

Interactions of charged particles with graphene

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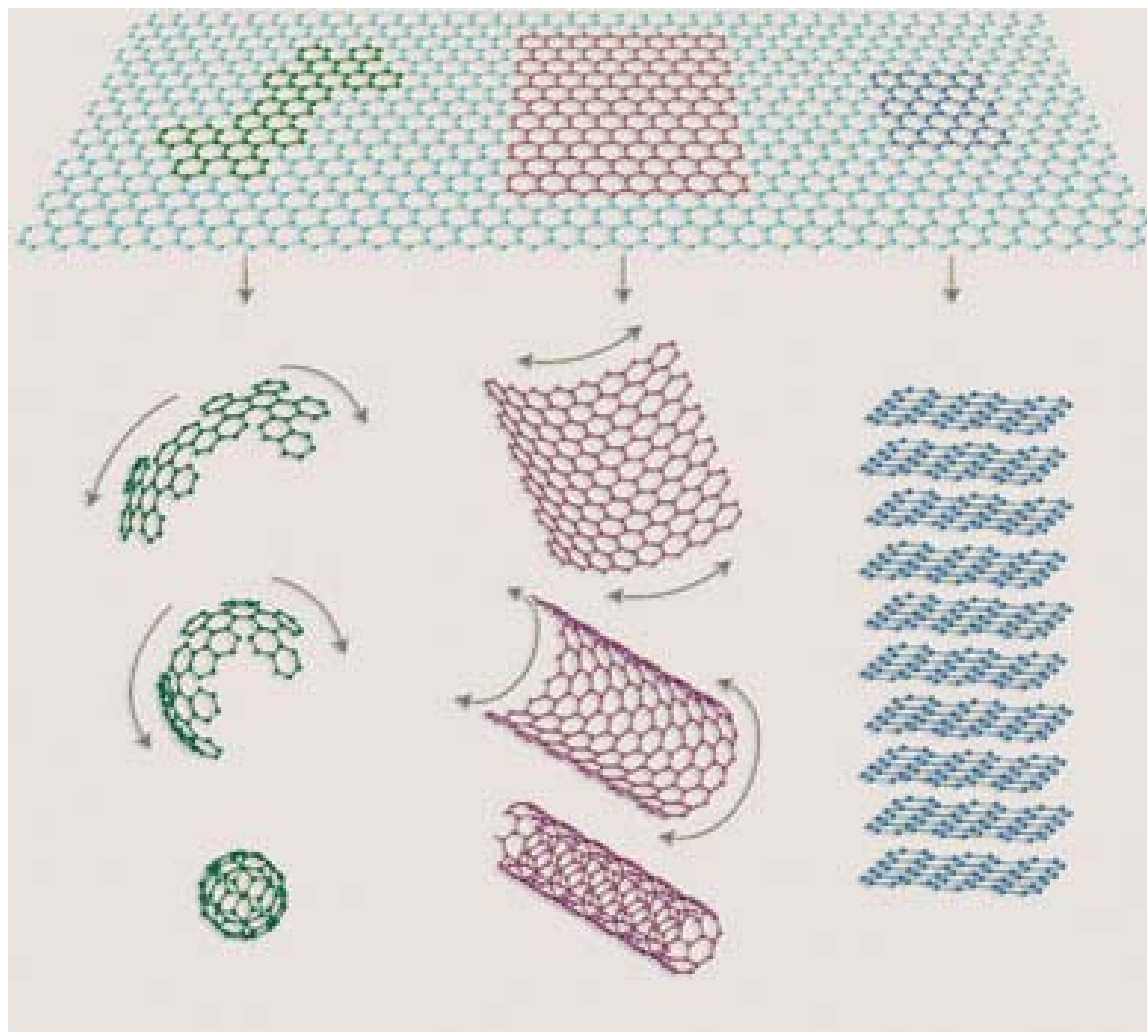
Support:  **NSERC
CRSNG**

Outline

- ❑ Introduction
- ❑ Ion interactions with graphene
 - Nonlinear screening
 - Image potential
 - Stopping force
 - Wake effect
- ❑ Plasmon excitations by electron beams
 - HREELS, energy $E \sim 20$ eV, substrate phonons
 - EELS, energy $E \sim 100$ keV, layered electron gas
- ❑ Outlook

Graphene as building block of carbon nanostructures

graphene



Fullerene, C₆₀

nanotube, CNT

graphite, HOPG

Why study carbon nanostructures ?

- Physical properties:
 - Electrical, mechanical, thermal
 - Dependent on: molecular structure, dielectric environment, local modification
- Applications:
 - Nanoelectronic devices
 - Biochemical sensors
 - New composite materials
 - Ion storage (H, Li)
 - Nanoelectromechanical systems (NEMS)

Electronic response to external charged particles (regime of electronic stopping for moving charges)

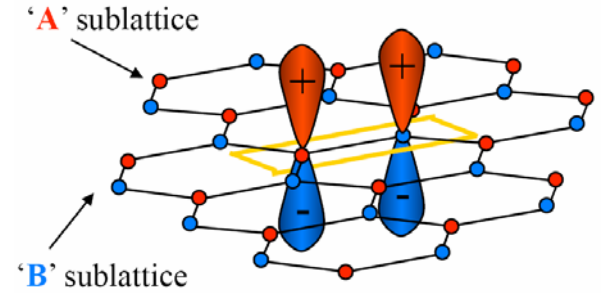
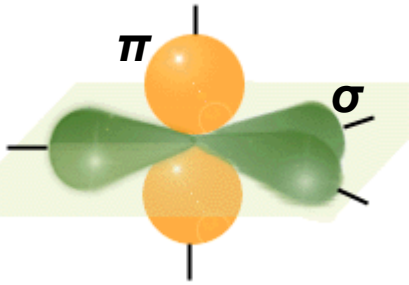
□ Electrons

- EELS in STEM: plasmon excitations of σ and π electrons in CNTs & graphene
- HREELS: plasmon excitations of π electrons in graphene
- image potential states: CNTs & graphene

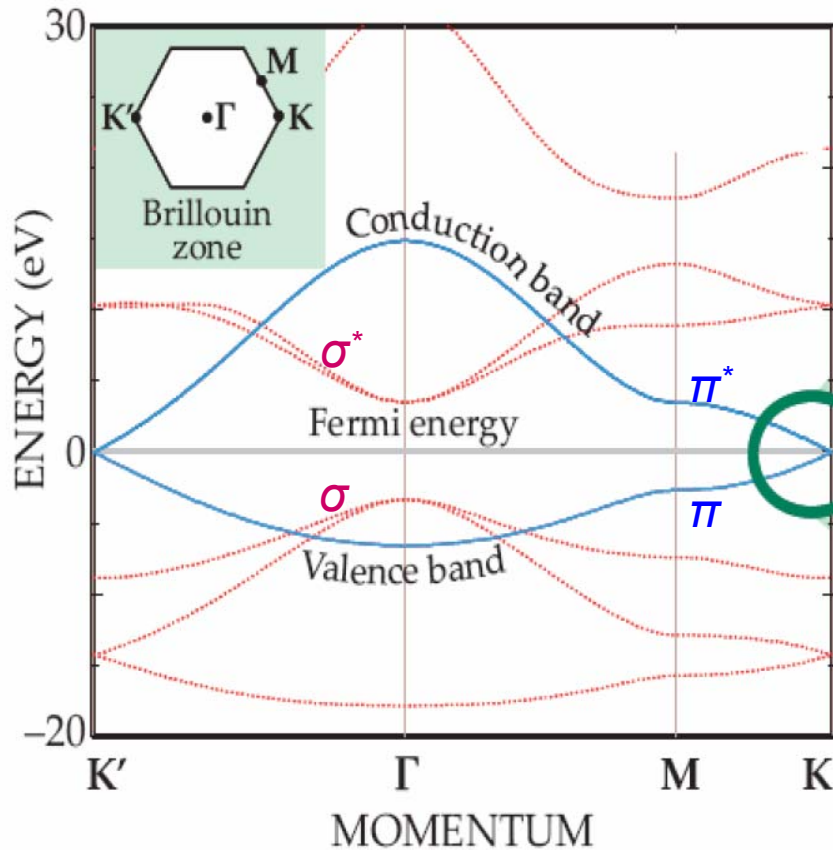
□ Ions

- ion channeling in CNTs
- grazing scattering of ions on graphene
- friction forces on slowly moving ions
- static screening of charged impurities

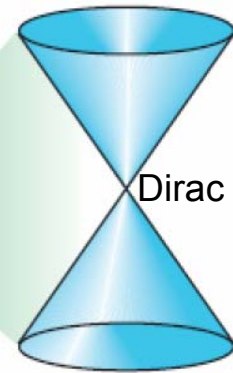
Electronic structure of graphene



Band structure



low-energy excitations involve π bands only

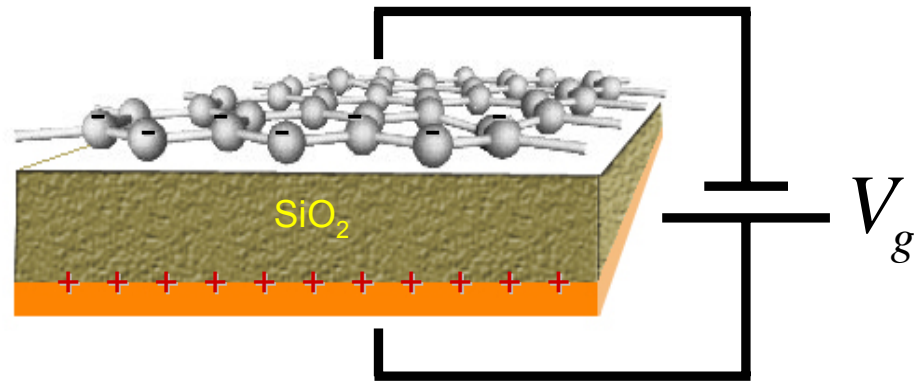
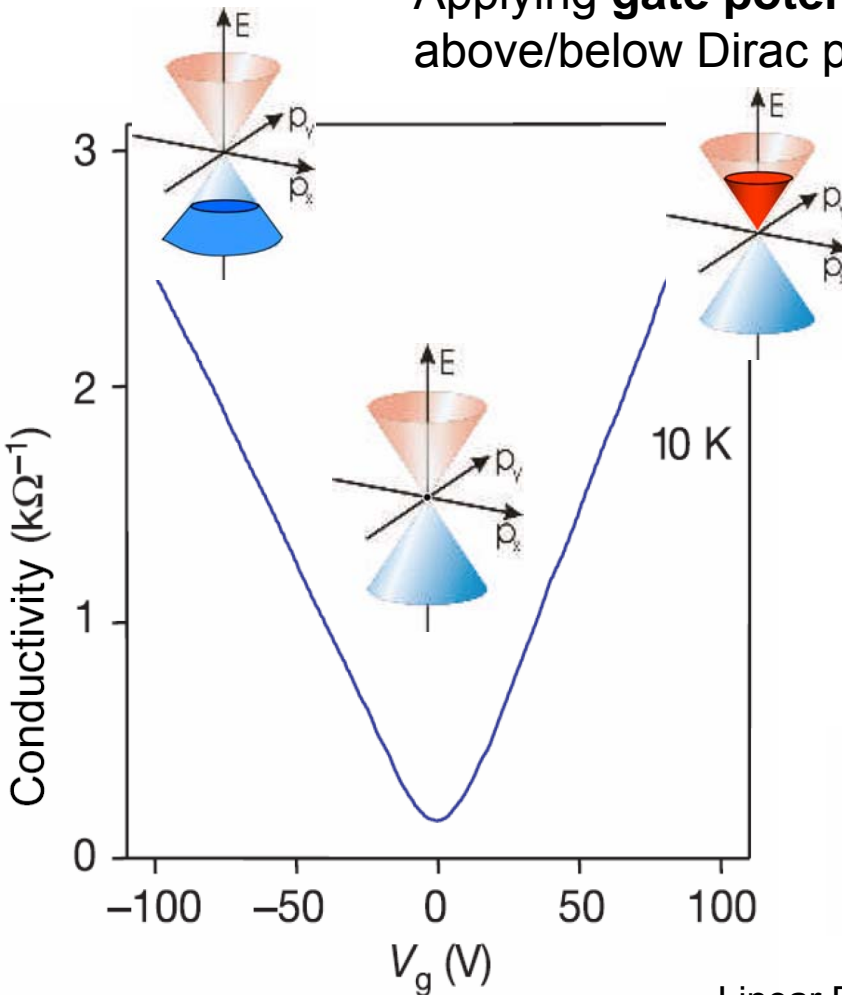


Dirac point

$$\epsilon_{\mathbf{k}} = \hbar v_F \|\mathbf{k}\|$$

Ambipolar electric field effect in single-layer graphene on top of an oxidized Si wafer

Applying **gate potential** V_g shifts Fermi energy E_F above/below Dirac point (electron or hole doping)



Equilibrium charge carrier density vs chemical potential μ may be >0 and <0

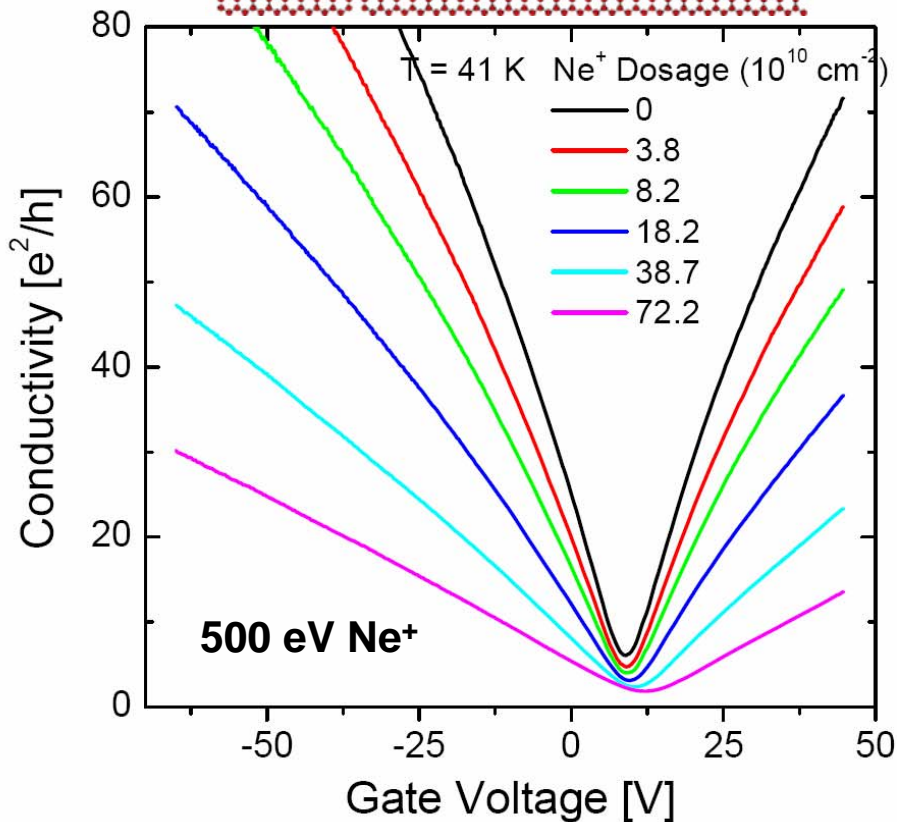
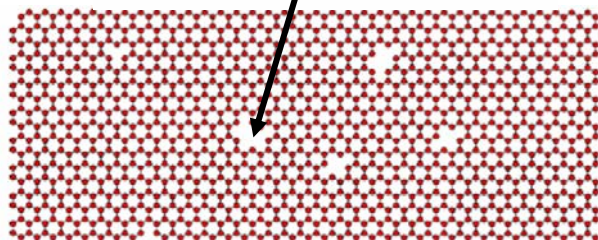
$$n(\mu) = \int_0^{\infty} d\varepsilon \rho(\varepsilon) \left[\frac{1}{1 + e^{\beta(\varepsilon - \mu)}} - \frac{1}{1 + e^{\beta(\varepsilon + \mu)}} \right]$$

Linear DOS: $\rho(\varepsilon) \approx \frac{g_s g_v}{2\pi} \frac{|\varepsilon|}{(\hbar v_F)^2}$, Fermi speed: $v_F \approx \frac{c}{300}$

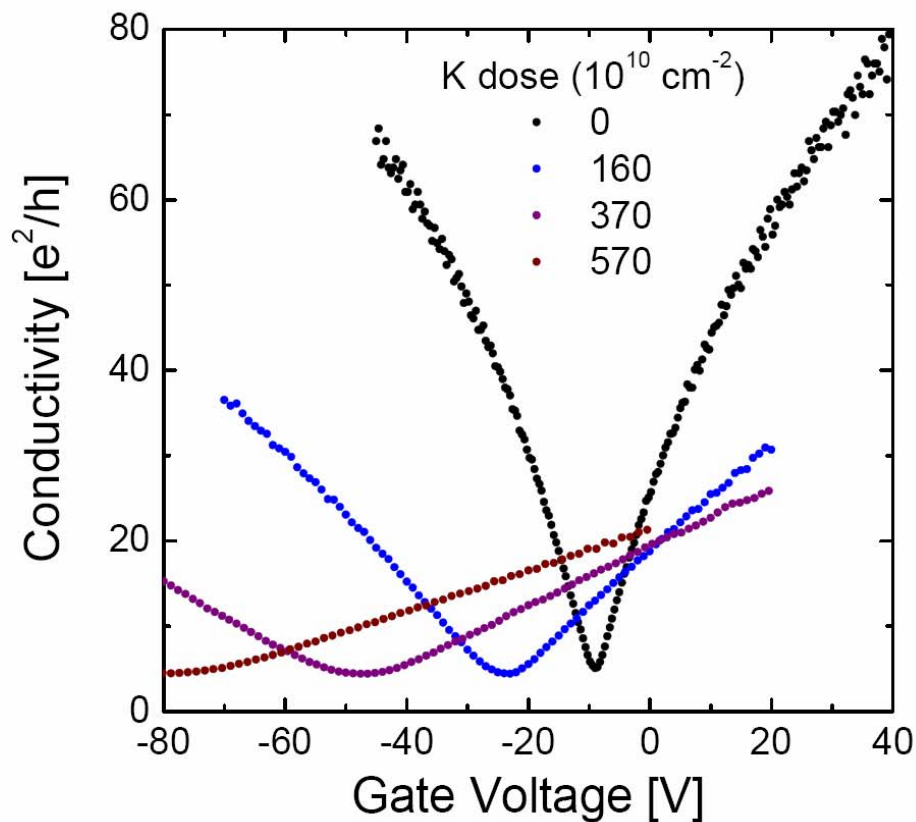
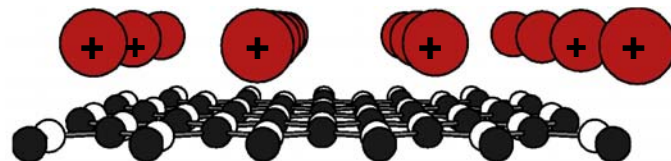
Ion bombardment vs ion adsorption on graphene

J.H. Chen *et al.*, *Phys. Rev. Lett.* (2009), *Nature Phys.* (2008)

Ne⁺ irradiation
(lattice defects)

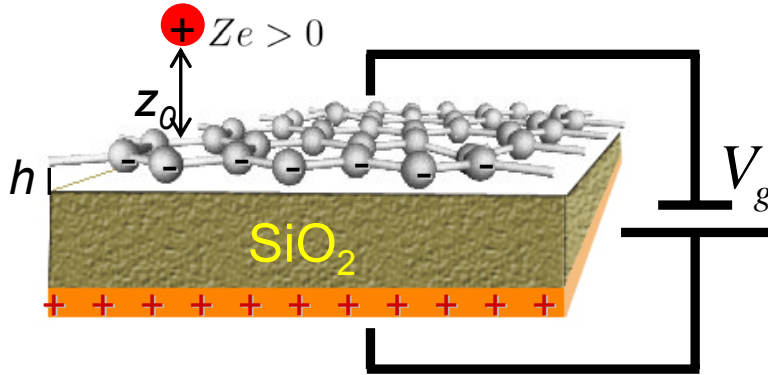


K doping
(charged impurities)



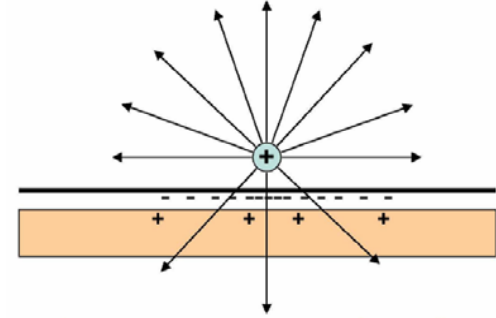
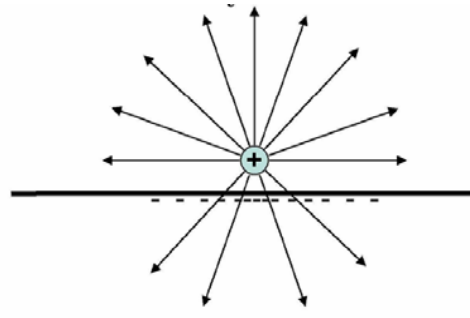
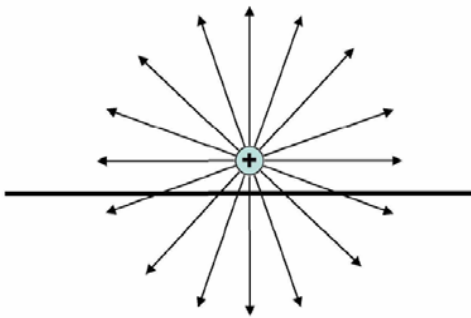
Nonlinear screening of external charge by graphene

M. Ghaznavi *et al.* *Phys. Rev. B* 81 (2010) 085416



Study several effects:

- doping via gate potential
- finite temperature
- exchange interaction effects
- distance z_0 from graphene
- size of gap h to the substrate



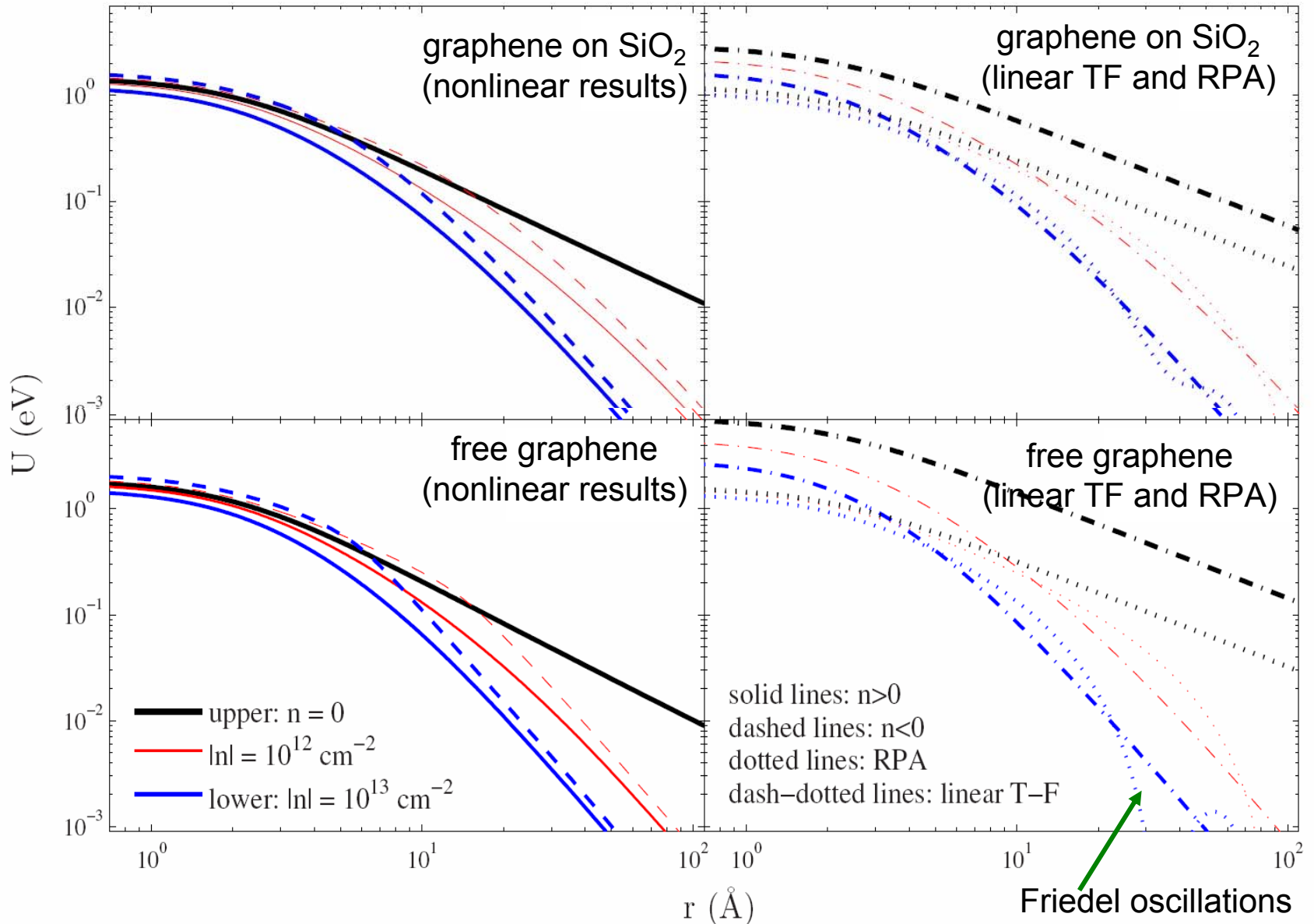
Thomas-Fermi model at finite T

$$U(\mathbf{r}) = U_{\text{ext}}(\mathbf{r}) - e^2 \int d^2\mathbf{r}' [n(\mu + U(\mathbf{r}')) - n(\mu)] \left[\frac{1}{\|\mathbf{r} - \mathbf{r}'\|} - \frac{\epsilon_s - 1}{\epsilon_s + 1} \frac{1}{\sqrt{(\mathbf{r} - \mathbf{r}')^2 + 4h^2}} \right]$$

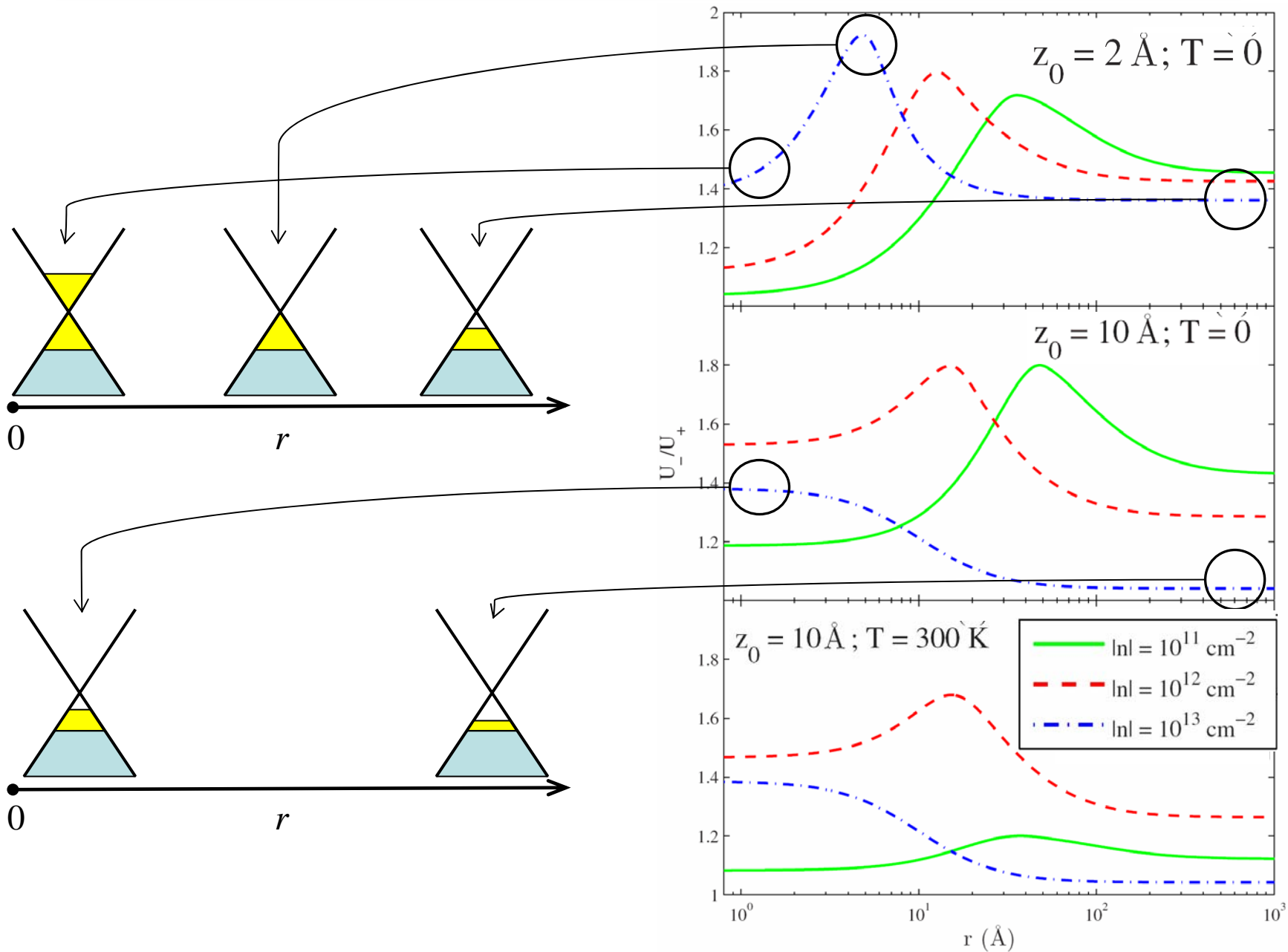
Self-consistent solution for screened potential $U_{\pm}(r)$ in graphene when $\mu \geq 0$

Screened potential in graphene due to external charge

M. Ghaznavi *et al.* *Phys. Rev. B* 81 (2010) 085416



Ratio $U_-(r)/U_+(r)$ of screened potentials for $n \lesssim 0$



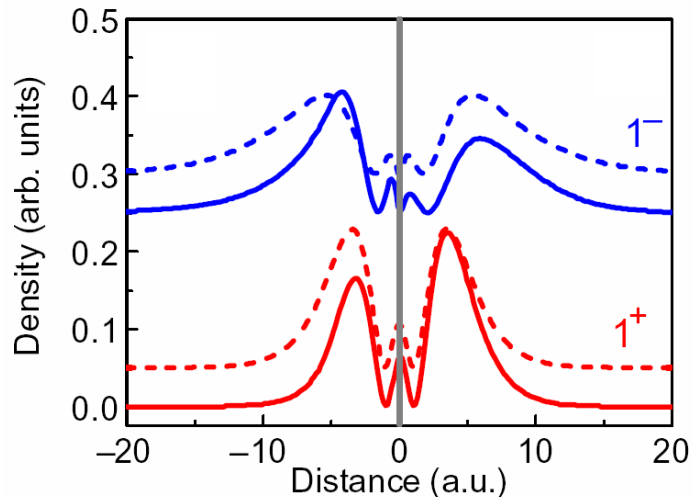
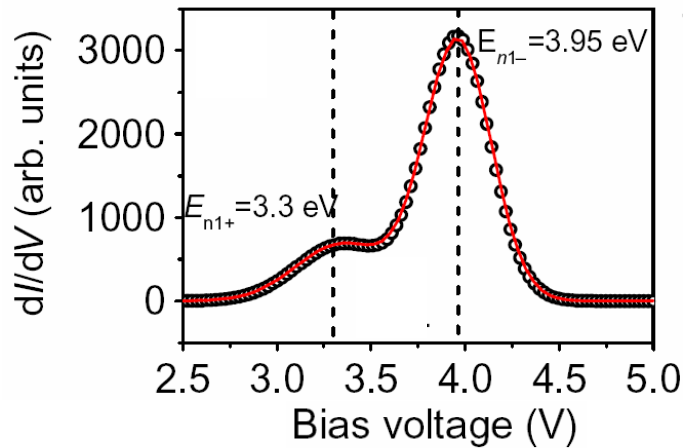
Nonlinear image interaction for external charge

M. Ghaznavi *et al. Phys. Rev. B* 81 (2010) 085416

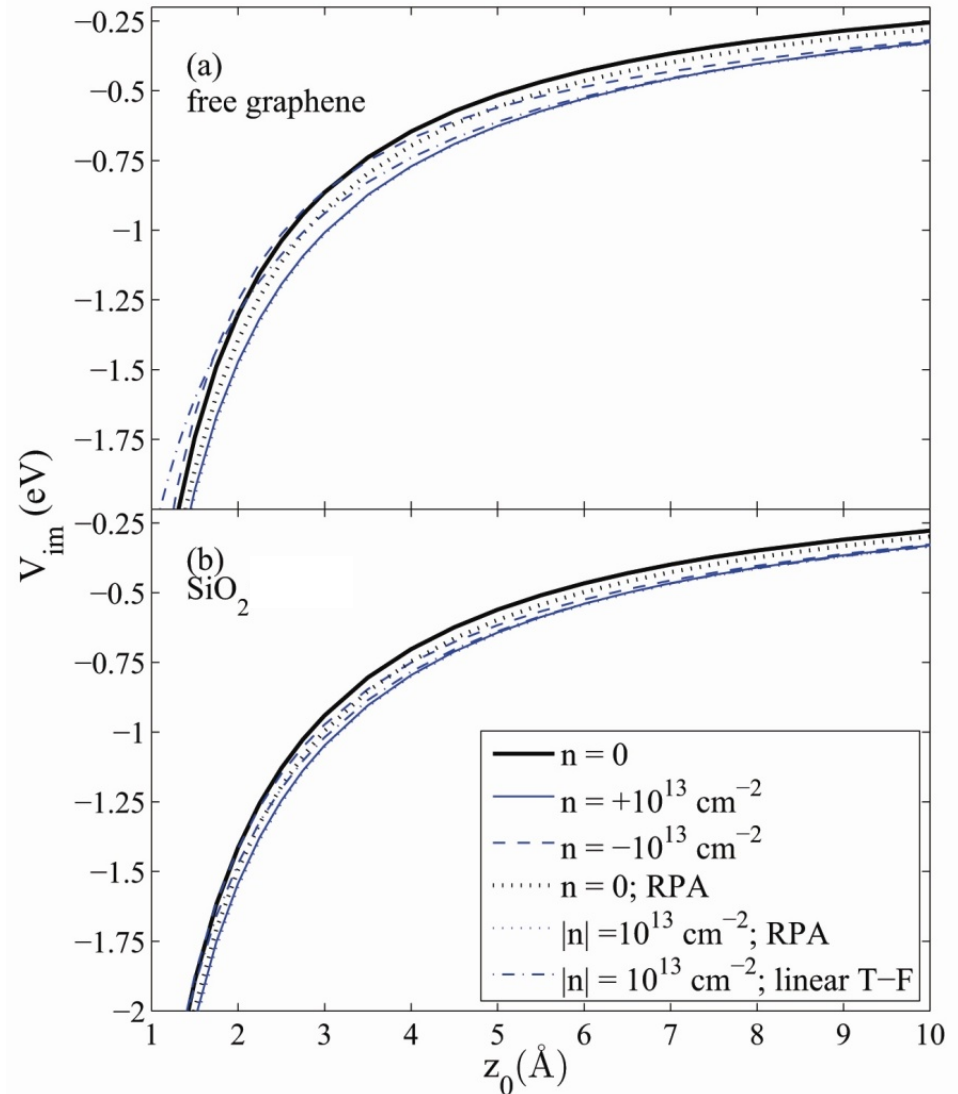
Observation of image potential states on epitaxial graphene:

S. Bose *et al. New J. Phys.*

12 (2010) 023028

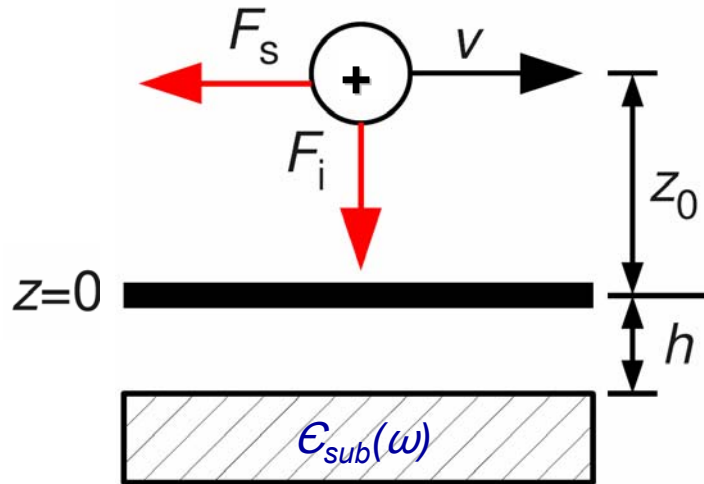


Nonlinear image potential



Interaction of moving external charge with graphene

K.F. Allison *et al.*, *Phys. Rev.* 80 (2009), *Nanotechnology* 21 (2010) 134017



Study several effects:

- doping via gate potential
- damping via Mermin approach
- local field effects
- phonons in polar substrate
- size of gap to the substrate

Dielectric function of graphene + substrate: Random Phase Approximation

$$\epsilon(q, \omega) = \left[1 - \frac{\epsilon_{\text{sub}}(\omega) - 1}{\epsilon_{\text{sub}}(\omega) + 1} e^{-2qh} \right]^{-1} + \frac{2\pi e^2}{q} \Pi_{\text{gra}}(q, \omega; \gamma)$$

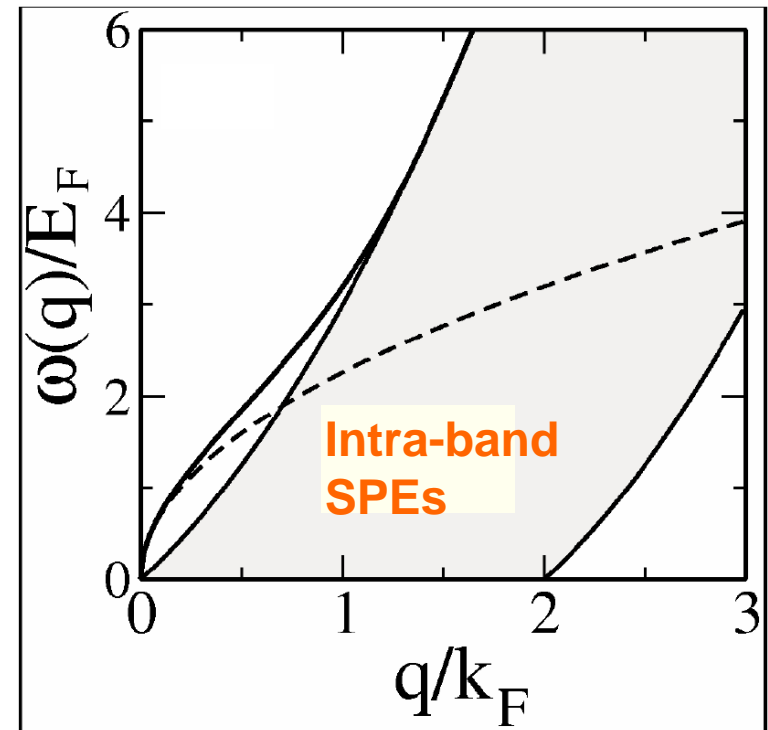
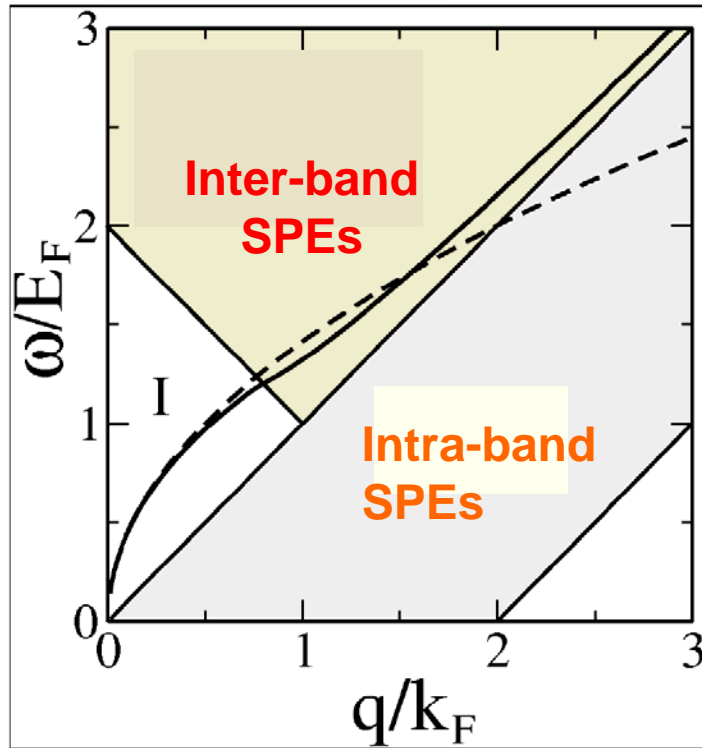
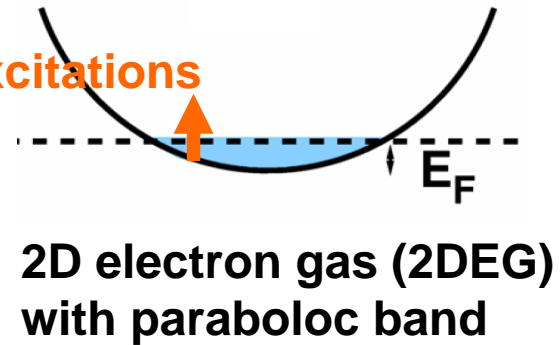
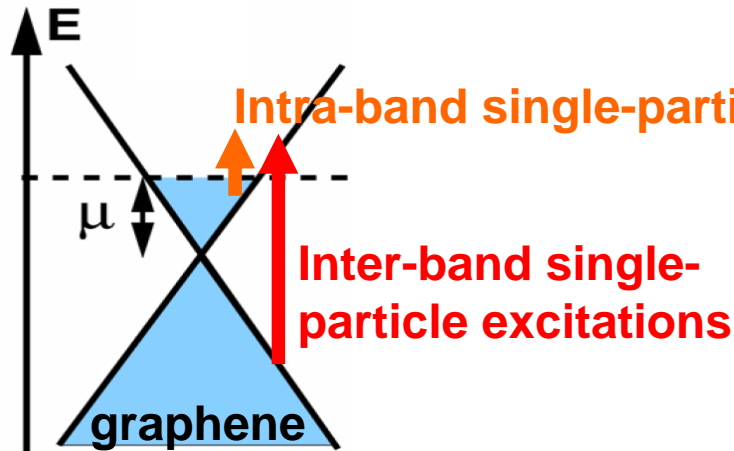
Stopping force

$$F_s = \frac{2}{\pi} \frac{Z^2 e^2}{v} \int_0^\infty dq e^{-2qz_0} \int_0^{qv} d\omega \frac{\omega}{\sqrt{q^2 v^2 - \omega^2}} \Im \left[\frac{1}{\epsilon(q, \omega)} \right]$$

Image force

$$F_i = \frac{2}{\pi} Z^2 e^2 \int_0^\infty dq q e^{-2qz_0} \int_0^{qv} \frac{d\omega}{\sqrt{q^2 v^2 - \omega^2}} \Re \left[\frac{1}{\epsilon(q, \omega)} - 1 \right]$$

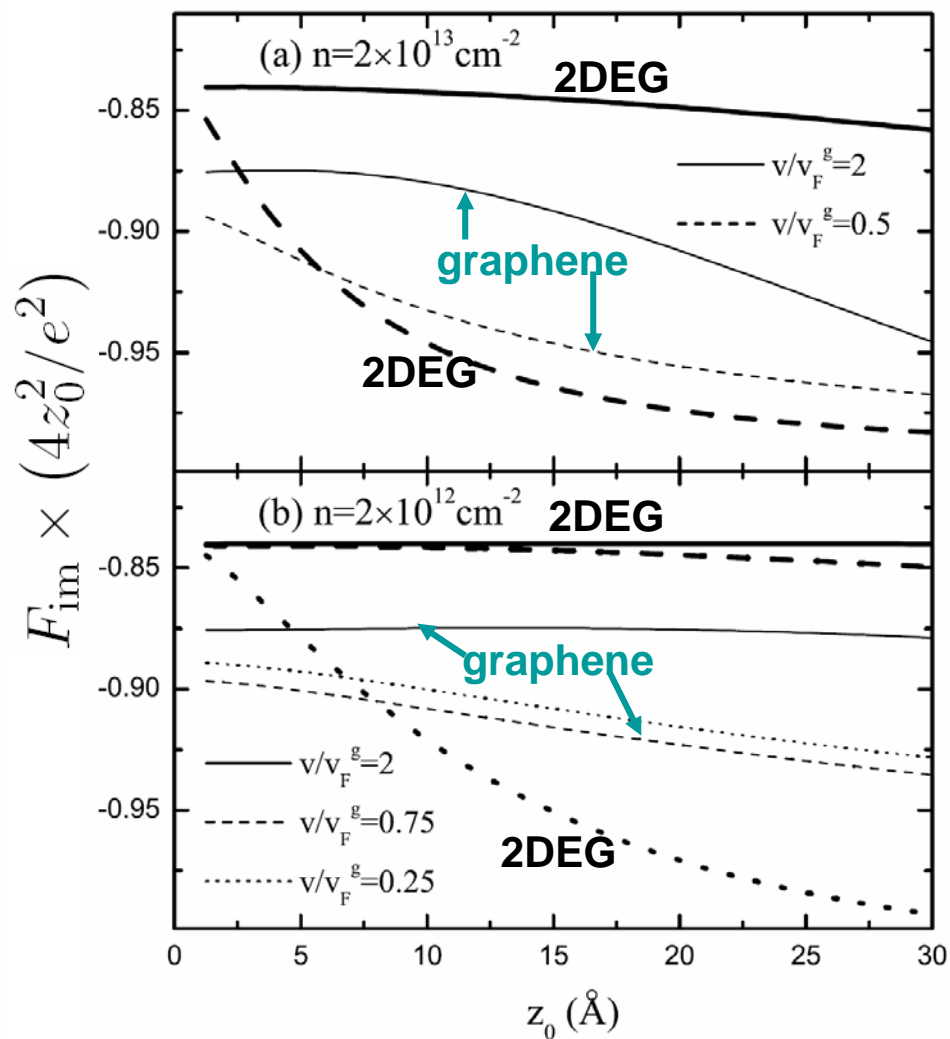
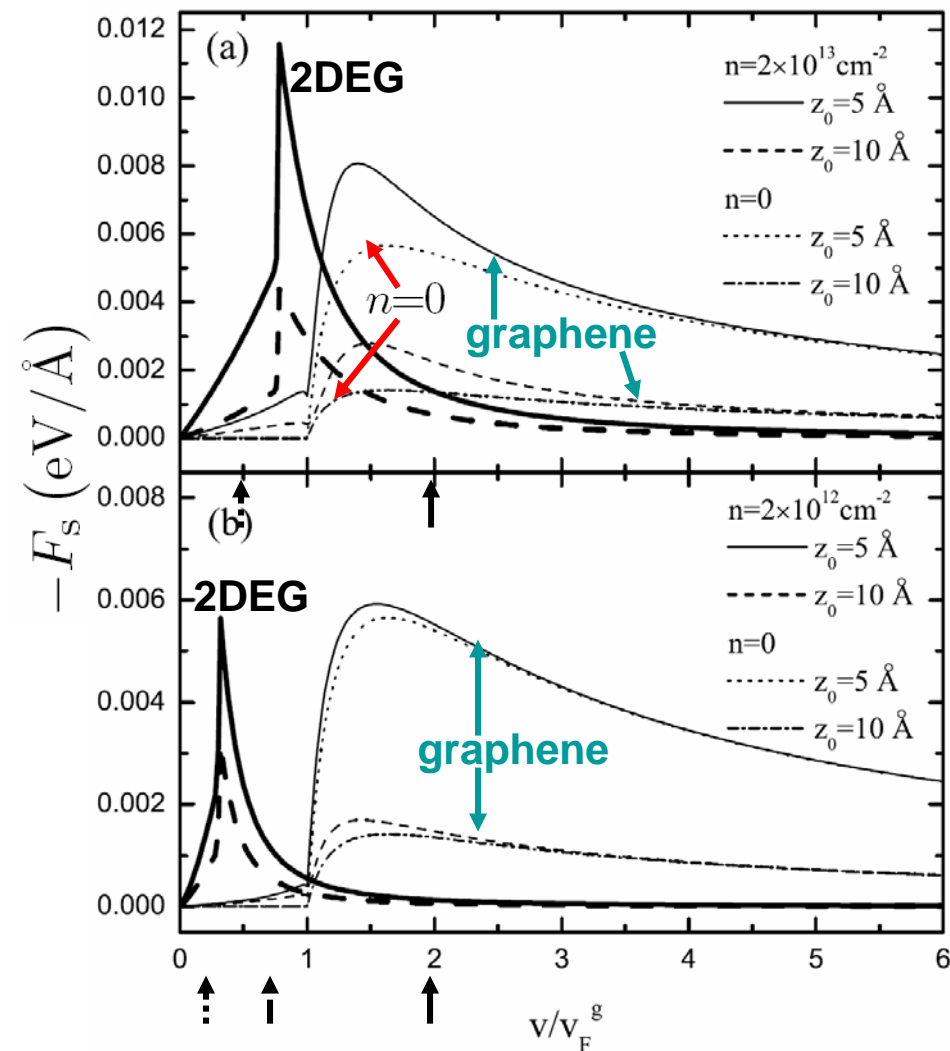
Loss Function $-\Im \{1/\epsilon(q, \omega)\}$ for graphene and 2DEG



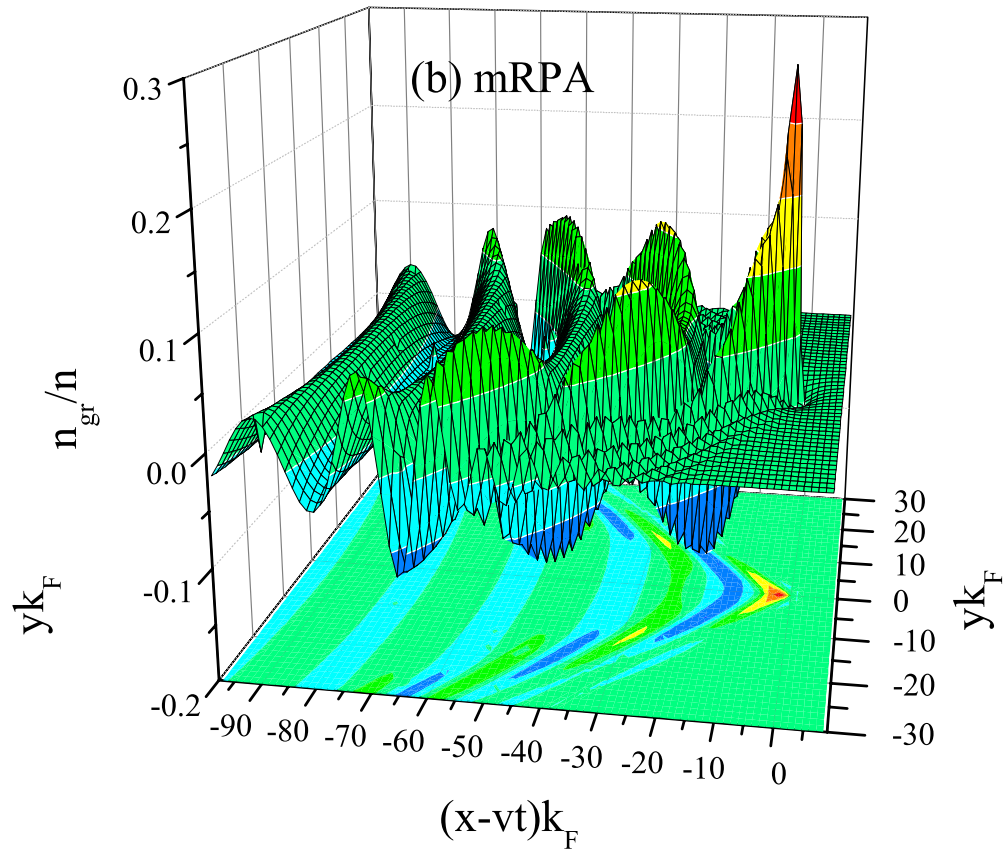
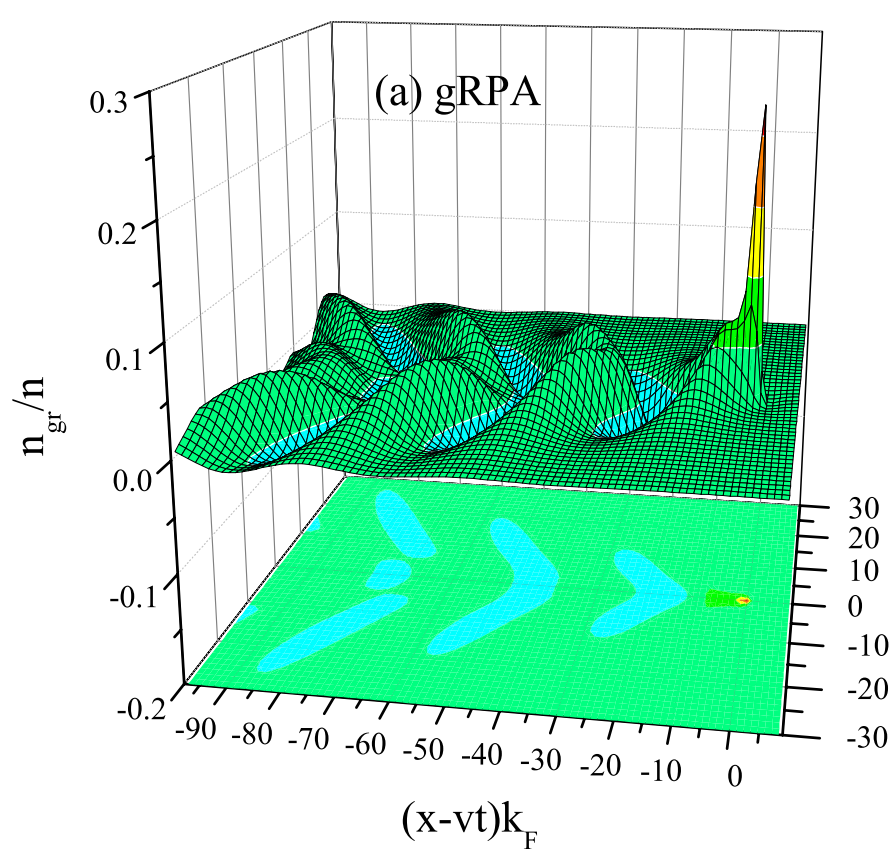
Stopping and image forces on a point charge moving over epitaxial graphene on SiC and 2DEG (Ag on Si)

Stopping force vs. speed

Image force (normalized) vs. distance



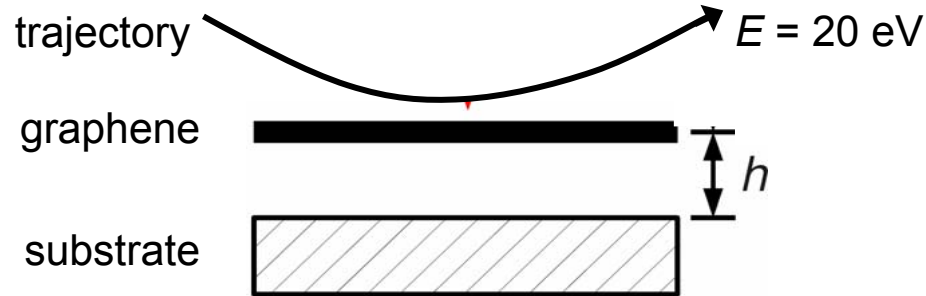
Wake effect due to charge moving at speed $v = 4v_F$ and distance $z_0 = 20 \text{ \AA}$ from free graphene & 2DEG with doping equilibrium density $n = 10^{13} \text{ cm}^{-2}$



Fermi wavenumber in graphene: $k_F = \sqrt{\pi n}$

HREEL spectra for graphene on a substrate

K.F. Allison *et al.*, *Nanotechnology* 21 (2010) 134017



Total energy loss:
$$E_{\text{loss}} = - \int_{-\infty}^{\infty} dt \int d^2\mathbf{r} \int_{-\infty}^{\infty} dz \rho_{\text{ext}}(\mathbf{r}, z, t) \frac{\partial}{\partial t} \Phi_{\text{ind}}(\mathbf{r}, z, t)$$

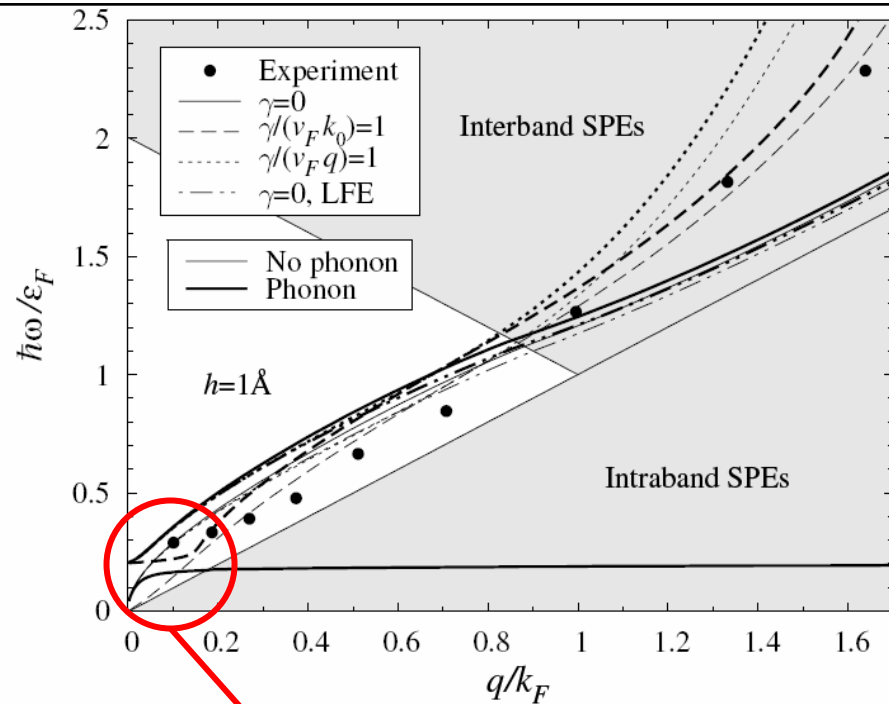
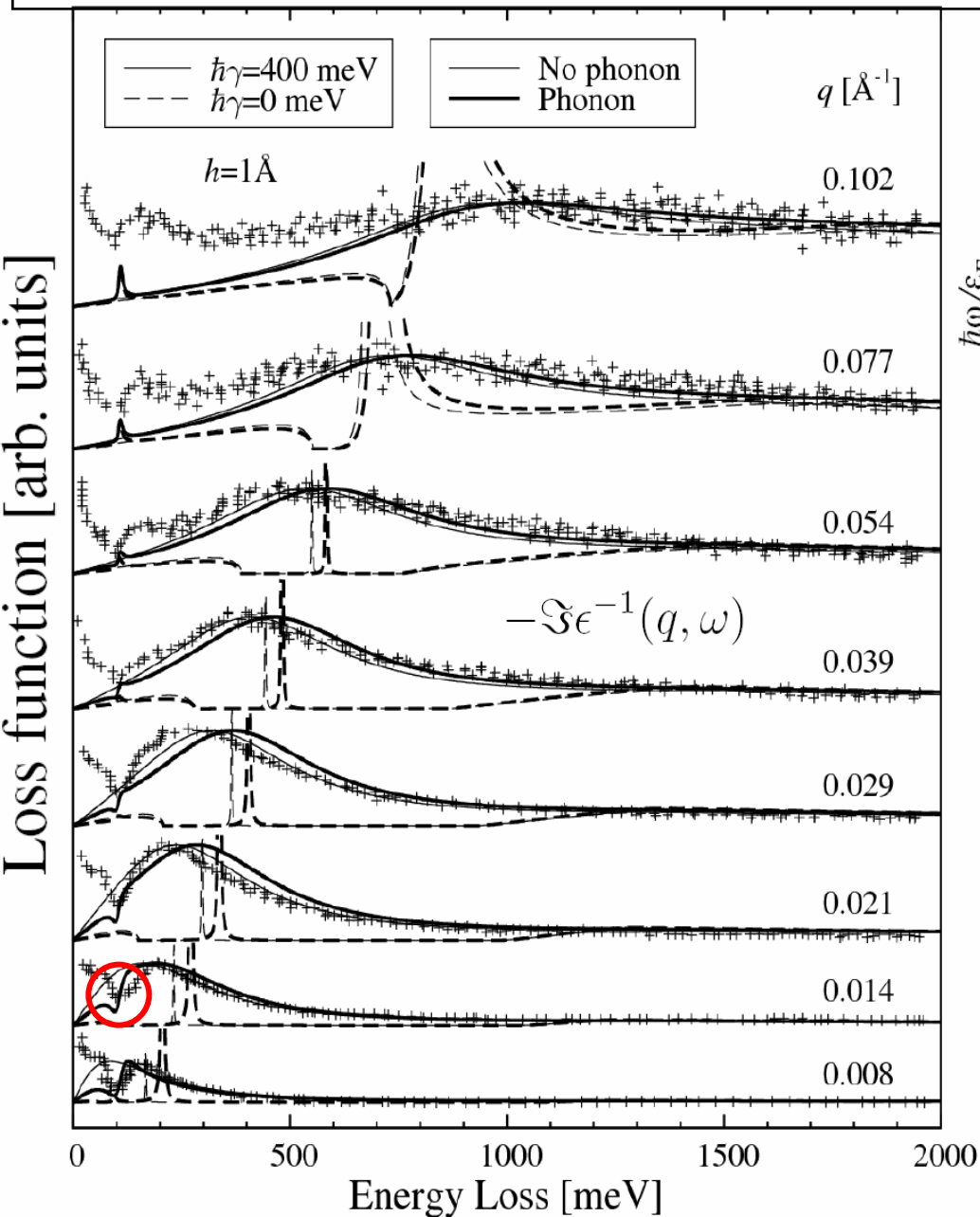
$$= \int_0^{\infty} d\omega \omega \int \frac{d^2\mathbf{q}}{2\pi^2} |S(\mathbf{q}, \omega)|^2 \Im \left[\frac{-1}{\epsilon(q, \omega)} \right] \rightarrow \text{Loss function}$$

Projectile structure factor:
$$S(\mathbf{q}, \omega) = \int_{-\infty}^{\infty} dt e^{i\omega t} \int d^2\mathbf{r} e^{-i\mathbf{q}\cdot\mathbf{r}} \int_{-\infty}^{\infty} dz e^{-q|z|} \rho_{\text{ext}}(\mathbf{r}, z, t)$$

$$\epsilon(q, \omega) = \left[1 - \frac{\epsilon_{\text{surf}}(q, \omega) - 1}{\epsilon_{\text{surf}}(q, \omega) + 1} e^{-2qh} \right]^{-1} + \frac{2\pi e^2}{q} \chi_{\text{gr}}(q, \omega)$$

Surface dielectric function:
$$\epsilon_{\text{surf}}(q, \omega) = \left[\frac{q}{\pi} \int_{-\infty}^{\infty} \frac{dq_z}{(q^2 + q_z^2) \epsilon_{\text{sub}}(\sqrt{q^2 + q_z^2}, \omega)} \right]^{-1}$$

HREELS spectra for graphene on SiC, exper. by Liu *et al.* PRB 78 (2008) 201403



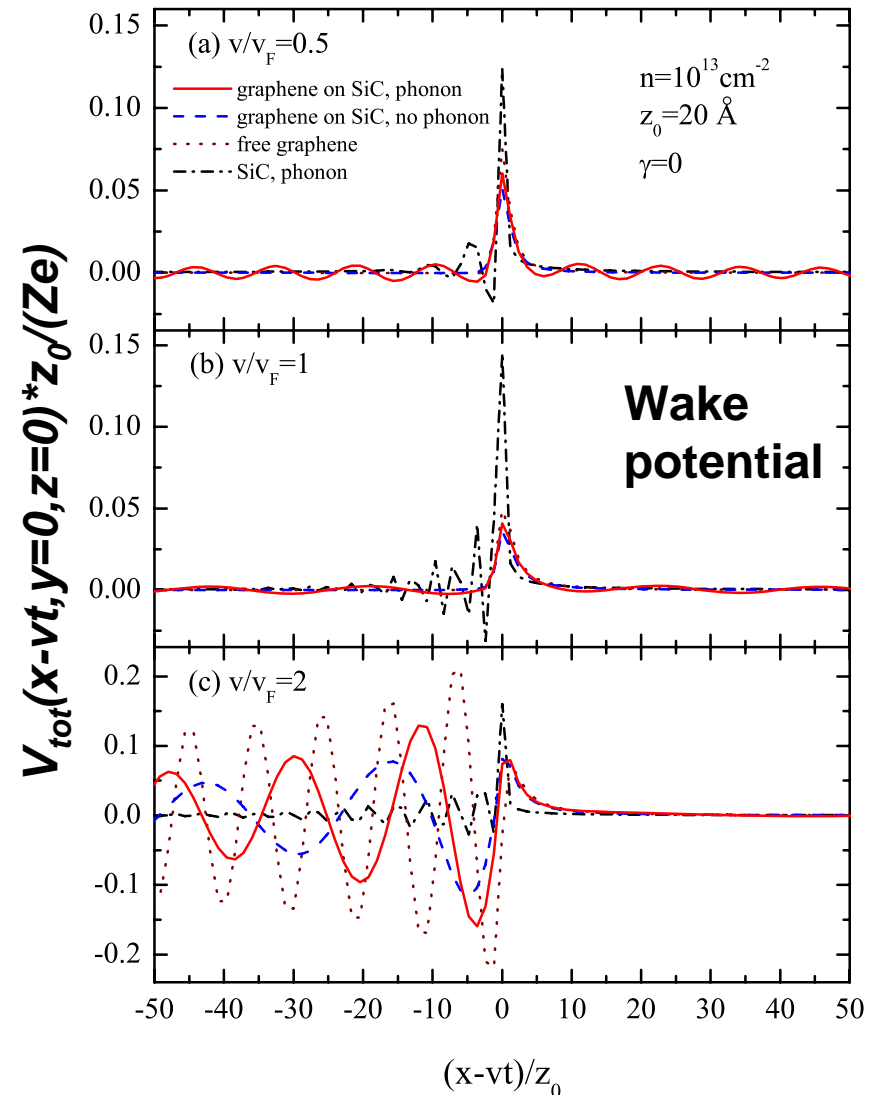
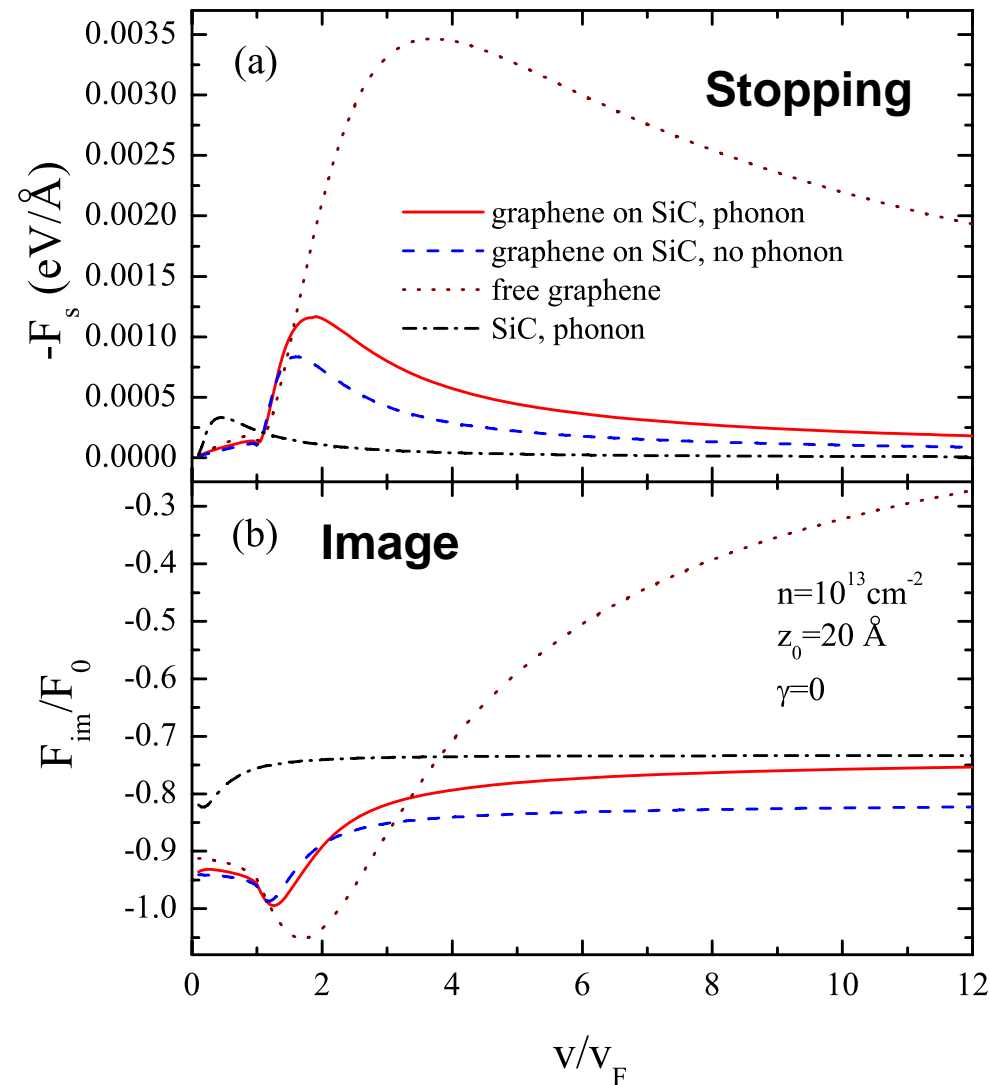
Plasmon coupling with substrate phonon

$$\epsilon(q, \omega) \approx 1 + \epsilon_{\text{sub}}(\omega) - 4 \frac{e^2}{\hbar} \frac{v_F q}{\omega^2} \sqrt{\pi n} = 0$$

$$\epsilon_{\text{sub}}(\omega) = \epsilon_{\infty} + (\epsilon_0 - \epsilon_{\infty}) \frac{\omega_{\text{TO}}^2}{\omega_{\text{TO}}^2 - \omega^2}$$

Stopping and image forces on projectile moving over graphene on SiC: effects of substrate phonons

CEPAS proceedings



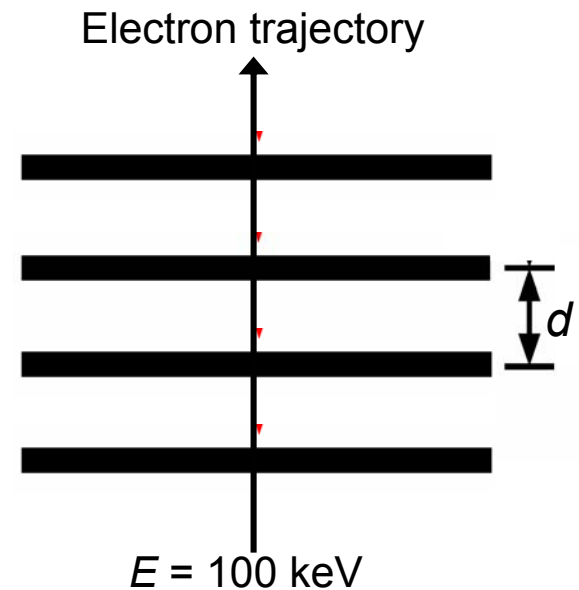
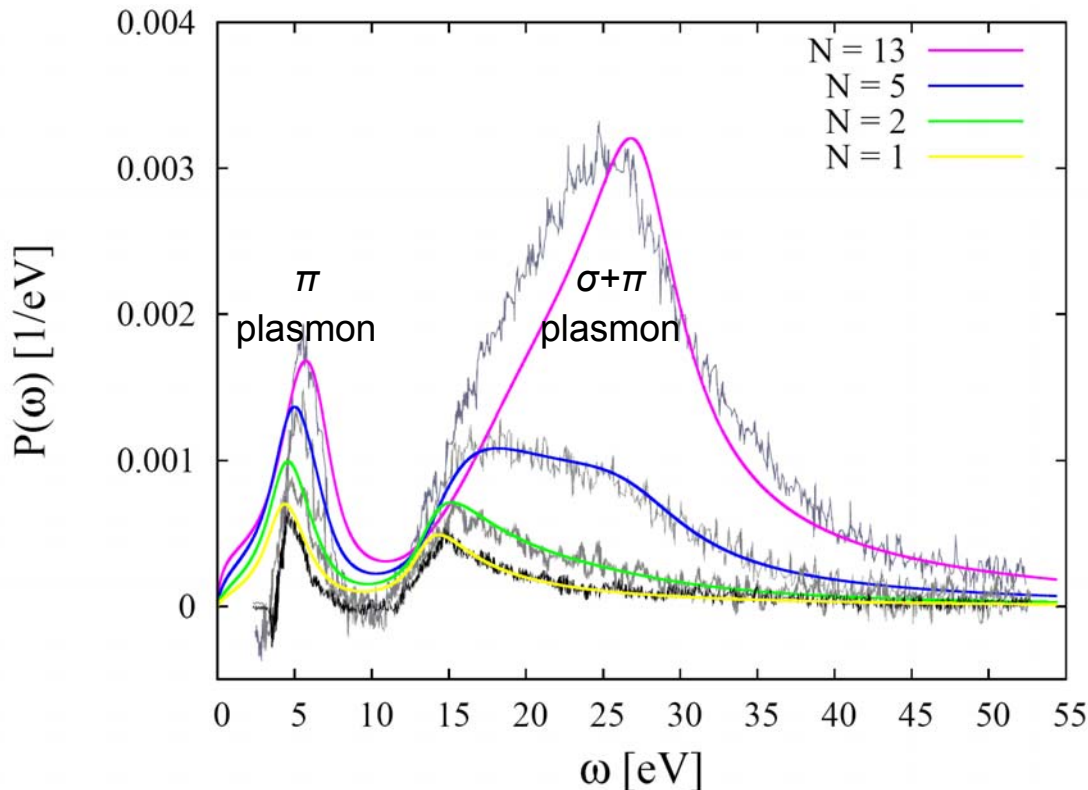
EEL spectra for free-standing N -layer graphene

experiment by: T. Eberlein *et al.*, *Phys. Rev. B* 77 (2008) 233406

$$P(\omega) = \frac{4e^2}{\pi v^2} \int_0^{q_{\max}} \frac{dq q^2}{\left[q^2 + \left(\frac{\omega}{v} \right)^2 \right]^2} \Im \left\{ V(q) \chi(q, \omega) \sum_{n=1}^N \sum_{n'=1}^N e^{-i(n-1)d\omega/v} (\mathcal{M}^{-1})_{nn'} e^{i(n'-1)d\omega/v} \right\}$$

$$\mathcal{M}_{nn'}(q, \omega) = \delta_{nn'} + (1 - \delta_{nn'}) V(q) \chi(q, \omega) e^{-qd|n'-n|}$$

$$\chi(q, \omega) = \chi_0(q, \omega) / [1 + V(q) \chi_0(q, \omega)], \quad V(q) = 2\pi e^2 / q$$

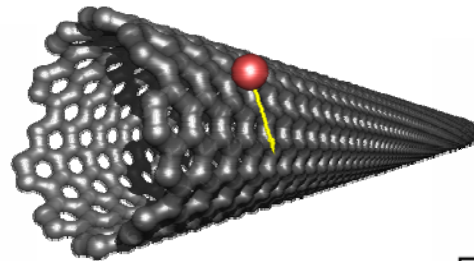


Two-fluid, 2D hydrodynamic model for σ and π electrons: used in single-wall carbon nanotube

experiment: Kramberger *et al.*, *Phys. Rev. Lett.* 100 (2008) 196803

theory: Mowbray *et al.*, *Phys. Rev. B* 82 (2010) 035405

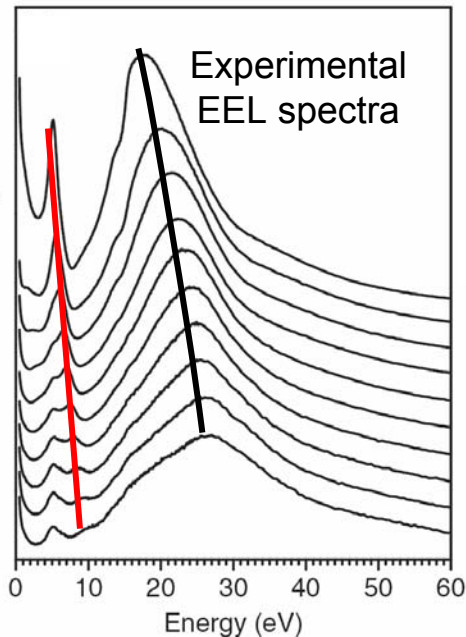
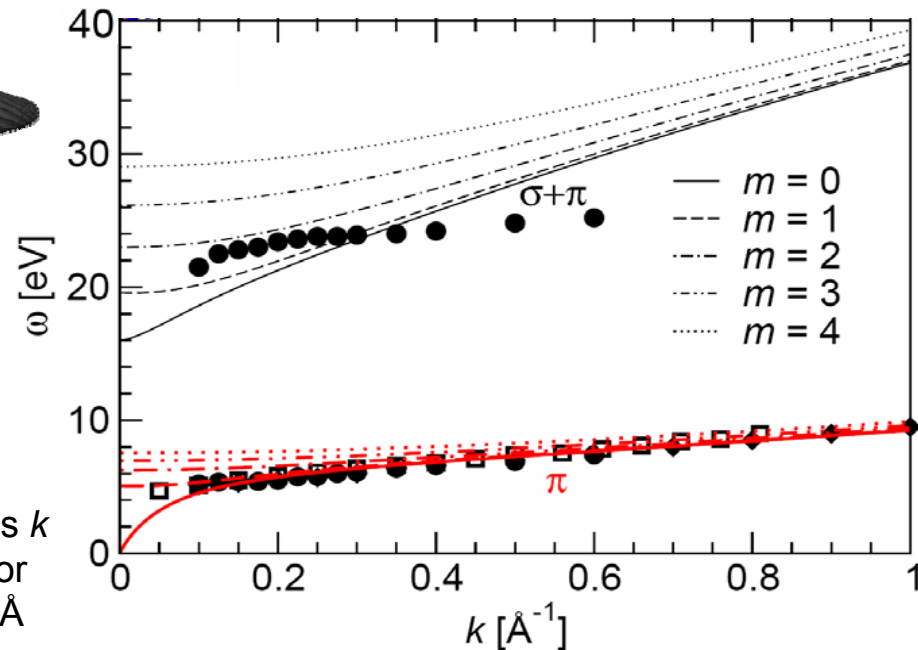
$$\chi_0(q, \omega) = \frac{n_{\pi}^0 q^2 / m_{\pi}^*}{s_{\pi}^2 q^2 + \omega_{\pi r}^2 - \omega(\omega + i\gamma_{\pi})} + \frac{n_{\sigma}^0 q^2 / m_{\sigma}^*}{s_{\sigma}^2 q^2 + \omega_{\sigma r}^2 - \omega(\omega + i\gamma_{\sigma})}$$



k [\AA^{-1}]
0.1
0.2
0.3
0.4
0.5
0.6
0.7
0.8
0.9
1.0

$$q^2 \rightarrow \frac{m^2}{a^2} + k^2$$

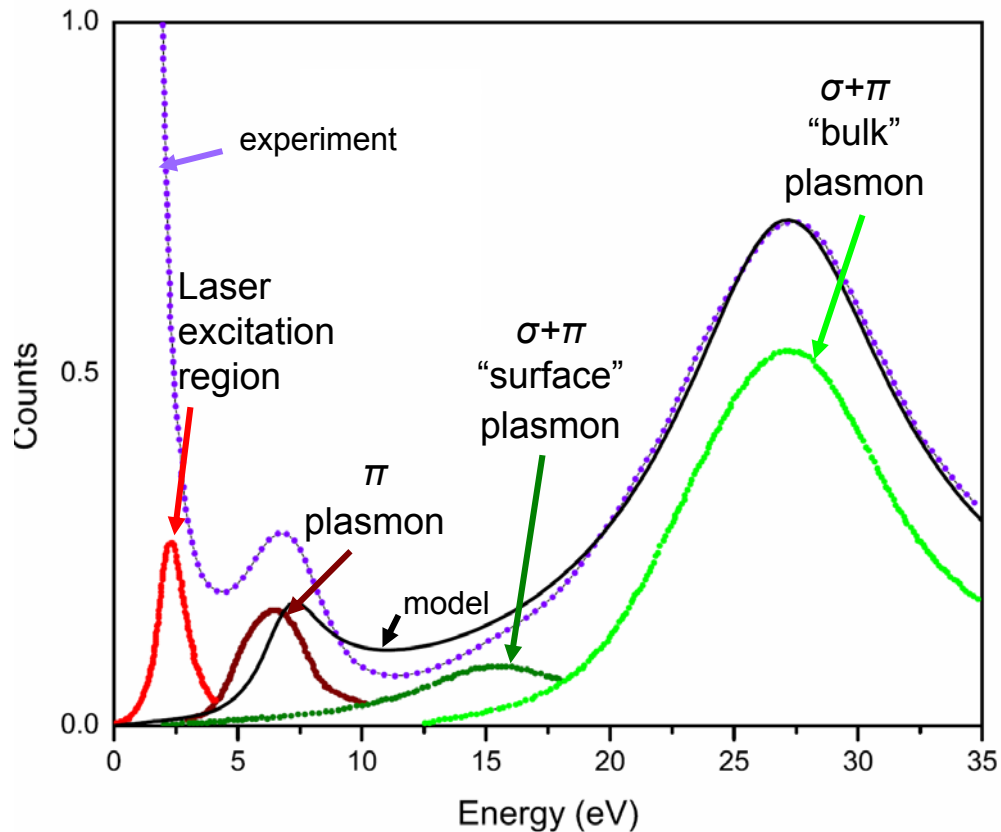
Plasmon dispersions vs k with angular mode m for nanotube radius $a = 7 \text{ \AA}$



EELS for free-standing graphite ($N \gg 1$)

F. Carbone *et al.*, *Chem. Phys. Lett.* 468 (2009) 107

$$P(\omega) \approx N \frac{4e^2}{\pi v^2} \int_0^{q_{\max}} \frac{dq q^2}{\left[q^2 + \left(\frac{\omega}{v} \right)^2 \right]^2} \Im \left[\frac{V(q) \chi_0(q, \omega)}{1 + \frac{\sinh(qd)}{\cosh(qd) - \cos(\omega d/v)} V(q) \chi_0(q, \omega)} \right]$$



Outlook

- Electronic energy loss and image interaction important for charged particle interactions with C-nanostructures: energy deposition & transport
- Plasmon excitations in HREELS and EELS: full theory needed for dielectric response of both σ and π electrons
- Ion charge states during scattering need to be considered
- Effects of dielectric environment, particularly substrate phonons
- Concept of friction and screening of mobile ions in aqueous solution

Thank you for your attention