversible.

Crucial in cryptography (later in the course) to learn when two states have the same partial trace, and if you see system S, what do you know about RS jointly.

We have defined the partial trace in the dirac notation.

Next we derive how it looks in matrix representation.

The partial trace (example for 2 qubits)

$$\mathbb{I} \otimes \langle 0| = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes [1 \ 0] = \begin{pmatrix} [1 \ 0] & [0 \ 0] \\ [0 \ 0] & [1 \ 0] \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$\mathbb{I} \otimes \langle 1| = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes [0 \ 1] = \begin{pmatrix} [0 \ 1] & [0 \ 0] \\ [0 \ 0] & [0 \ 1] \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

The partial trace (example for 2 qubits)

$$I \otimes \langle 01 | = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes [1 \ 0] = \begin{pmatrix} [1 \ 0] & [0 \ 0] \\ [0 \ 0] & [1 \ 0] \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

$$I \otimes \langle 11 | = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes [0 \ 1] = \begin{pmatrix} [0 \ 1] & [0 \ 0] \\ [0 \ 0] & [0 \ 1] \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$(I \otimes \langle 01 |) M (I \otimes | 0 \rangle) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \\ 0 & 1 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} m_{11} & m_{13} \\ m_{31} & m_{33} \end{pmatrix}$$

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$$(I \otimes \langle 0|) M (I \otimes |0\rangle) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix} \begin{pmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \\ 0 & 1 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} m_{11} & m_{13} \\ m_{31} & m_{33} \end{pmatrix}$$

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$$\text{tr}_B M = \sum_{i=1}^d (I \otimes \langle i |) M (I \otimes |i\rangle) = \begin{pmatrix} m_{11} + m_{22} & m_{13} + m_{24} \\ m_{31} + m_{42} & m_{33} + m_{44} \end{pmatrix}$$

$$\text{tr}_B M = \begin{pmatrix} \text{tr} \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} & \text{tr} \begin{pmatrix} m_{13} & m_{14} \\ m_{23} & m_{24} \end{pmatrix} \\ \text{tr} \begin{pmatrix} m_{31} & m_{32} \\ m_{41} & m_{42} \end{pmatrix} & \text{tr} \begin{pmatrix} m_{33} & m_{34} \\ m_{43} & m_{44} \end{pmatrix} \end{pmatrix} \xleftarrow{\text{tracing each block}} \begin{pmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{pmatrix} = M$$

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consistent with the idea entries in each block
 rep states in B while the block labels corr to
 states in A

	00	01	10	11
00	$\begin{pmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{pmatrix}$			
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Exercise:

$$\text{tr}_A M = \text{tr}_A \begin{pmatrix} m_{11} & m_{12} & m_{13} & m_{14} \\ m_{21} & m_{22} & m_{23} & m_{24} \\ m_{31} & m_{32} & m_{33} & m_{34} \\ m_{41} & m_{42} & m_{43} & m_{44} \end{pmatrix} = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} + \begin{pmatrix} m_{33} & m_{34} \\ m_{43} & m_{44} \end{pmatrix}$$

summing diagonal blocks

Example: A, B are 3- and 2-dim respectively. (M: 6x6)

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Remark:

The trace of an r -dim system is a linear map from $r \times r$ matrices to real numbers.

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Remark:

The trace of an r -dim system is a linear map from $r \times r$ matrices to real numbers.

The partial trace of an r -dim system is a linear map from $r_s \times r_s$ matrices to $s \times s$ matrices where the trace is applied to R , and the identity map on S .

It acts on tensor product matrices as:

$$\text{tr}_R M_R \otimes M_S = \underbrace{(\text{tr } M_R)}_{\text{scalar}} \cdot \underbrace{M_S}_{\text{scalar product}}$$

and extends to any $r_s \times r_s$ matrix (by linearity).

Revised formulation of QM:

Revised description of quantum state:

$$|\Psi\rangle = \sum_{\bar{i}} a_i |\bar{i}\rangle |\eta_i\rangle \longrightarrow |\Psi\rangle\langle\Psi| \longrightarrow \sum_{\bar{i}} a_i^2 |\eta_i\rangle\langle\eta_i| = \rho_S$$

1. outer product 2. partial trace

Checking the new definition of partial trace on $|\Psi\rangle\langle\Psi|$:

$$\text{Tr}_R |\Psi\rangle\langle\Psi| = \sum_{\bar{i}=1}^{\dim R} (\langle\bar{i}|\otimes I) |\Psi\rangle\langle\Psi| (|\bar{i}\rangle\otimes I)$$

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$$\text{So } \text{tr}_R |\Psi\rangle\langle\Psi| = \sum_{\bar{i}=1}^{\dim R} a_i^2 |\eta_i\rangle\langle\eta_i| \text{ same as wanted above.}$$

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$$\text{tr}_R |\Psi_k\rangle\langle\Psi_k| = \sum_{i=1}^{\dim R} \frac{a_i P_k |\eta_i\rangle}{\|I \otimes P_k |\Psi\rangle\|} \frac{a_i \langle\eta_i| P_k}{\|I \otimes P_k |\Psi\rangle\|}$$

$$\text{tr}_R |\Psi_K\rangle\langle\Psi_K| = \sum_{i=1}^{\dim R} \frac{a_i P_K |\eta_i\rangle}{\|I \otimes P_K |\Psi\rangle\|} \frac{a_i \langle\eta_i| P_K}{\|I \otimes P_K |\Psi\rangle\|}$$

$$= \sum_{i=1}^{\dim R} \frac{a_i^2 P_K |\eta_i\rangle\langle\eta_i| P_K}{\text{tr} P_K \beta_S}$$

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revised description of
post-measurement state

Revised formulation of QM:

Revised description of evolution by unitary U on S:

$$|\Psi\rangle = \sum_i a_i |\bar{i}\rangle |\eta_i\rangle$$

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$$\text{So, } \beta_S \longrightarrow U \beta_S U^\dagger.$$

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Revised rules coincide with alternative interpretation of density matrix as ensemble of quantum states (later).

We took our 5 postulates, applied to composite system RS and derived QM as viewed from S .

We now deduce some properties for the states and evolution/measurement.

We can instead start with these properties and formulate all QM from scratch based on them ...

Properties of the density matrix and partial trace:

1. For any $|\psi\rangle$ on RS , for any unitary U on R ,

$$\text{tr}_R U_R \otimes I_S |\psi\rangle\langle\psi| U_R^\dagger \otimes I_S = \text{tr}_R |\psi\rangle\langle\psi|$$

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Remarks:

(a) This restates the non-signalling principle in the presence of entanglement: if 1 doesn't hold, a party holding R can affect the density matrix of S and its measurement statistics, thereby partially communicating whether U is applied or not.

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Remarks:

- (a) This restates the non-signalling principle in the presence of entanglement: if 1 doesn't hold, a party holding R can affect the density matrix of S and its measurement statistics, thereby partially communicating whether U is applied or not.
- (b) Or, the above can be proved directly as an alternative proof to the non-signalling principle.

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Proof: Let r, s be dim of R, S . $I_S \otimes U_R =$

U	0	0	
0	U		
0			
\vdots			
			U

$\leftarrow r \rightarrow$
 $\leftarrow s \text{ blocks} \rightarrow$

note the switch
of the ordering

Properties of the density matrix and partial trace:

1. For any $|\psi\rangle$ on RS , for any unitary U on R ,

$$\text{tr}_R U_R \otimes I_S |\psi\rangle\langle\psi| U_R^\dagger \otimes I_S = \text{tr}_R |\psi\rangle\langle\psi|$$

Proof: Let r, s be dim of R, S . $I_S \otimes U_R =$

$$M = \begin{pmatrix} M_{11} & M_{12} & M_{13} & \dots \\ M_{21} & M_{22} & & \\ M_{31} & & & \\ \vdots & & & \\ \vdots & & & M_{ss} \end{pmatrix}, \quad M_{ij} : r\text{-by-}r \text{ matrices.}$$

$$I_S \otimes U_R = \begin{pmatrix} U & 0 & 0 & & \\ 0 & U & & & \\ 0 & & & & \\ \vdots & & & & \\ \vdots & & & & U \end{pmatrix}$$

← s blocks →

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$$\begin{pmatrix} U & 0 & 0 & & \\ 0 & U & & & \\ 0 & & & & \\ \vdots & & & & \\ \vdots & & & & U \end{pmatrix}$$

$\leftarrow r \rightarrow$
 $\leftarrow s \text{ blocks} \rightarrow$

$$(I_S \otimes U_R) M (I_S \otimes U_R^\dagger) = \begin{pmatrix} U & 0 & 0 & & \\ 0 & U & & & \\ 0 & & & & \\ \vdots & & & & \\ \vdots & & & & U \end{pmatrix} \begin{pmatrix} M_{11} & M_{12} & M_{13} & \dots & \\ M_{21} & M_{22} & & & \\ M_{31} & & & & \\ \vdots & & & & \\ \vdots & & & & M_{ss} \end{pmatrix} \begin{pmatrix} U^\dagger & 0 & 0 & & \\ 0 & U^\dagger & & & \\ 0 & & & & \\ \vdots & & & & \\ \vdots & & & & U^\dagger \end{pmatrix}$$

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U	0	0	
0	U		
0			
\vdots			
			U

$M =$

M_{11}	M_{12}	$M_{13} \dots$	
M_{21}	M_{22}		
M_{31}			
\vdots			
			M_{ss}

, M_{ij} : r -by- r matrices.

$\leftarrow r \rightarrow$

U	0	0	
0	U		
0			
\vdots			
			U

$\leftarrow s \text{ blocks} \rightarrow$

$$(I_S \otimes U_R) M (I_S \otimes U_R^\dagger) =$$

U	0	0	
0	U		
0			
\vdots			
			U

M_{11}	M_{12}	$M_{13} \dots$	
M_{21}	M_{22}		
M_{31}			
\vdots			
			M_{ss}

U^\dagger	0	0	
0	U^\dagger		
0			
\vdots			
			U^\dagger

$$=$$

$U M_{11} U^\dagger$	$U M_{12} U^\dagger$	$U M_{13} U^\dagger$	\dots
$U M_{21} U^\dagger$	$U M_{22} U^\dagger$	$U M_{23} U^\dagger$	\dots
$U M_{31} U^\dagger$	\dots		
\vdots			
			$U M_{ss} U^\dagger$

(exercise)

$$\text{tr}_R (I_S \otimes U_R) M (I_S \otimes U_R^\dagger) = \text{tr}_R \left(\begin{array}{c|c|c|c} UM_{11}U^\dagger & UM_{12}U^\dagger & UM_{13}U^\dagger & \dots \\ \hline UM_{21}U^\dagger & UM_{22}U^\dagger & UM_{23}U^\dagger & \dots \\ \hline UM_{31}U^\dagger & \dots & & \\ \vdots & & & \\ & & & UM_{ss}U^\dagger \end{array} \right)$$

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$$= \left(\begin{array}{c|c|c|c|c} \text{tr} UM_{11}U^\dagger & \text{tr} UM_{12}U^\dagger & \text{tr} UM_{13}U^\dagger & \dots & \\ \hline \text{tr} UM_{21}U^\dagger & \text{tr} UM_{22}U^\dagger & \text{tr} UM_{23}U^\dagger & \dots & \\ \hline \text{tr} UM_{31}U^\dagger & \dots & & & \\ \vdots & & & & \text{tr} UM_{ss}U^\dagger \end{array} \right)$$

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2. Corollary: the partial trace of R can be taken along any basis of R (just as the trace is basis indep).

Proof (reading exercise for W25):

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and any basis of R can be written as $\{U_R^\dagger |i\rangle\}_i$
for some U_R .

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2. Corollary: the partial trace of R can be taken along any basis of R (just as the trace is basis indep).

3. Partial tracing a system R has physical meaning of not accessing R . One example is losing access to R irreversibly (disgarded or corrupted by noise).

Properties of the density matrix and partial trace:

4. The partial trace is:

(a) linear

(b) trace preserving

(c) completely positive (i.e., applied to 1 out of 2 systems, it takes positive semidefinite (PSD) matrices to PSD matrices)

hermitian with non-negative eigenvalues

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$$(b) \operatorname{tr}_R M = \sum_{i=1}^r (\mathbb{I} \otimes \langle i |) M_{SR} (\mathbb{I} \otimes | i \rangle)$$

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matrix on S

vector on S

matrices on S

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vector on S

$$= \sum_{j=1}^s \sum_{i=1}^r \left(\underbrace{\langle j |}_S \otimes \underbrace{\langle i |}_R \right) M_{SR} \left(| j \rangle \otimes | i \rangle \right) = \operatorname{tr} M$$

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Proof: (c) if M is PSD, $M = \sum_j \lambda_j |\psi_j\rangle\langle\psi_j|$ where $\lambda_j > 0$
on RS

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$$\begin{aligned} \text{tr}_R M &= \text{tr}_R \sum_j \lambda_j |\psi_j\rangle\langle\psi_j| = \sum_j \lambda_j \text{tr}_R |\psi_j\rangle\langle\psi_j| \\ &= \sum_j \sum_{i=1}^r \lambda_j \underbrace{\mathcal{V}_{ij} \mathcal{V}_{ij}^\dagger}_{\text{PSD}} \geq 0 \end{aligned}$$

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6. Corollary: $\rho_S = \sum_K \mu_K |\Psi_K\rangle\langle\Psi_K|$ where $\mu_K \geq 0$, $\sum_K \mu_K = 1$
|
eigenvector of ρ_S

Properties of the density matrix and partial trace:

7. For $\rho_S = \sum_K \mu_K |\Psi_K\rangle\langle\Psi_K|$ where $\mu_K \geq 0$, $\sum_K \mu_K = 1$

a measurement on S has same statistics as drawing k with prob μ_K , preparing $|\Psi_K\rangle$, and measuring.

Proof:

Properties of the density matrix and partial trace:

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a measurement on S has same statistics as drawing k with prob μ_K , preparing $|\Psi_K\rangle$, and measuring.

Proof: if measurement is given by projectors $\{P_\ell\}$

$$\begin{aligned} \text{prob}(\ell) &= \text{tr } P_\ell \rho_S \\ &= \text{tr } P_\ell \sum_K \mu_K |\Psi_K\rangle\langle\Psi_K| \end{aligned}$$

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□

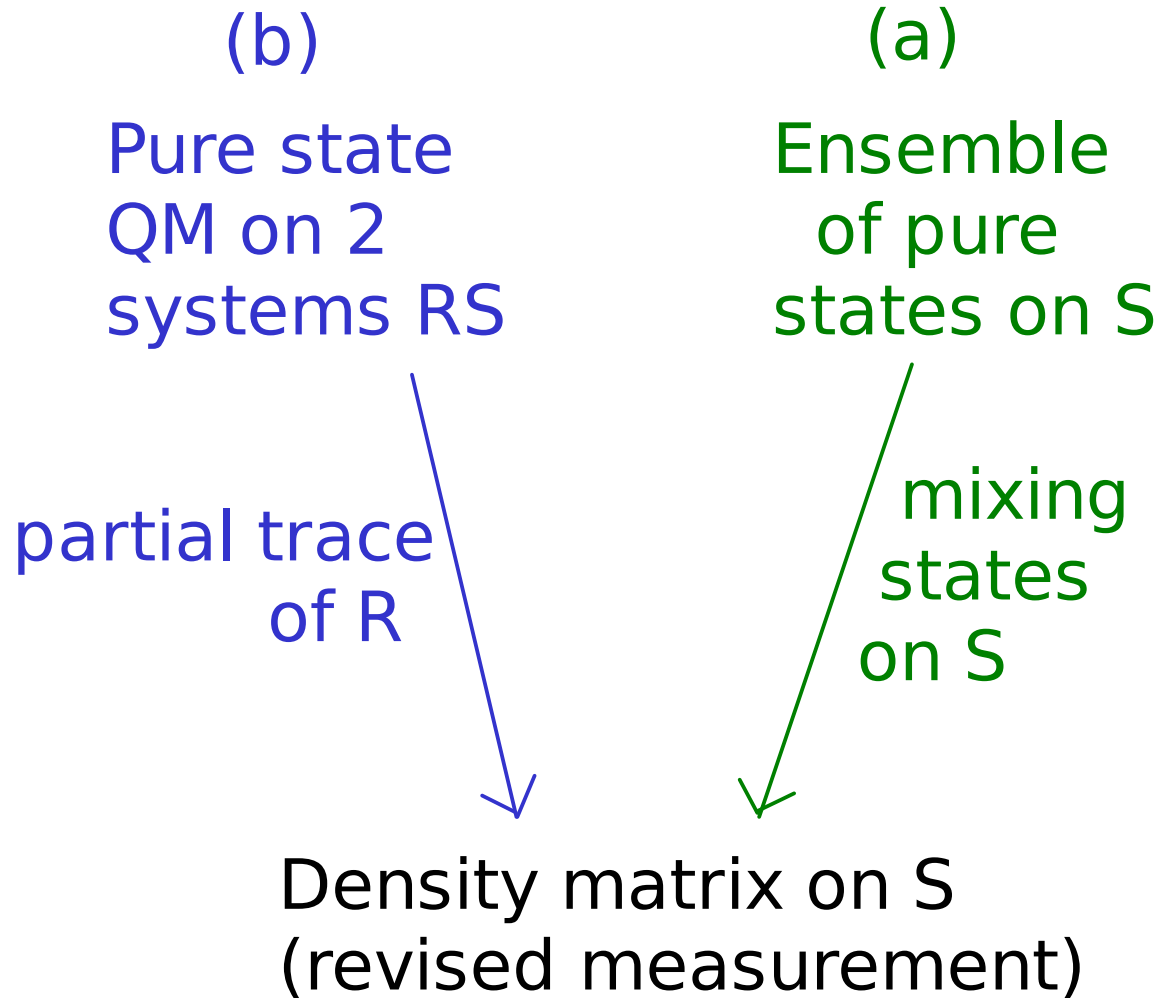
Property 7 gives a second interpretation of density matrix as a probabilistic mixture (or ensemble) of quantum states (later).

We took our 5 postulates, applied to composite system RS and derived QM as viewed from S .

✓ We now deduce some properties for the states and evolution/measurement.

→ We can instead start with these properties and formulate all QM from scratch based on them ...

Alternative approach to deriving the density matrix



Mixed state quantum mechanics

From vectors to density matrices

The density matrix of a pure state $|\psi\rangle$ is $\rho = |\psi\rangle\langle\psi|$.

Definition: if a system is in the state $|\psi_0\rangle$ with prob p_0 and $|\psi_1\rangle$ with prob p_1 , $p_0 + p_1 = 1$, then the state of the system is given by the density matrix

$$\rho = p_0 |\psi_0\rangle\langle\psi_0| + p_1 |\psi_1\rangle\langle\psi_1|.$$

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Terminology: this is called a mixed state.

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The definition extends to arbitrary mixtures, over any number of states, which can also be mixed.

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Terminology: this is called a mixed state.

The definition extends to arbitrary mixtures, over any number of states, which can also be mixed.

The prescription: "with probability p_i the state is $|\psi_i\rangle$ " is called an ensemble, or a probabilistic mixture, of states.

Example 1

If the state is $|0\rangle$ with prob $2/3$, $|+\rangle$ with prob $1/3$,
what is the density matrix ?

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Answer: $|0\rangle\langle 0| = \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}$

$$|+\rangle\langle +| = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \times \frac{1}{\sqrt{2}} [1 \ 1] = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$$

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So the density matrix of interest is

$$\rho = \frac{2}{3} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \frac{1}{3} \cdot \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} = \begin{bmatrix} \frac{5}{6} & \frac{1}{6} \\ \frac{1}{6} & \frac{1}{6} \end{bmatrix}$$

Example 1 (continued)

Applying spectral decomposition to ρ

$$\begin{aligned}\rho &= \begin{bmatrix} \frac{5}{6} & \frac{1}{6} \\ \frac{1}{6} & \frac{1}{6} \end{bmatrix} = \begin{bmatrix} 0.23 & 0.97 \\ -0.97 & 0.23 \end{bmatrix} \begin{bmatrix} 0.13 & 0 \\ 0 & 0.87 \end{bmatrix} \begin{bmatrix} 0.23 & 0.97 \\ -0.97 & 0.23 \end{bmatrix} \\ &= 0.13 \begin{bmatrix} 0.23 \\ -0.97 \end{bmatrix} \begin{bmatrix} 0.23 & -0.97 \end{bmatrix} + 0.87 \begin{bmatrix} 0.97 \\ 0.23 \end{bmatrix} \begin{bmatrix} 0.97 & 0.23 \end{bmatrix}\end{aligned}$$

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It means that, for a different mixture, where the state is $\begin{bmatrix} 0.23 \\ -0.97 \end{bmatrix}$ with prob 0.13, $\begin{bmatrix} 0.97 \\ 0.23 \end{bmatrix}$ with prob 0.87, the density matrix is also ρ .

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So, two different mixtures can give rise to the same density matrix.

Example 1 (continued)

If the state is $|0\rangle$ with prob $2/3$, $|+\rangle$ with prob $1/3$,
is the state same as $\sqrt{\frac{2}{3}}|0\rangle + \sqrt{\frac{1}{3}}|+\rangle$?

Vote: (a) yes, (b) no.

Example 1 (continued)

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Renormalizing $\sqrt{\frac{2}{3}}|0\rangle + \sqrt{\frac{1}{3}}|+\rangle$ gives $0.95|0\rangle + 0.32|+\rangle =: |\Psi\rangle$

$$|\Psi\rangle\langle\Psi| = \begin{bmatrix} 0.9 & 0.3 \\ 0.3 & 0.1 \end{bmatrix} \neq \rho = \begin{bmatrix} \frac{5}{6} & \frac{1}{6} \\ \frac{1}{6} & \frac{1}{6} \end{bmatrix} !$$

Example 1 (continued)

If the state is $|0\rangle$ with prob $2/3$, $|+\rangle$ with prob $1/3$,
is the state same as $\sqrt{\frac{2}{3}}|0\rangle + \sqrt{\frac{1}{3}}|+\rangle$?

Vote: (a) yes, (b) no.

Renormalizing $\sqrt{\frac{2}{3}}|0\rangle + \sqrt{\frac{1}{3}}|+\rangle$ gives $0.95|0\rangle + 0.32|1\rangle =: |\Psi\rangle$

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Note: $|\Psi\rangle \propto \sqrt{p_0}|4_0\rangle + \sqrt{p_1}|4_1\rangle$ and

$\rho = p_0|4_0\rangle\langle 4_0| + p_1|4_1\rangle\langle 4_1|$ are very different!

A superposition is a sum over vectors that exhibit interference. A mixture is a sum over density matrices, and the summands do not exhibit interference.

Alternative state postulate:

For a d -dim system, states are

(a) trace 1

(b) positive semidefinite

(c) $d \times d$

matrices.

By spectral decomp, states are ensembles of pure states.

Evolution of density matrices by unitaries and meas:

1. For an initial mixture:

$|\psi_0\rangle$ with prob p_0 , $|\psi_1\rangle$ with prob p_1

a unitary U evolves it to a new mixture:

$U|\psi_0\rangle$ with prob p_0 , $U|\psi_1\rangle$ with prob p_1

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The initial density matrix $\rho = p_0 |\psi_0\rangle\langle\psi_0| + p_1 |\psi_1\rangle\langle\psi_1|$
becomes $p_0 U|\psi_0\rangle\langle\psi_0|U^\dagger + p_1 U|\psi_1\rangle\langle\psi_1|U^\dagger = U\rho U^\dagger$.

2. For a measurement given by projectors $\{ P_j \}$,
the probability to get outcome j is

$$\text{Tr } P_j |\psi_0\rangle\langle\psi_0| \text{ for the initial state } |\psi_0\rangle$$

$$\text{Tr } P_j |\psi_1\rangle\langle\psi_1| \text{ for the initial state } |\psi_1\rangle$$

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which is equal to $\text{Tr } P_j \rho$ for $\rho = p_0 |\psi_0\rangle\langle\psi_0| + p_1 |\psi_1\rangle\langle\psi_1|$.

The postmeasurement state given outcome j is

$$\frac{P_j |\psi_0\rangle\langle\psi_0| P_j}{\text{Tr } P_j |\psi_0\rangle\langle\psi_0|} \quad \text{for the initial state } |\psi_0\rangle$$

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$$\rho_0 \times \text{Pr}(j | |\psi_0\rangle) + \rho_1 \times \text{Pr}(j | |\psi_1\rangle)$$

$$= \frac{\rho_0 \times P_j |\psi_0\rangle\langle\psi_0| P_j + \rho_1 P_j |\psi_1\rangle\langle\psi_1| P_j}{\rho_0 \times \text{Tr } P_j |\psi_0\rangle\langle\psi_0| + \rho_1 \text{Tr } P_j |\psi_1\rangle\langle\psi_1|} = \frac{P_j \rho P_j}{\text{Tr } P_j \rho}$$

Remarks:

1. A density matrix corresponds to a pure state if and only if it is rank-1 (in which case also a projector).
2. The overall phase of a vector disappears when we calculate its density matrix, another way to see why it is irrelevant.

Example 1 (continued)

Can the two mixtures:

$|0\rangle$ with prob $2/3$, $|+\rangle$ with prob $1/3$

$\begin{bmatrix} 0.23 \\ -0.97 \end{bmatrix}$ with prob 0.13 , $\begin{bmatrix} 0.97 \\ 0.23 \end{bmatrix}$ with prob 0.87

be distinguished by operating and measuring the given 2-dim system?

Vote: (a) yes, (b) no.

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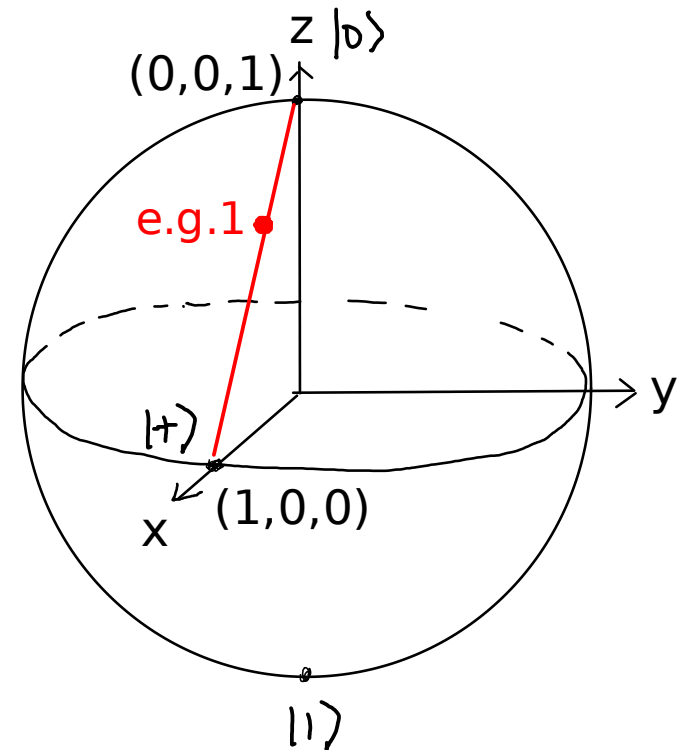
Answer: no. The subsequent density matrices and measurement outcomes only depend on ρ , so, all decompositions of ρ into a convex sum of states are indistinguishable.

Bloch sphere revisited

For a 2-dimensional quantum system, any density matrix can be written as: $\rho = \frac{1}{2} \{ I + aX + bY + cZ \}$.

Exercise:

Show that $\rho \geq 0$ iff $a^2 + b^2 + c^2 \leq 1$
with $\text{rank}(\rho) = 1$ iff $a^2 + b^2 + c^2 = 1$.



Example: classical marginal distribution

A classical random variable U with range $\{1, \dots, d\}$ can be represented as a density matrix $\sum_{u=1}^d p_u |u\rangle\langle u|$, $\{|u\rangle\}$ basis.

Classical information is represented by diagonal density matrices !

Let ρ denote the density matrix for a joint distribution on XY and carried by systems AB :

$$\rho = \sum_{x=1}^{d_A} \sum_{y=1}^{d_B} P_{xy} |x\rangle\langle x|_A \otimes |y\rangle\langle y|_B$$

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Then,

$$\begin{aligned} \text{tr}_B \rho &= \sum_{i=1}^{d_B} (\mathbb{I} \otimes \langle i|) \rho (\mathbb{I} \otimes |i\rangle) \\ &= \sum_{i=1}^{d_B} (\mathbb{I} \otimes \langle i|) \sum_{x=1}^{d_A} \sum_{y=1}^{d_B} P_{xy} |x\rangle\langle x|_A \otimes |y\rangle\langle y|_B (\mathbb{I} \otimes |i\rangle) \end{aligned}$$

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turns B into 1-dim sys

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which is indeed the margin distribution X !

Density matrices on a composite system

Consider a bipartite system AB , with $\dim d_A, d_B$.

The most general state on AB is a mixture of pure states on AB , each is $d_A d_B$ -dimensional.

Reading exercise for W25?

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These states arises if Alice draws x with prob p_x records x in system A (classical information), prepares ρ_x on system B and gives it to Bob.

Exercise:

Show that if $\rho = \sum_x p_x |x\rangle\langle x| \otimes \rho_x$

then $\text{tr}_A \rho = \underbrace{\sum_x p_x \rho_x}$

density matrix for
the mixture ρ_x w.p. p_x .